### Implications of the S<sub>8</sub> tension for decaying dark matter

#### Guillermo Franco Abellán

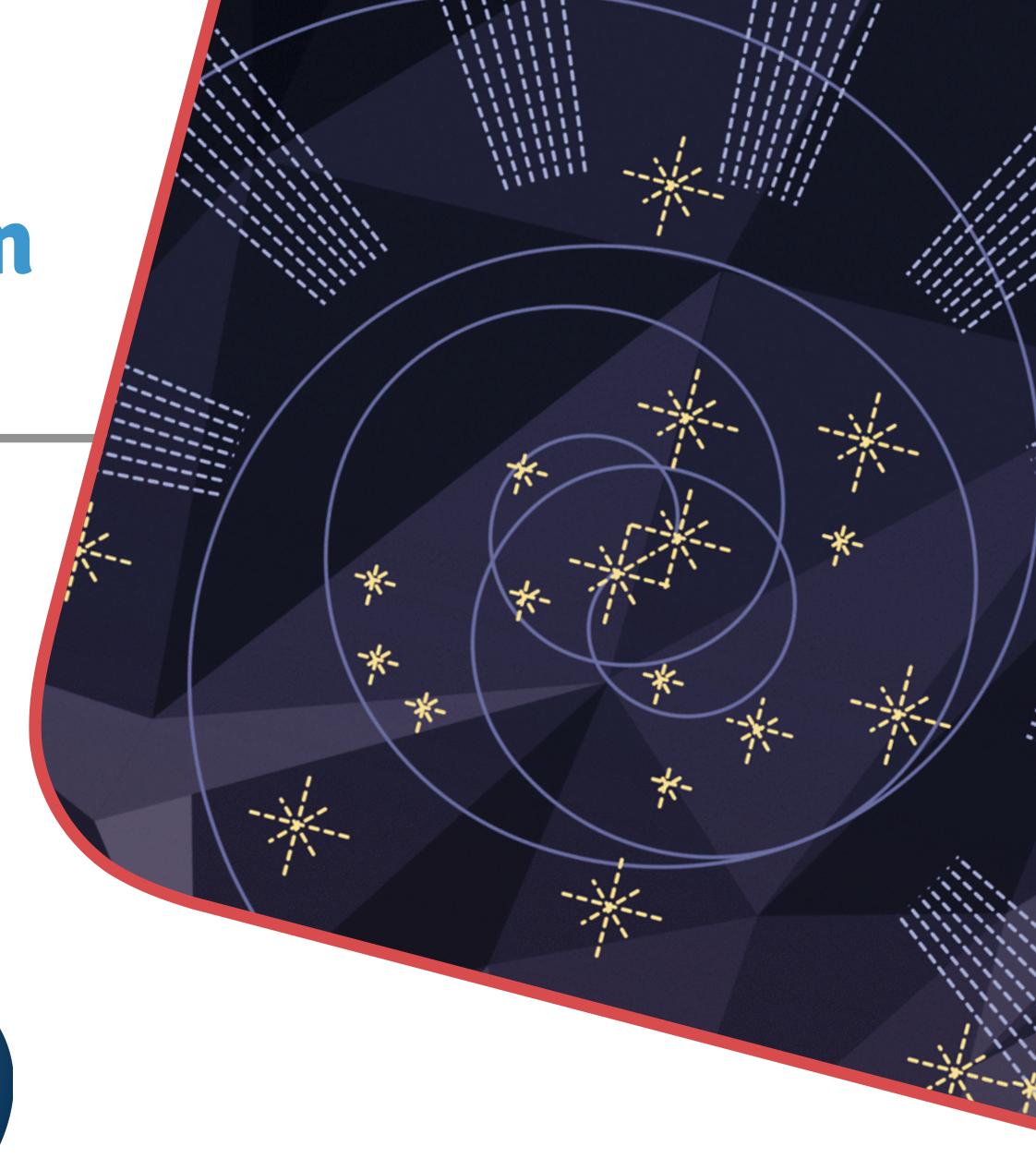
w/ R. Murgia, V. Poulin, PRD 104 (2021)

w/R. Murgia, V. Poulin, J. Lavalle, PRD 105 (2022)









$$S_8 = \sigma_8 \sqrt{\Omega_m/0.3}$$

$$\sigma_8^2 = \int P_m(k, z = 0) W_R^2(k) d\ln k$$

Ω<sub>m</sub> should be left unchanged (well constrained by SNIa & galaxy clustering)

Suppress power at scales 
$$k \sim 0.1 - 1 \ h/{\rm Mpc}$$

$$S_8 = \sigma_8 \sqrt{\Omega_m/0.3}$$

$$\sigma_8^2 = \int P_m(k, z = 0) W_R^2(k) d\ln k$$

$$S_8 = \sigma_8 \sqrt{\Omega_m/0.3}$$

$$\sigma_8^2 = \int P_m(k, z = 0) W_R^2(k) d\ln k$$

Ω<sub>m</sub> should be left unchanged (well constrained by SNIa & galaxy clustering)

Suppress power at scales  $k \sim 0.1 - 1 \ h/{\rm Mpc}$ 

Modify only perturbation properties (expansion history well constrained by low-z probes)

Ω<sub>m</sub> should be left unchanged (well constrained by SNIa & galaxy clustering)

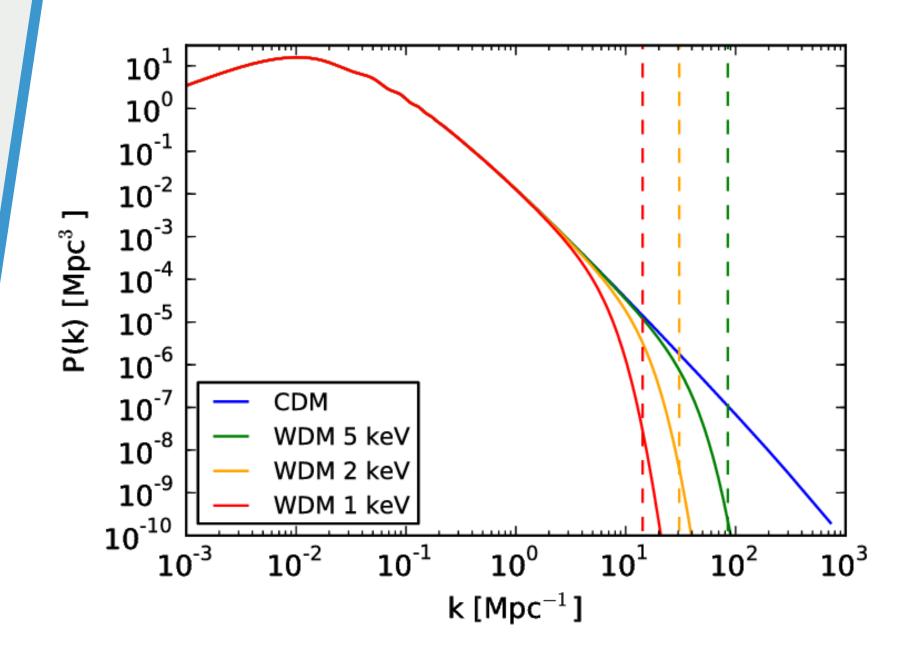
Suppress power at scales  $k \sim 0.1 - 1 \ h/{\rm Mpc}$ 

Modify only perturbation properties (expansion history well constrained by low-z probes)

$$S_8 = \sigma_8 \sqrt{\Omega_m/0.3}$$

$$\sigma_8^2 = \int P_m(k, z = 0) W_R^2(k) d\ln k$$

#### **Ex:** Warm dark matter



Very constrained by Ly-α! [Iršič+ 17]

#### Decaying Dark Matter (DDM)

Well motivated theoretically (ex: R-parity violation)

#### Decaying Dark Matter (DDM)

- Well motivated theoretically (ex: R-parity violation)
- **Decay products?** 
  - To SM particles  $\mbox{Model-dependent, strongly constrained } \Gamma^{-1} \gtrsim 10^7 10^{10} \ t_U$  [Blanco+ 18]
  - To dark radiation  $\mbox{Model-independent, less constrained} \ \Gamma^{-1} \gtrsim 10 \ t_U$  [Nygaard+ 20]

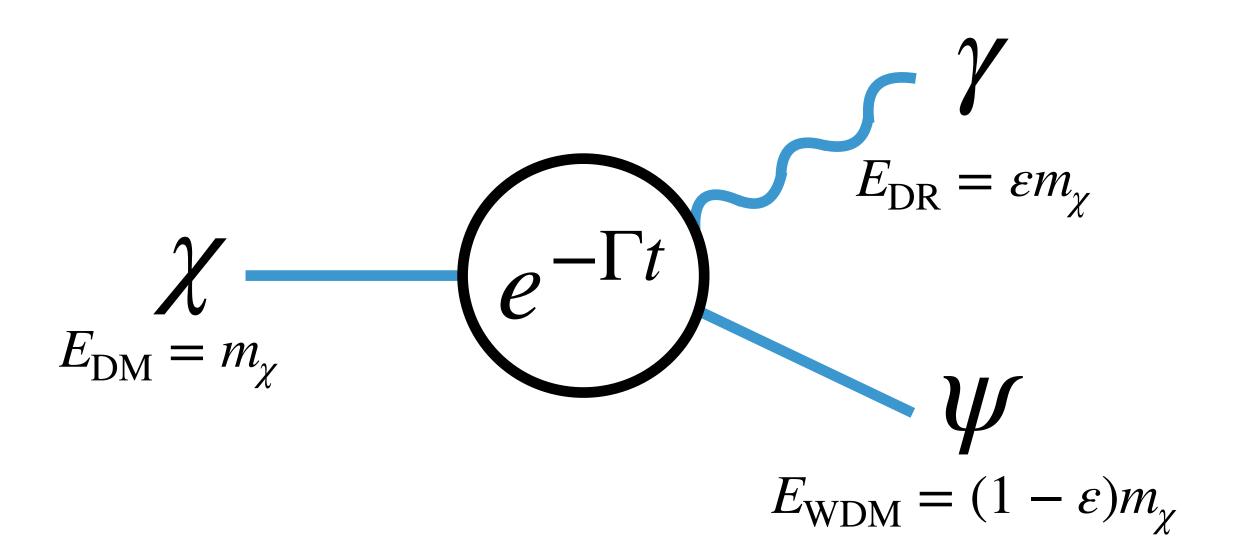
#### Decaying Dark Matter (DDM)

- Well motivated theoretically (ex: R-parity violation)
- Decay products?
  - To SM particles  $\mbox{Model-dependent, strongly constrained } \Gamma^{-1} \gtrsim 10^7 10^{10} \ t_U$  [Blanco+ 18]
  - To dark radiation  $\mbox{Model-independent, less constrained} \ \, \Gamma^{-1} \gtrsim 10 \,\, t_U$  [Nygaard+ 20]

What about massive products?

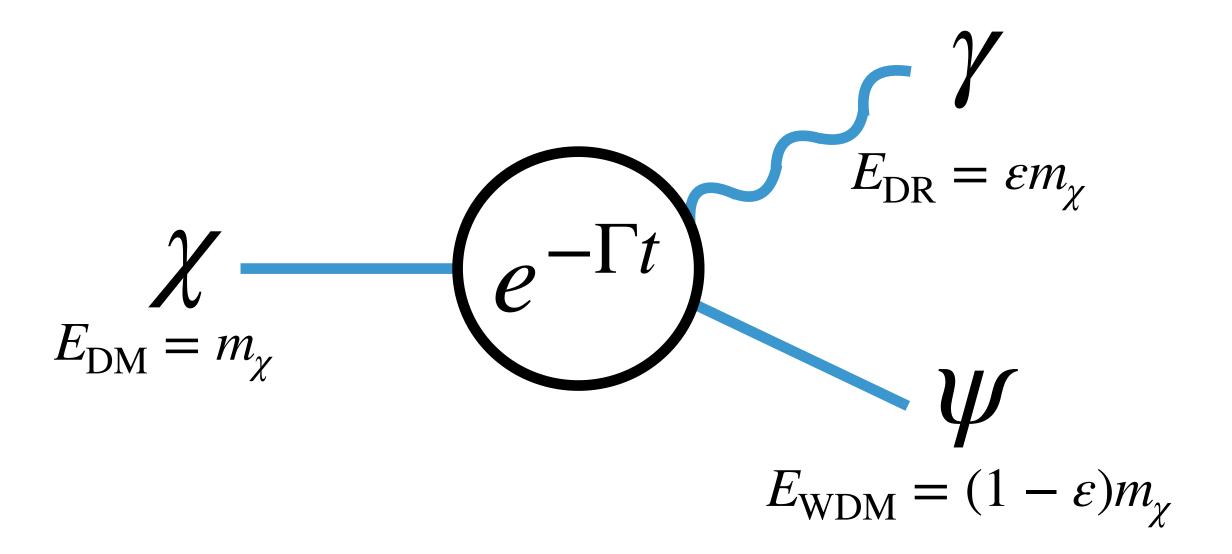
#### DDM with massive decay products

We explore DM decays to massless (Dark Radiation) and massive (Warm Dark Matter) particles



#### DDM with massive decay products

We explore DM decays to massless (Dark Radiation) and massive (Warm Dark Matter) particles



#### 2 extra parameters:

Decay rate  $\Gamma$ DR energy fraction  $\mathcal E$ 

$$\varepsilon = \frac{1}{2} \left( 1 - \frac{m_{\psi}^2}{m_{\chi}^2} \right) \begin{cases} = 0 & (\Lambda \text{CDM}) \\ = 1/2 & (DM \to DR) \end{cases}$$

#### GOAL

Perform a parameter scan by including full treatment of linear perts., in order to assess the impact on the S<sub>8</sub> tension

#### **Evolution of DDM perturbations**

Track  $\delta_i$ ,  $\theta_i$  and  $\sigma_i$  for i = dm, dr, idm

Boltzmann hierarchy of eqs., dictate evolution of p.s.d. multipoles  $\delta f_{\ell}$  (q, k, T)

#### **Evolution of DDM perturbations**

Track  $\delta_i$ ,  $\theta_i$  and  $\sigma_i$  for i = dm, dr, idm

- Boltzmann hierarchy of eqs., dictate evolution of p.s.d. multipoles  $\delta f_{\ell}$  (q, k,  $\tau$ )
  - For DM and DR, momentum d.o.f. are integrated out
  - For WDM, need to follow full evolution in phase space Computationally prohibitive,  $\mathcal{O}(10^8)$  ODEs to solve!

#### New fluid equations for the WDM species

#### Based on previous approximation for massive neutrinos

[Lesgourgues+ 11]

$$\delta'_{\text{wdm}} = -3aH(c_{\text{syn}}^2 - w)\delta_{\text{wdm}} - (1+w)\left(\theta_{\text{wdm}} + \frac{h'}{2}\right) + a\Gamma(1-\varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}(\delta_{\text{dm}} - \delta_{\text{wdm}})$$

$$\theta'_{\text{wdm}} = -aH(1 - 3c_a^2)\theta_{\text{wdm}} + \frac{c_{\text{syn}}^2}{1 + w}k^2\delta_{\text{wdm}} - k^2\sigma_{\text{wdm}} - a\Gamma(1 - \varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}\frac{1 + c_a^2}{1 + w}\theta_{\text{wdm}}$$

#### New fluid equations for the WDM species

#### Based on previous approximation for massive neutrinos

[Lesgourgues+ 11]

$$\delta'_{\text{wdm}} = -3aH(c_{\text{syn}}^2 - w)\delta_{\text{wdm}} - (1+w)\left(\theta_{\text{wdm}} + \frac{h'}{2}\right) + a\Gamma(1-\varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}(\delta_{\text{dm}} - \delta_{\text{wdm}})$$

$$\theta_{\text{wdm}}' = -aH(1 - 3c_a^2)\theta_{\text{wdm}} + \frac{c_{\text{syn}}^2}{1 + w}k^2\delta_{\text{wdm}} - k^2\sigma_{\text{wdm}} - a\Gamma(1 - \varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}\frac{1 + c_a^2}{1 + w}\theta_{\text{wdm}}$$

#### New fluid equations for the WDM species

#### Based on previous approximation for massive neutrinos

[Lesgourgues+ 11]

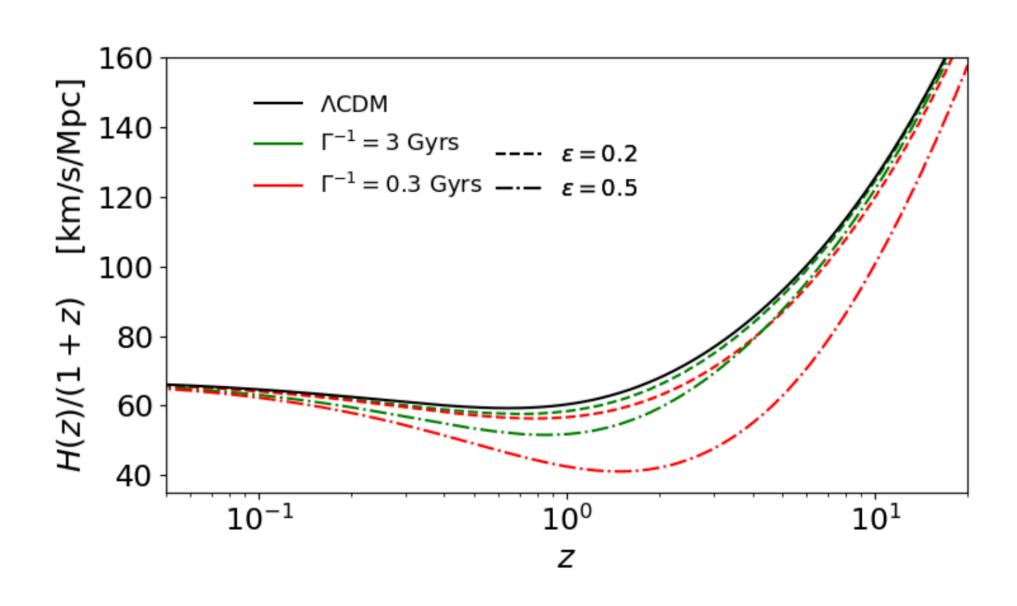
$$\delta'_{\text{wdm}} = -3aH(c_{\text{syn}}^2 - w)\delta_{\text{wdm}} - (1+w)\left(\theta_{\text{wdm}} + \frac{h'}{2}\right) + a\Gamma(1-\varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}(\delta_{\text{dm}} - \delta_{\text{wdm}})$$

$$\theta'_{\text{wdm}} = -aH(1 - 3c_a^2)\theta_{\text{wdm}} + \frac{c_{\text{syn}}^2}{1 + w}k^2\delta_{\text{wdm}} - k^2\sigma_{\text{wdm}} - a\Gamma(1 - \varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}\frac{1 + c_a^2}{1 + w}\theta_{\text{wdm}}$$

CPU time reduced from
~ 1 day to ~ 1 minute !!

#### H(z) more affected by the DR:

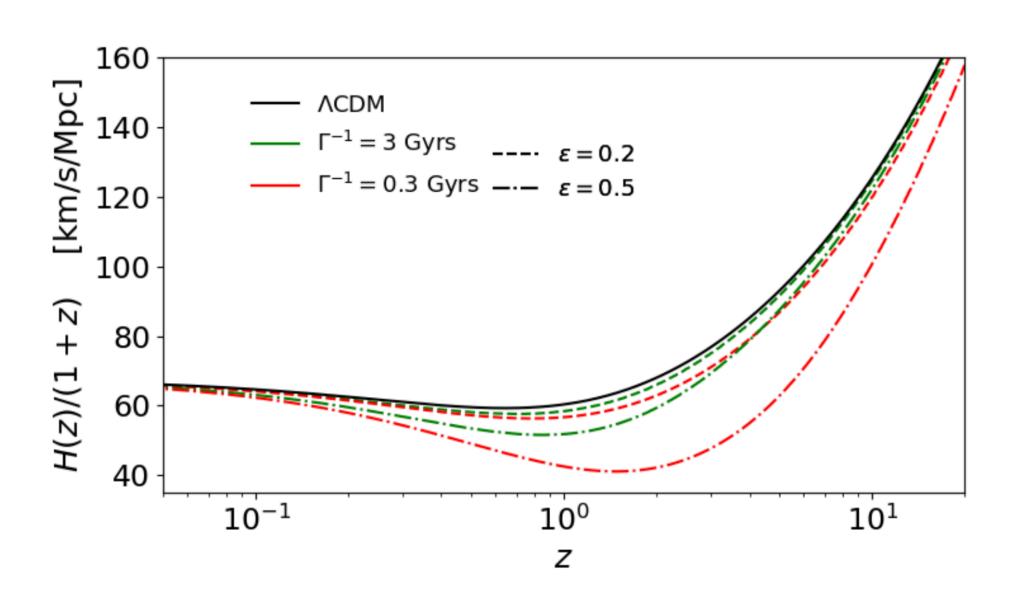
 $\Gamma^{\uparrow}$   $\mathcal{E}^{\uparrow}$ 



Impact on background

#### H(z) more affected by the DR:

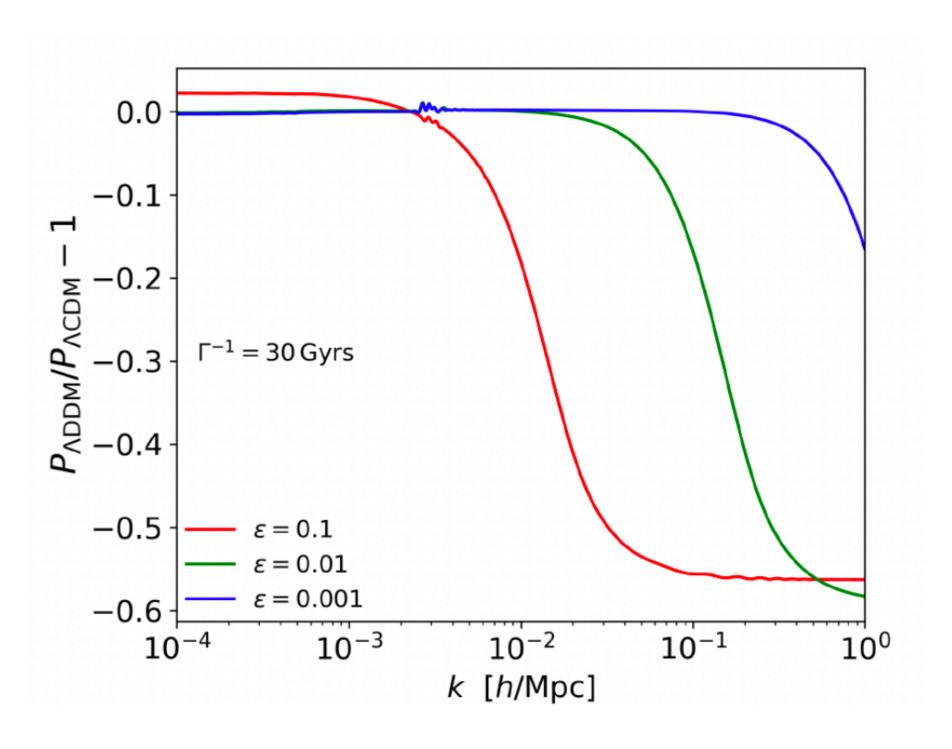
$$\Gamma$$
  $\varepsilon$ 



# Impact on background t on perturbations

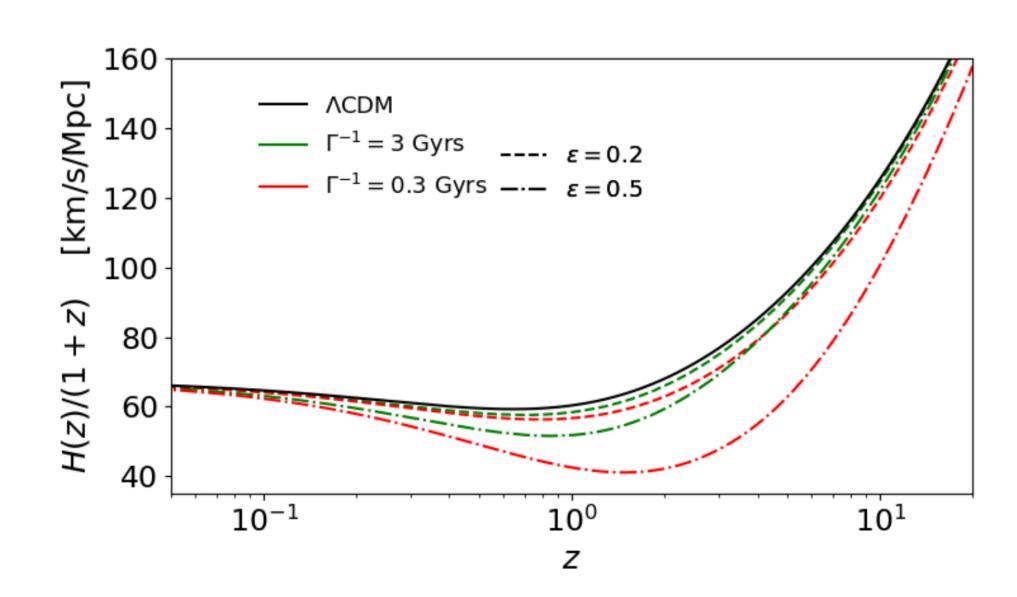
#### P(k) more affected by the WDM (suppression at $k > k_{fs}$ ):

$$\Gamma$$
 $\mathcal{E}$ 



#### H(z) more affected by the DR:

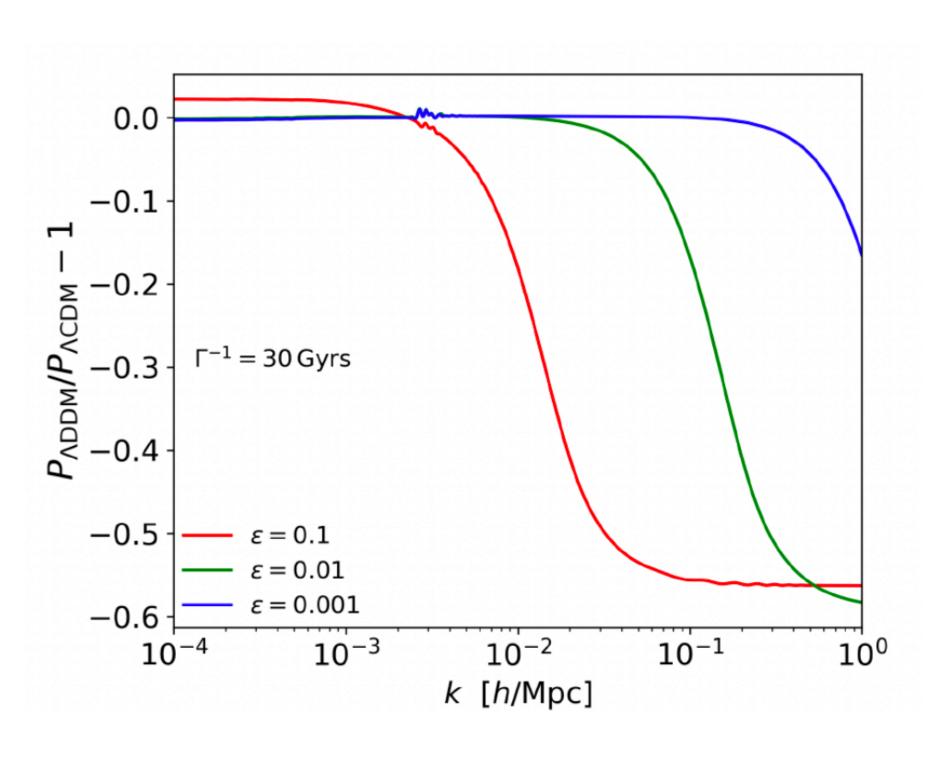
$$\Gamma$$
 $\mathcal{E}$ 



With large  $\Gamma$  and small  $\epsilon$ , we can achieve a P(k) suppression while leaving H(z) unaffected

P(k) more affected by the WDM (suppression at  $k > k_{fs}$ ):

$$\Gamma$$
 $\mathcal{E}$ 



Impact on background

t on perturbations

Impaci

To compare against weak-lensing data, we need the non-linear prediction

To compare against weak-lensing data, we need the non-linear prediction

## BUT this would require to run many expensive ADDM simulations

To compare against weak-lensing data, we need the non-linear prediction

#### BUT

this would require to run many expensive ADDM simulations

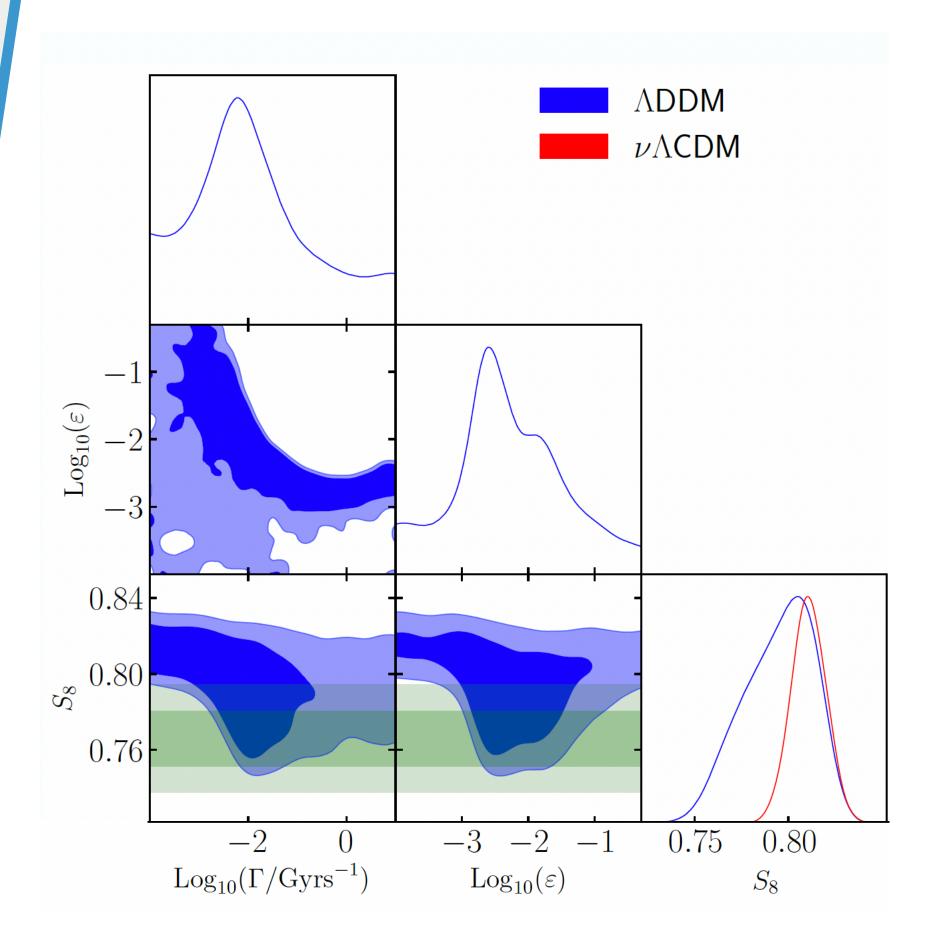
Use a S<sub>8</sub> prior instead (very simplistic, but should be seen as a minimal test)



#### Explaining the S<sub>8</sub> tension

Reconstructed S<sub>8</sub> values are in excellent agreement with WL data

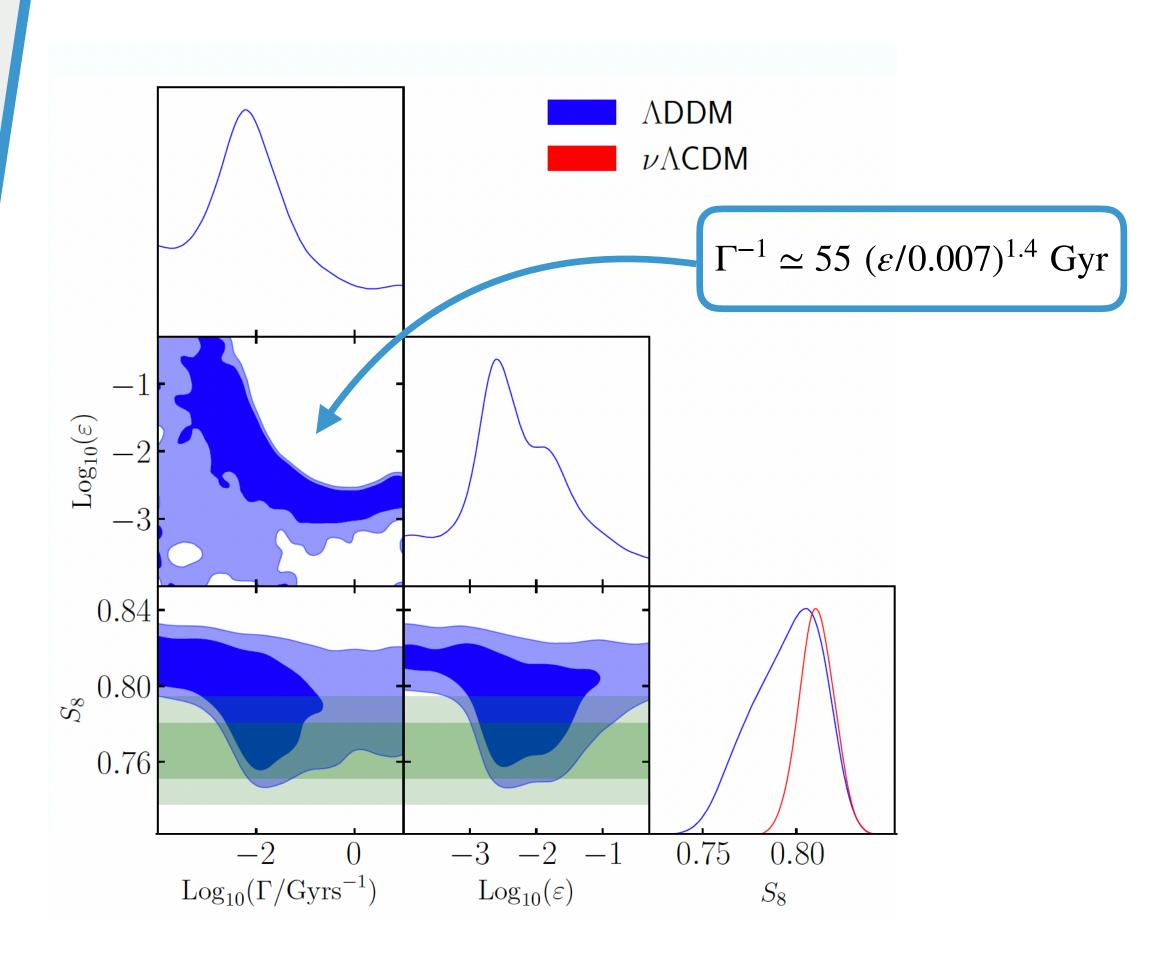
Planck18 + BAO + SNIa + S<sub>8</sub> (KiDS+BOSS+2dfLenS):



#### Explaining the S<sub>8</sub> tension

Reconstructed S<sub>8</sub> values are in excellent agreement with WL data

Planck18 + BAO + SNIa + S<sub>8</sub> (KiDS+BOSS+2dfLenS):



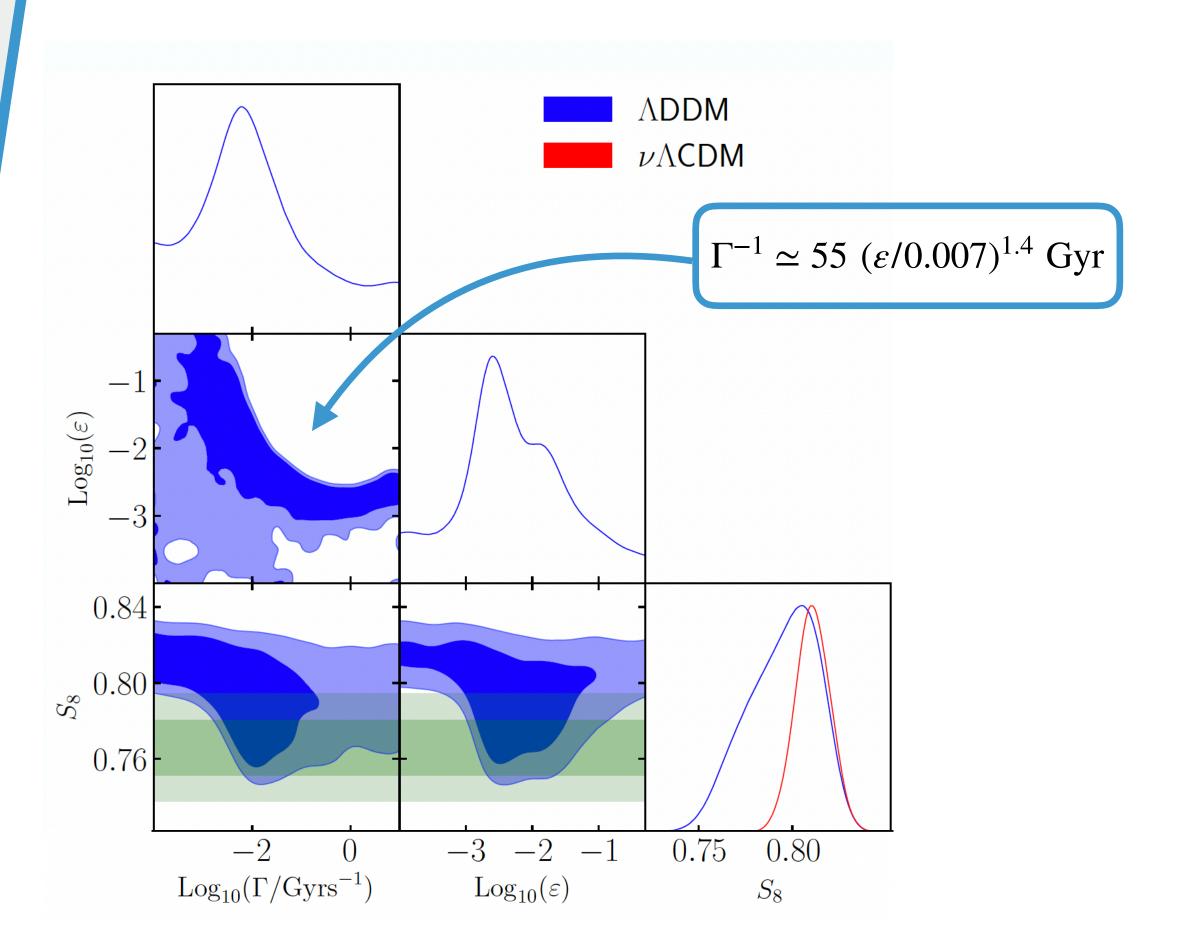
#### Explaining the S<sub>8</sub> tension

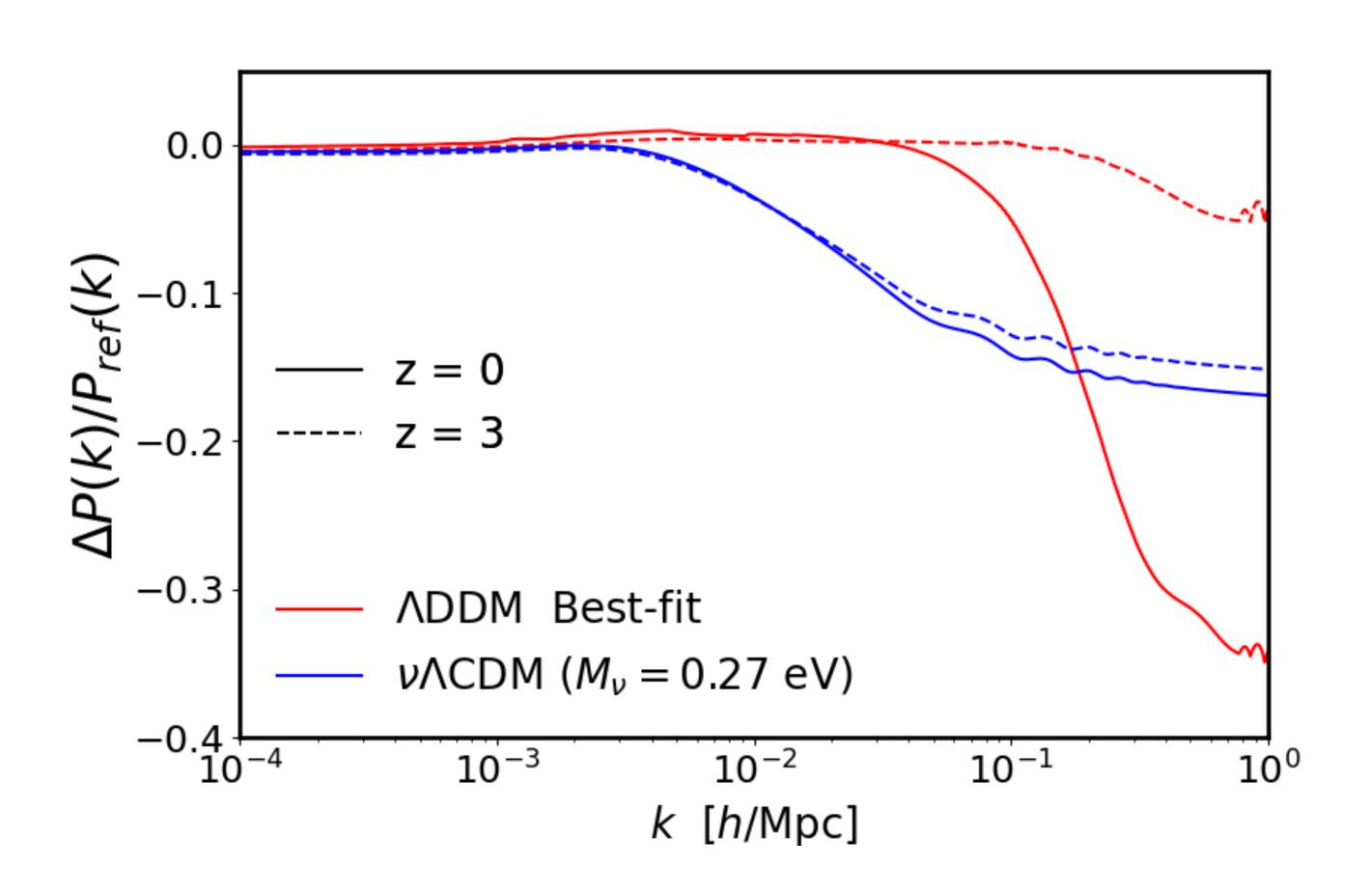
Reconstructed S<sub>8</sub> values are in excellent agreement with WL data

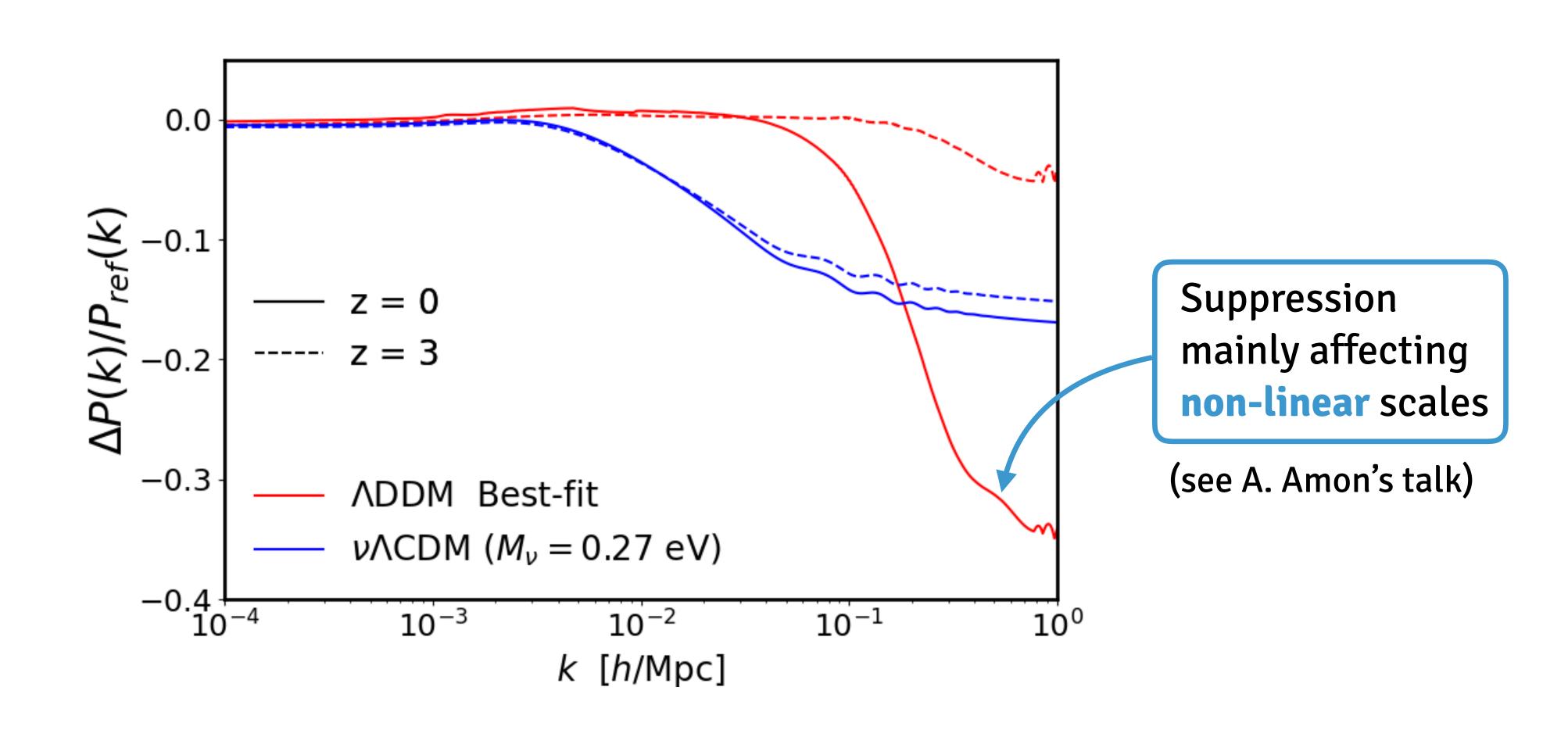
	νΛCDM	<b>ADDM</b>
$\chi^2_{\rm CMB}$	1015.9	1015.2
$\chi^2_{S_8}$	5.64	0.002

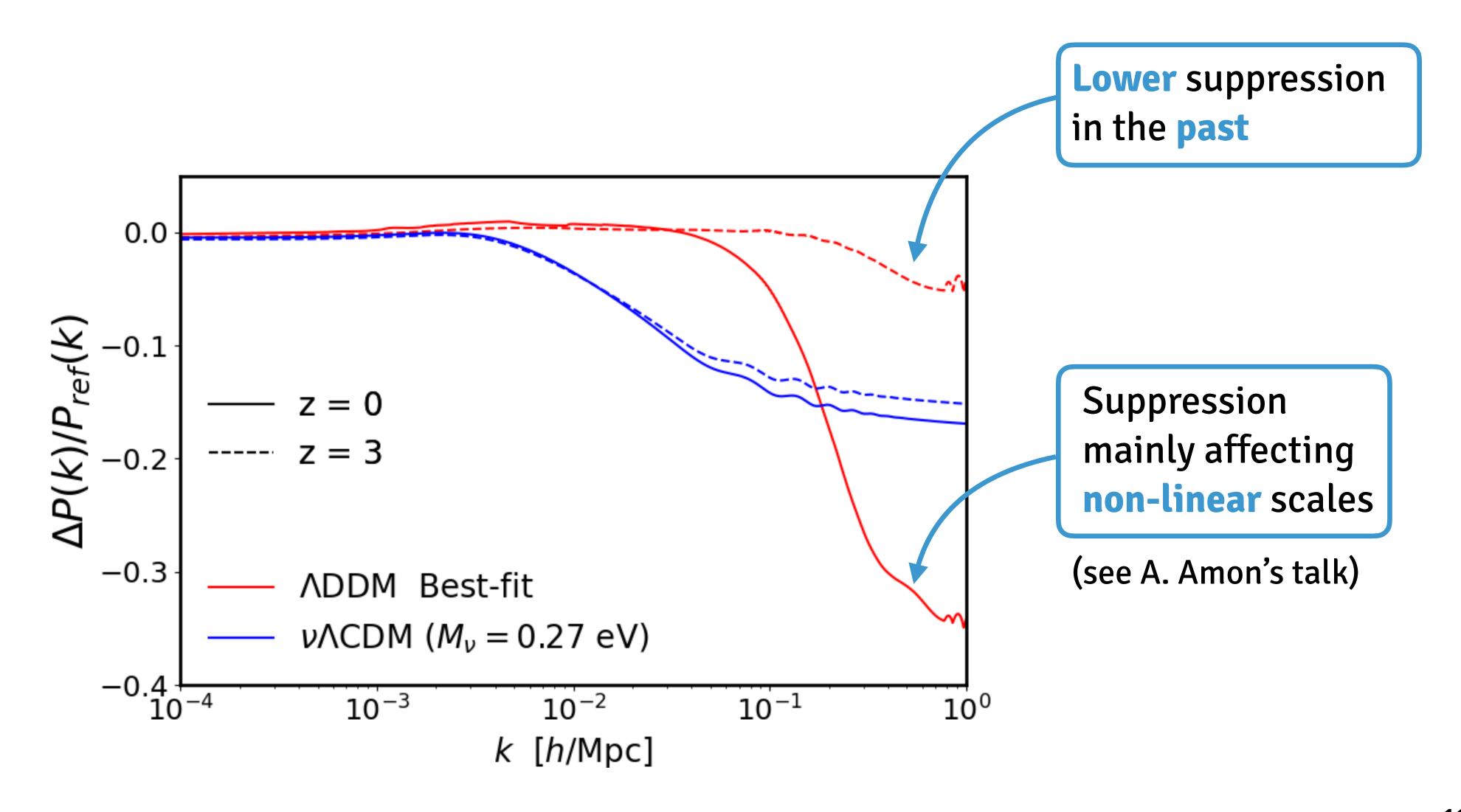
$$\Delta \chi_{\min}^2 = -5.5$$

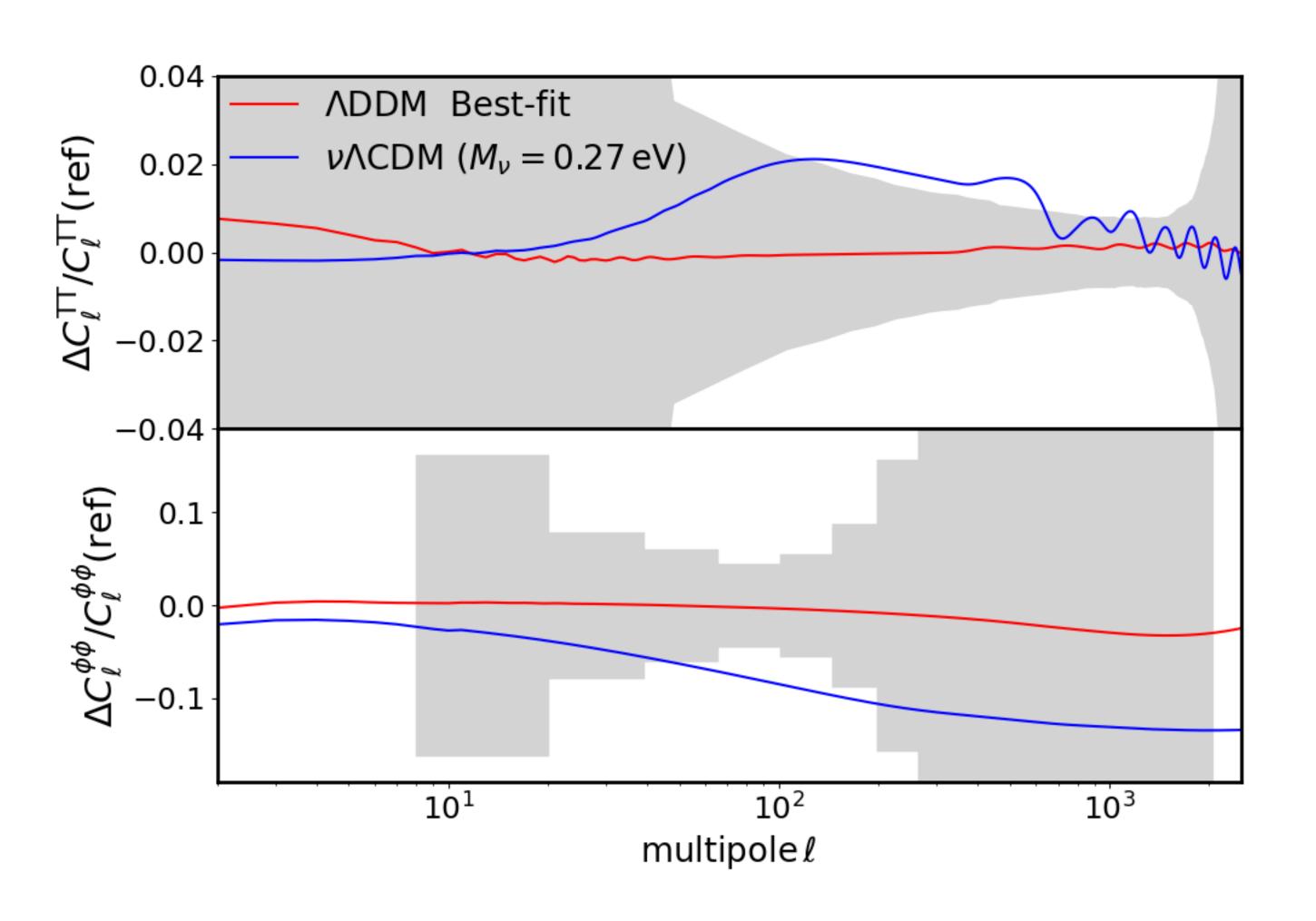
Planck18 + BAO + SNIa + S<sub>8</sub> (KiDS+BOSS+2dfLenS):







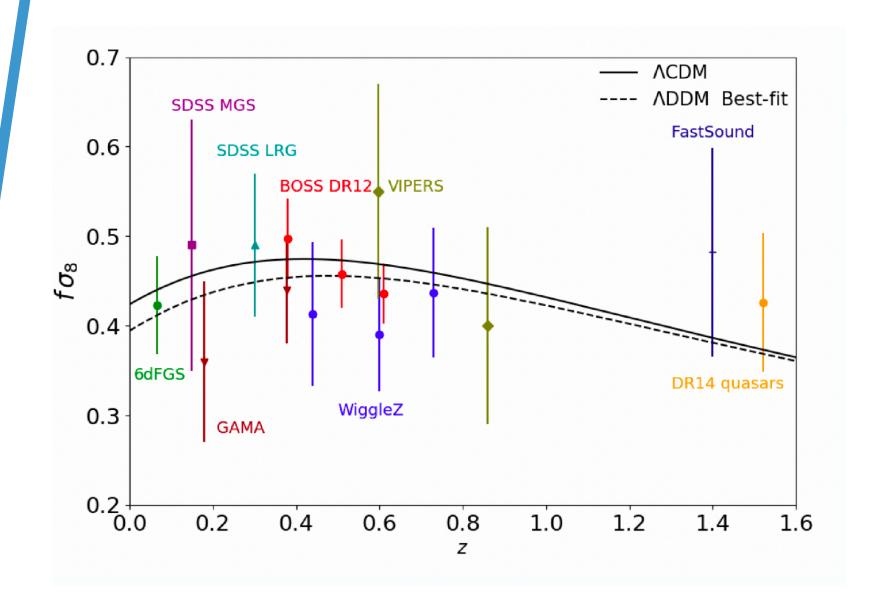


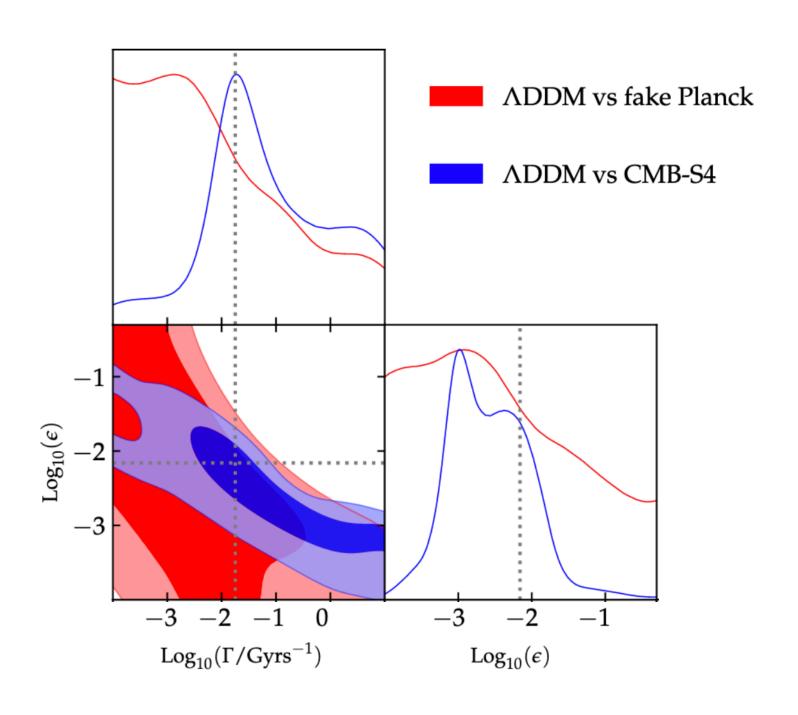


Time-dependence of DDM suppression allows for a better fit to CMB data

#### **Prospects for DDM**

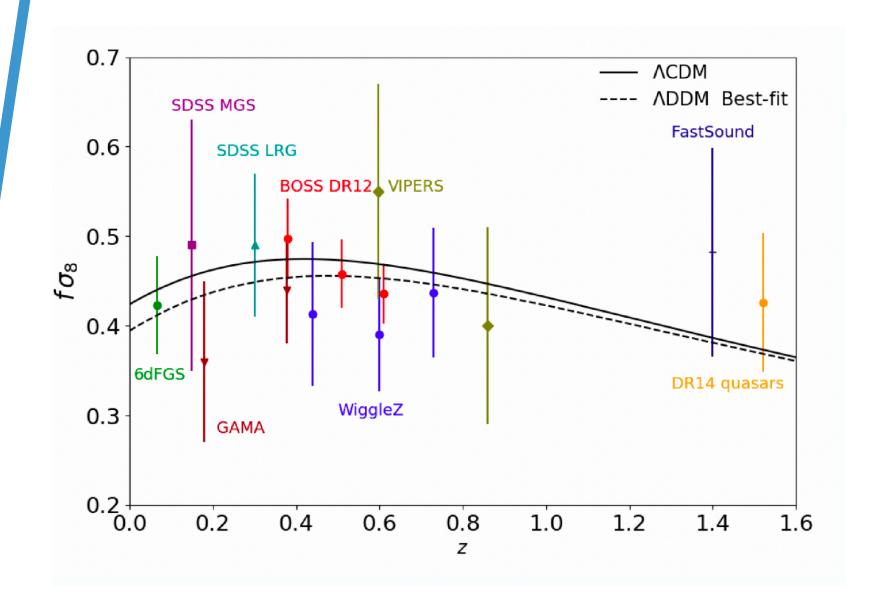
Future accurate fσ<sub>8</sub> and CMB data will be able to capture DDM signature

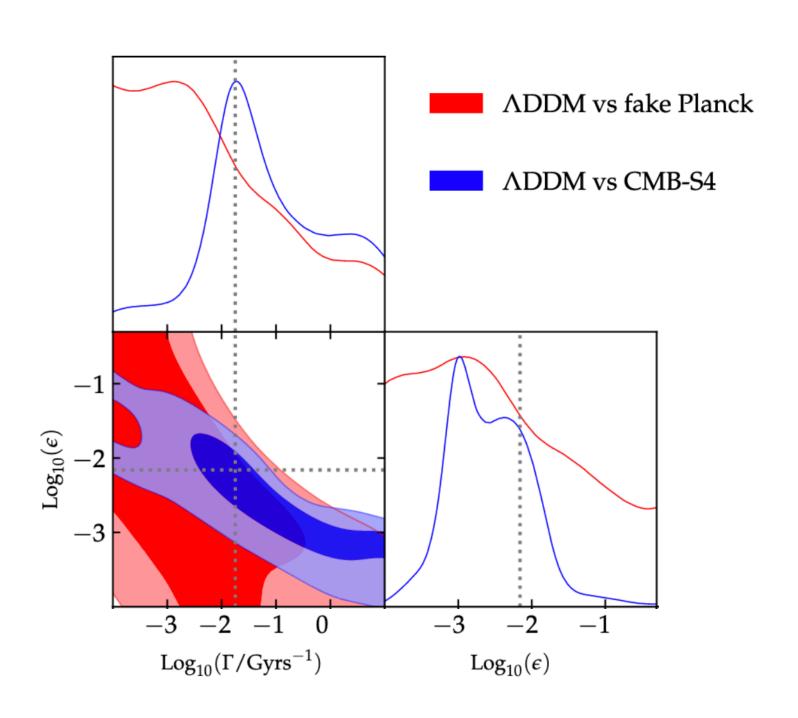




#### **Prospects for DDM**

- Future accurate  $f\sigma_8$  and CMB data will be able to capture DDM signature
- Run DDM simulations, to test model against non-linear observables like Cosmic Shear or Lyman-α forest

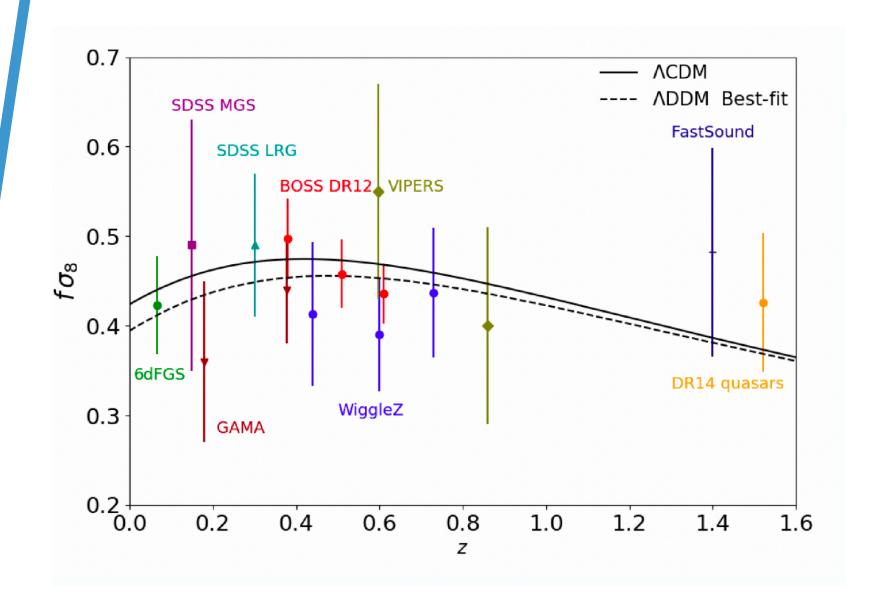


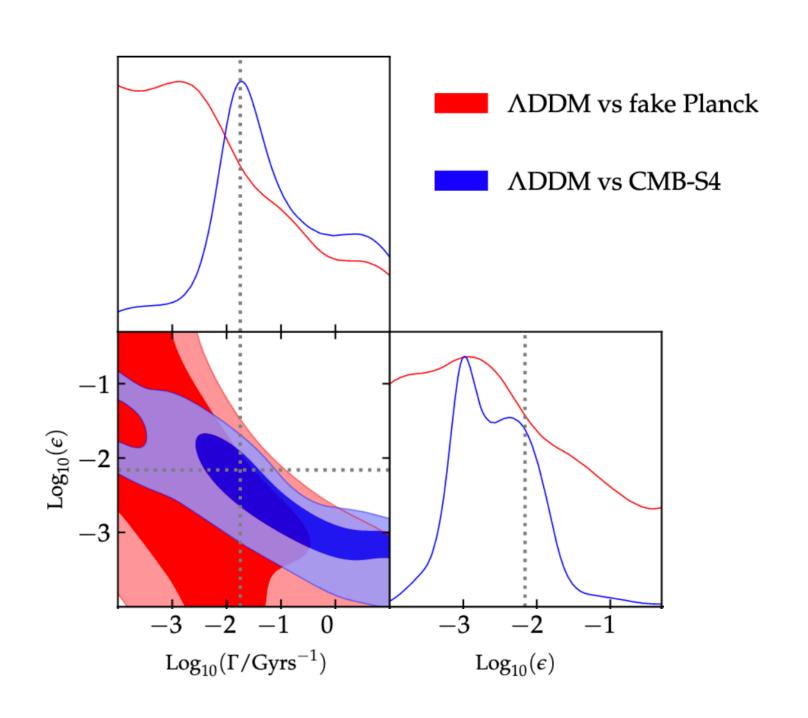


#### **Prospects for DDM**

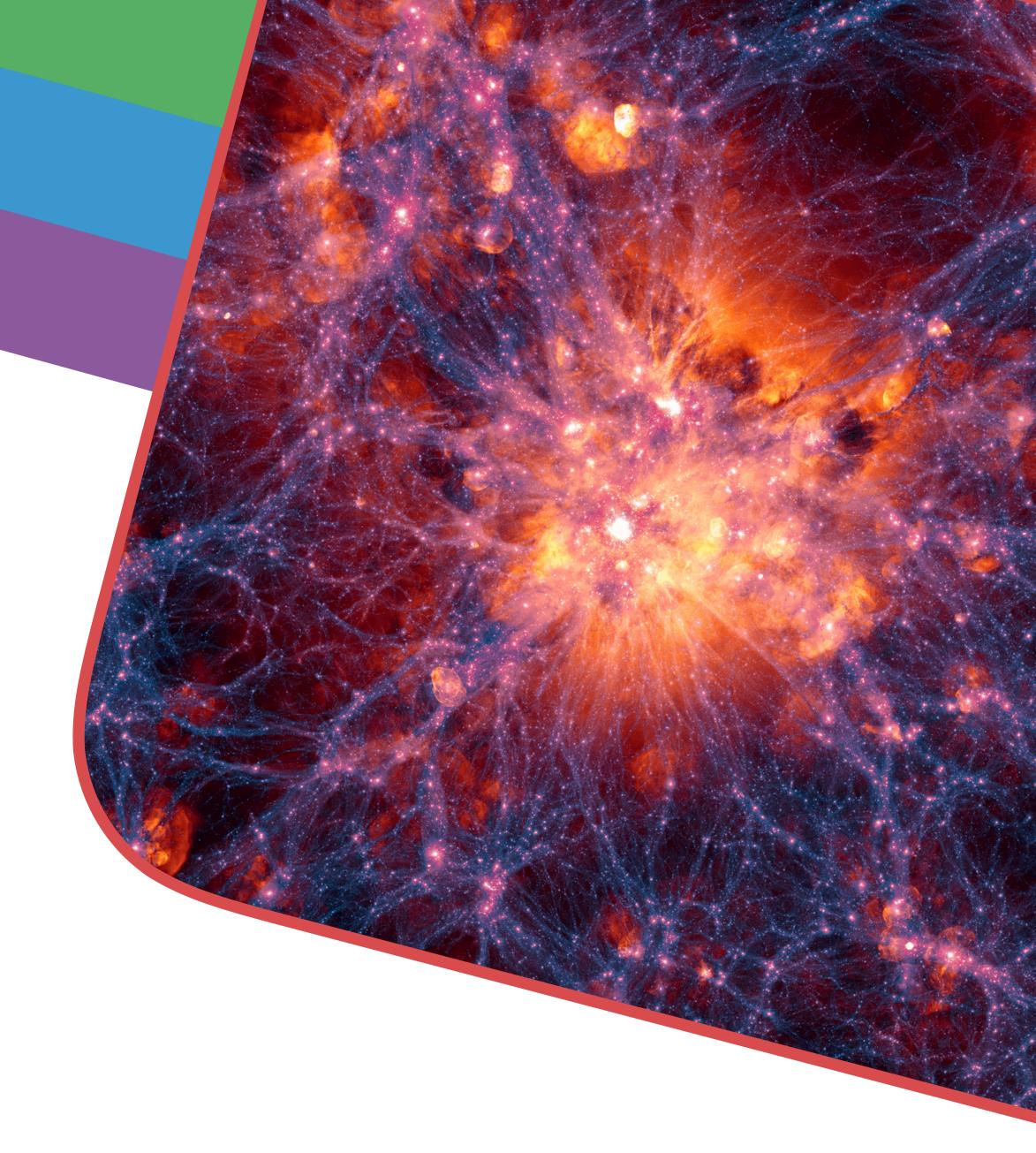
- Future accurate  $f\sigma_8$  and CMB data will be able to capture DDM signature
- Run DDM simulations, to test model against non-linear observables like Cosmic Shear or Lyman-α forest
- Mildly non-linear analysis using the EFTofBOSS data already improves constraints on lifetime

[Simon+ 22]

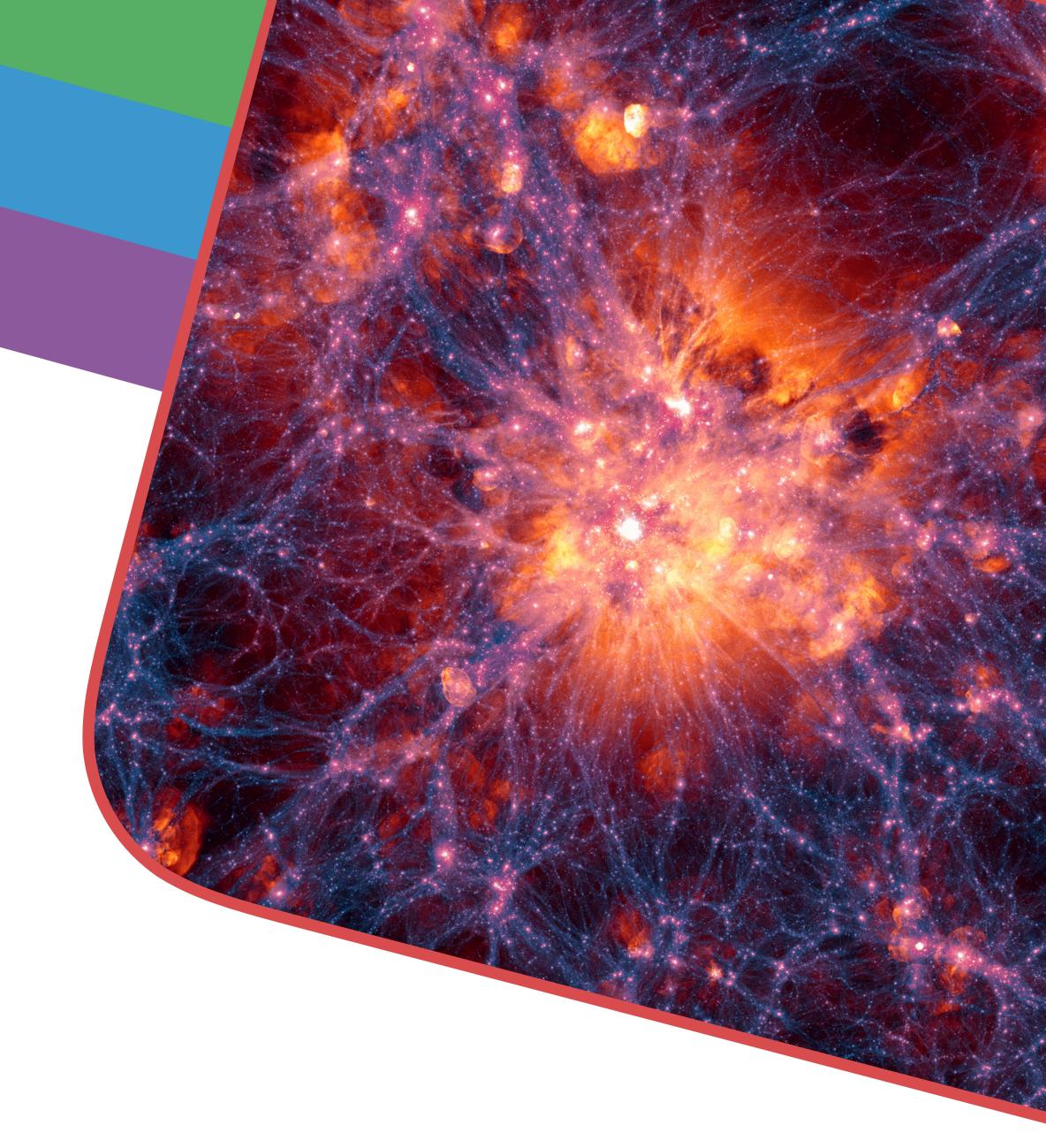




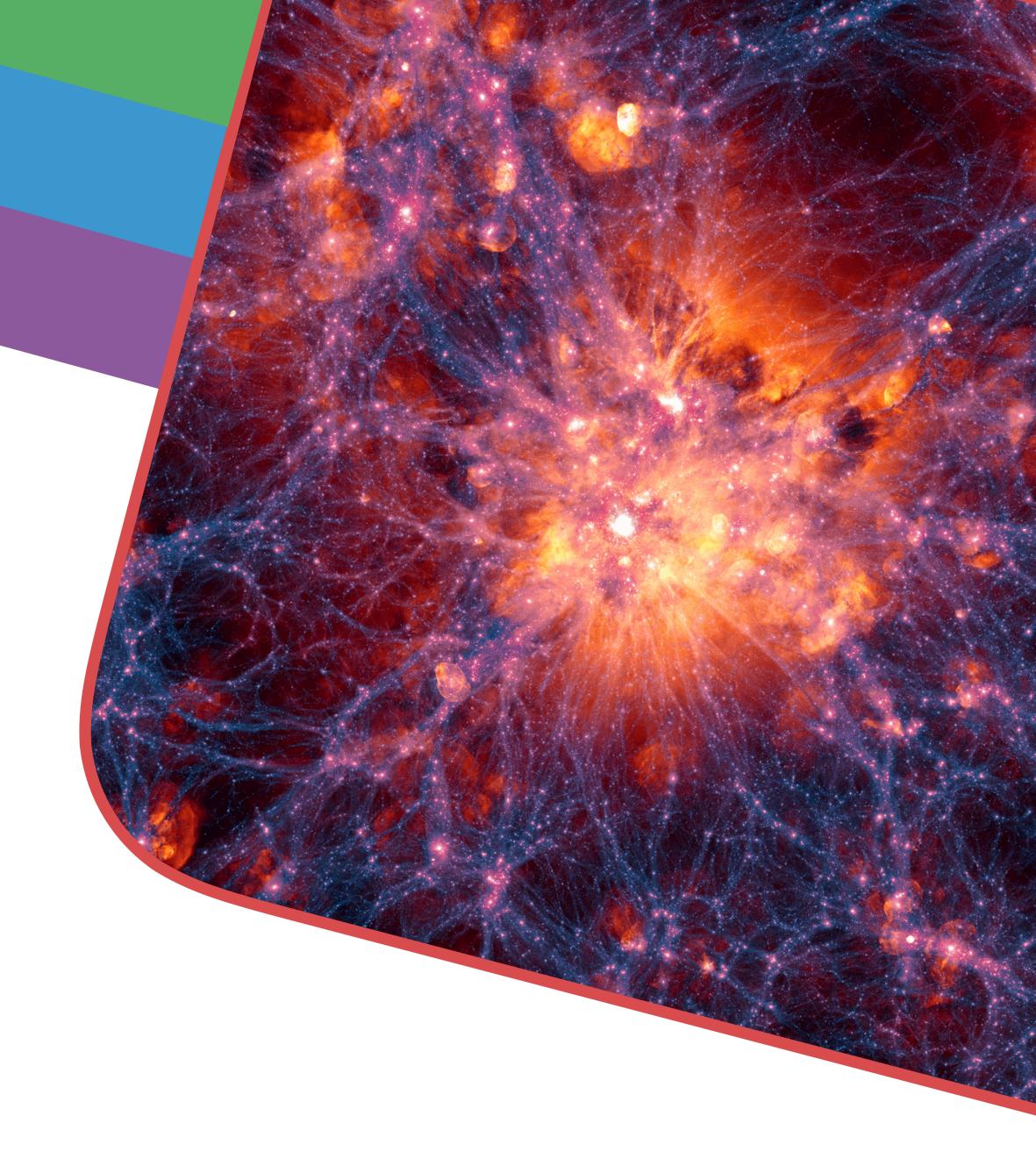
First thorough cosmological analysis of this 2-body DM decay scenario by including a full treatment of perturbations



- First thorough cosmological analysis of this 2-body DM decay scenario by including a full treatment of perturbations
- It can successfully explain the S<sub>8</sub> anomaly while providing a good fit to CMB, BAO and SNIa data

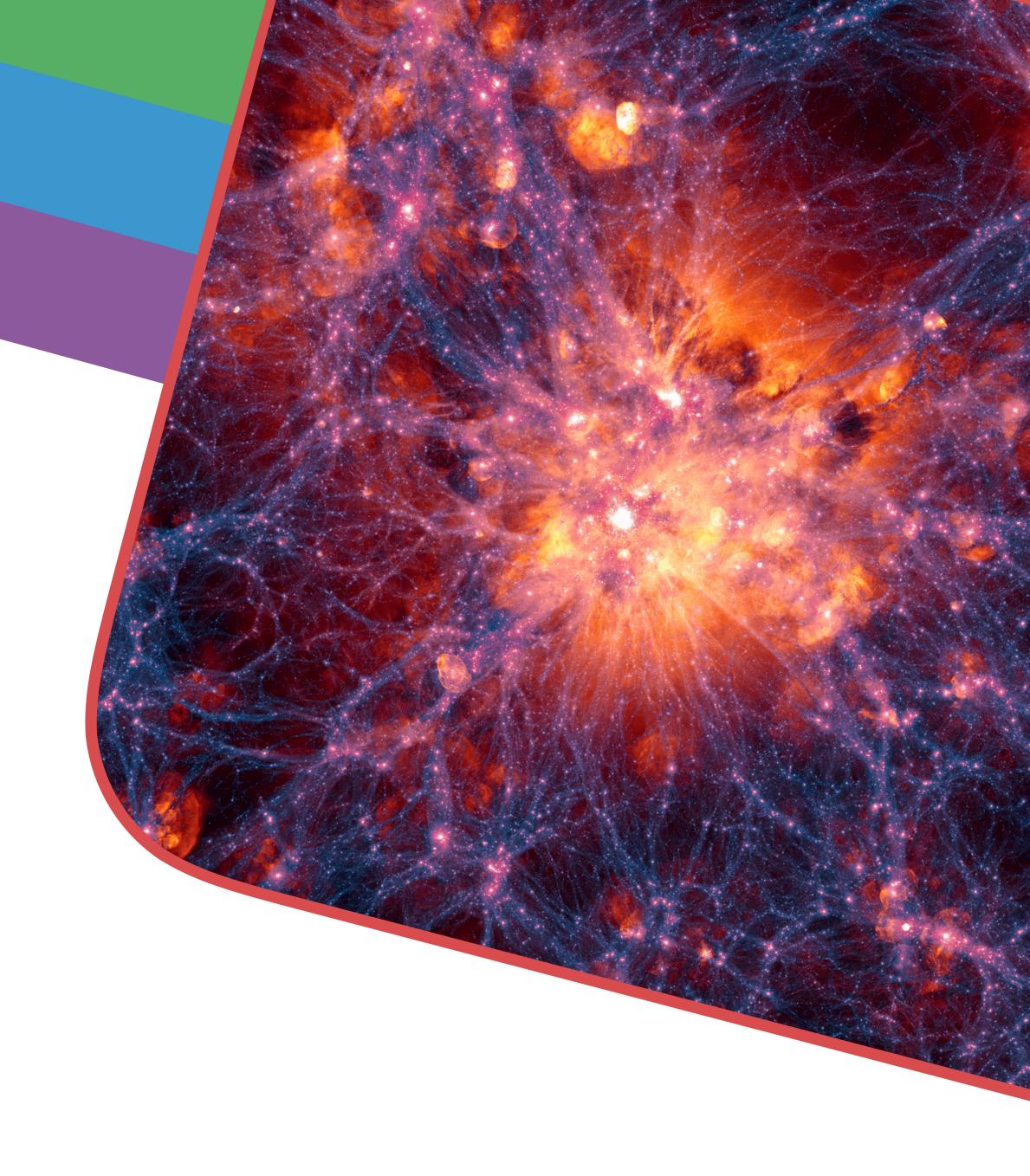


- First thorough cosmological analysis of this 2-body DM decay scenario by including a full treatment of perturbations
- It can successfully explain the S<sub>8</sub> anomaly while providing a good fit to CMB, BAO and SNIa data
- Future accurate growth factor and CMB data will be able to further test this scenario



THANKS FOR YOUR ATTENTION

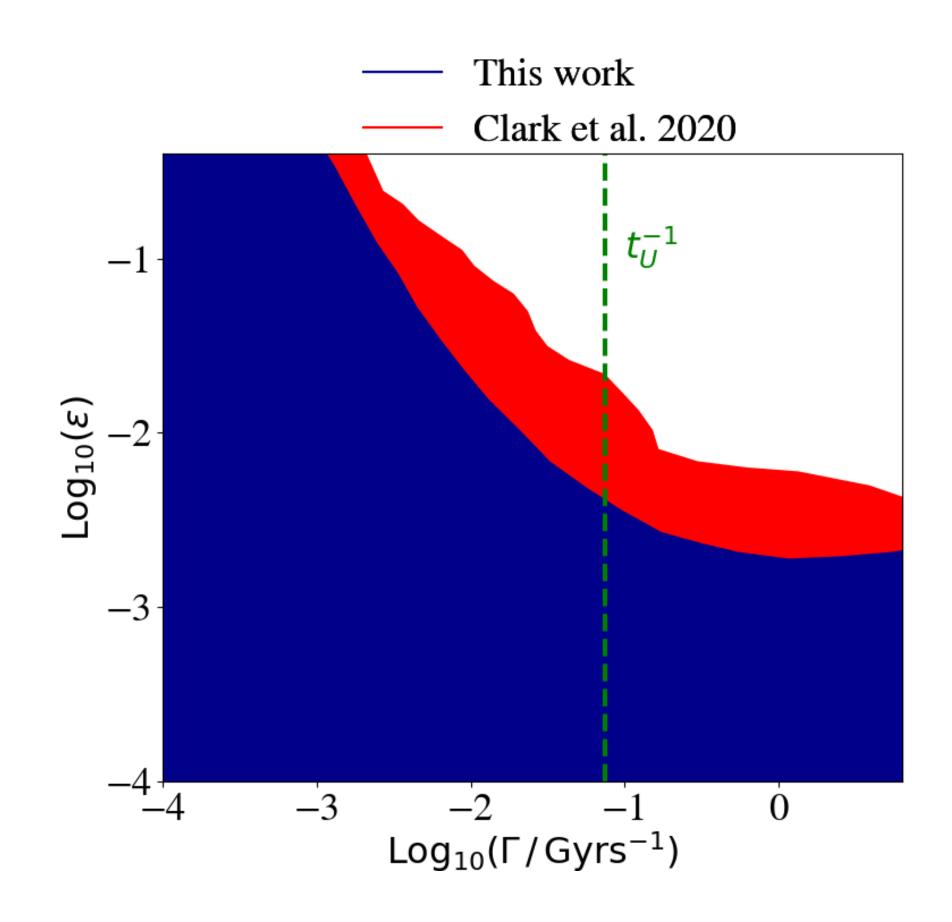
guillermo.franco-abellan@umontpellier.fr



# BACK-UP

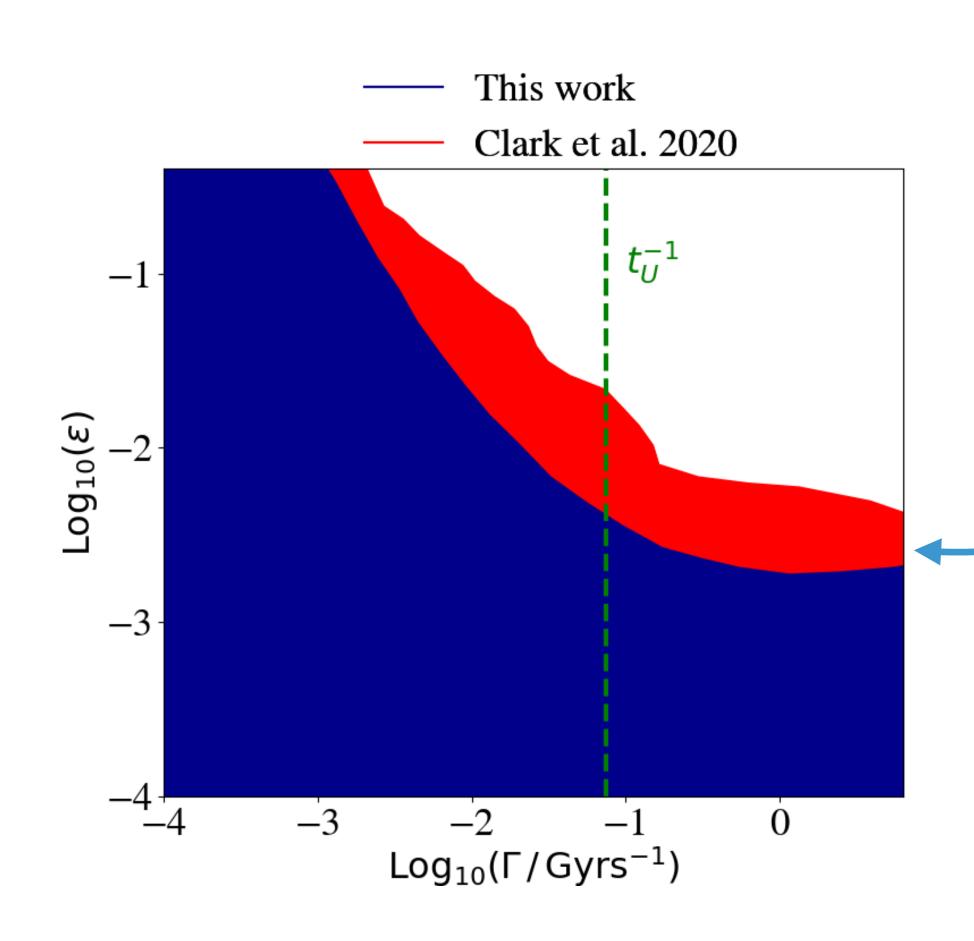
### **General constraints**

### Planck18 + BAO + SNIa:



### **General constraints**

### Planck18 + BAO + SNIa:



Constraints up to 1 order of magnitude stronger than former works due to the inclusion of WDM perts.

# Interesting implications

Model building

Why  $\epsilon << 1/2$ , i.e.  $m_{wdm} \sim m_{dm}$ ? **Ex:** Supergravity

[Choi+ 21]

### Interesting implications

Model building

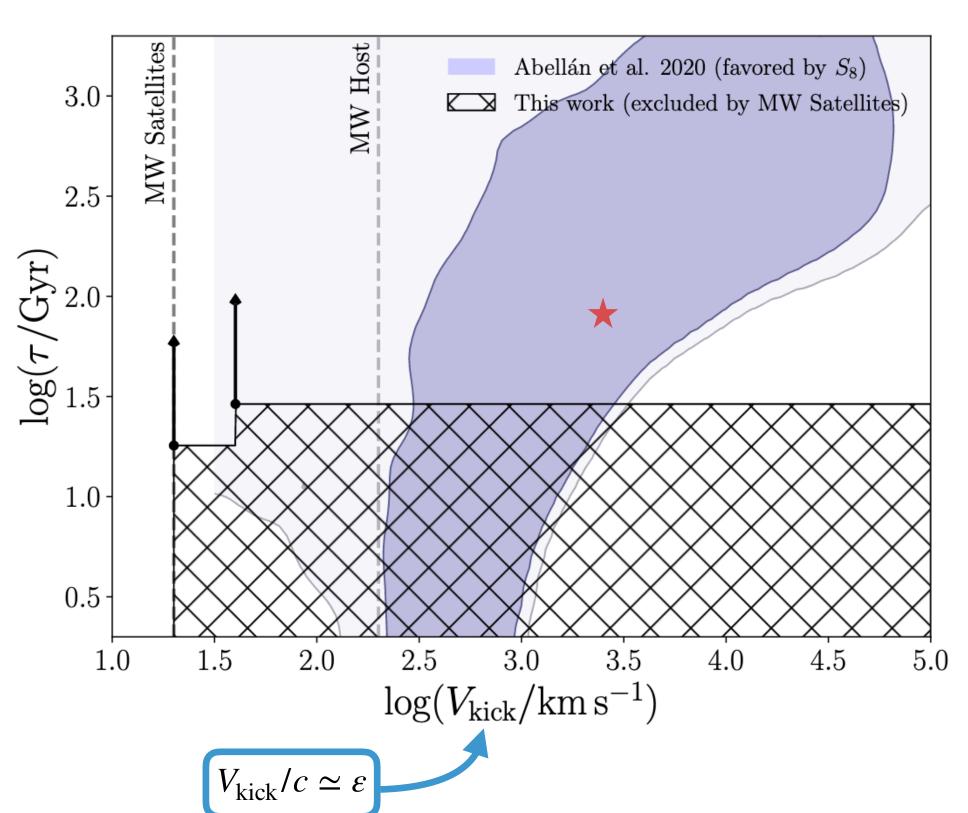
Why  $\epsilon << 1/2$ , i.e.  $m_{wdm} \sim m_{dm}$ ?

Ex: Supergravity

[Choi+ 21]

Reduction in the abundance of subhalos, can be constrained by observations of MW satellites

[DES 22]



# The full Boltzmann hierarchy

$$f(q, k, \mu, \tau) = \overline{f}(q, \tau) + \delta f(q, k, \mu, \tau)$$

Expand  $\delta f$  in multipoles. The Boltzmann eq. leads to the following hierarchy (in synchronous gauge comoving with the mother)

$$\frac{\partial}{\partial \tau} \left( \delta f_0 \right) = -\frac{\mathbf{q}k}{\mathbf{a}\mathbf{E}} \delta f_1 + q \frac{\partial \bar{f}}{\partial q} \frac{\dot{h}}{6} + \frac{\Gamma \bar{N}_{dm}(\tau)}{4\pi q^3 H} \delta(\tau - \tau_q) \delta_{dm},$$

$$\frac{\partial}{\partial \tau} \left( \delta f_1 \right) = \frac{\mathbf{q}k}{3\mathbf{a}\mathbf{E}} \left[ \delta f_0 - 2\delta f_2 \right],$$

$$\frac{\partial}{\partial \tau} \left( \delta f_2 \right) = \frac{\mathbf{q}k}{5\mathbf{a}\mathbf{E}} \left[ 2\delta f_1 - 3\delta f_3 \right] - q \frac{\partial \bar{f}}{\partial q} \frac{(\dot{h} + 6\dot{\eta})}{15},$$

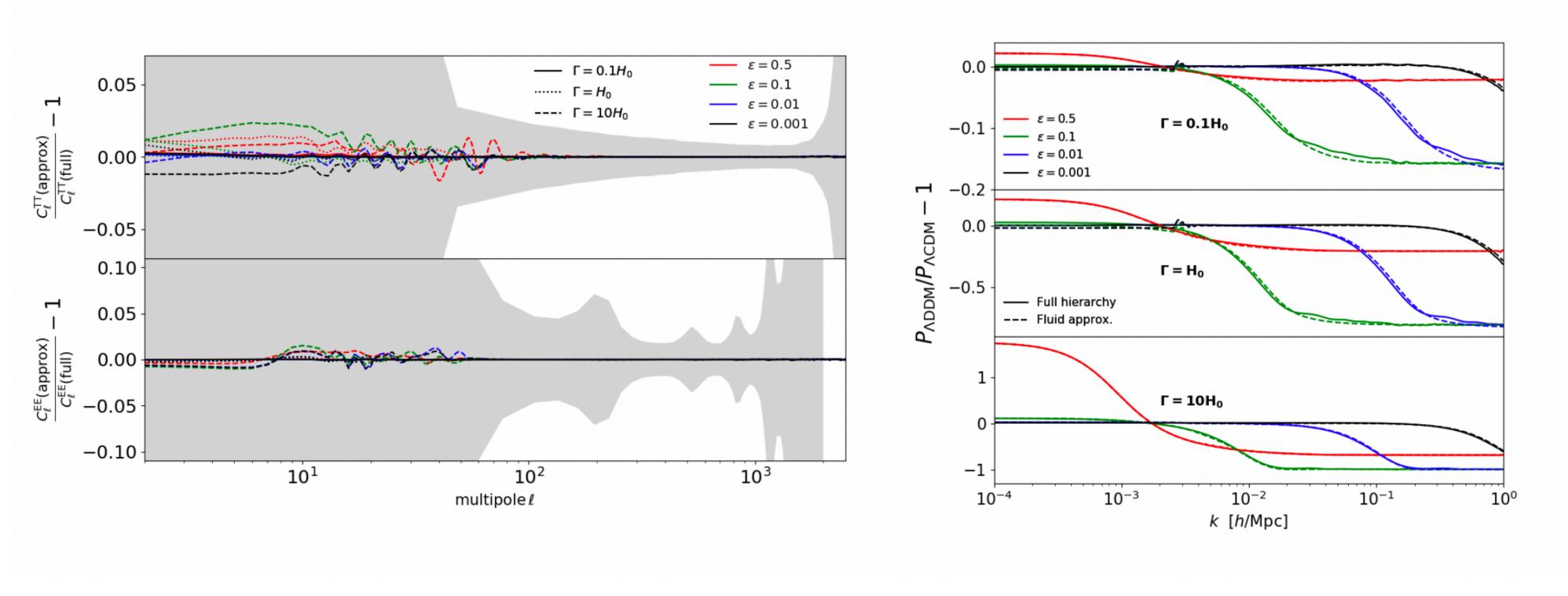
$$\frac{\partial}{\partial \tau} \left( \delta f_\ell \right) = \frac{\mathbf{q}k}{(2\ell + 1)\mathbf{a}\mathbf{E}} \left[ \ell \delta f_{\ell-1} - (\ell + 1)\delta f_{\ell+1} \right] \qquad \text{(for } \ell \geq 3).$$

where  $q = a(\tau_q)p_{\text{max}}$  . In the relat. limit  $\mathbf{q}/\mathbf{aE} = \mathbf{1}$  , so one can take

$$F_{\ell} \equiv \frac{4\pi}{\rho_c} \int dq \ q^3 \delta f_{\ell}$$
 and integrate out the dependency on q

# Checking the accuracy of the WDM fluid approx.

We compare the full Boltzmann hierarchy calculation with the WDM fluid approx.

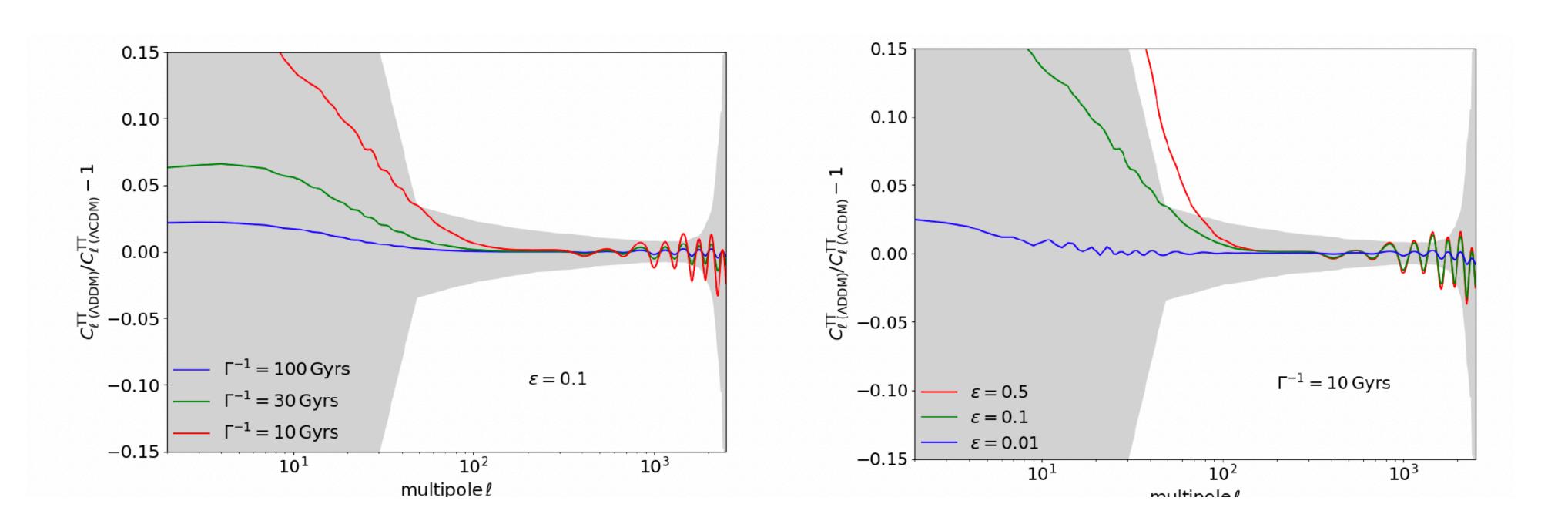


The max. error on  $S_8$  is ~0.65 %, smaller than the ~1.8 % error of the measurement from BOSS+KiDS+2dfLenS

### Impact of DDM on the CMB temperature spectrum

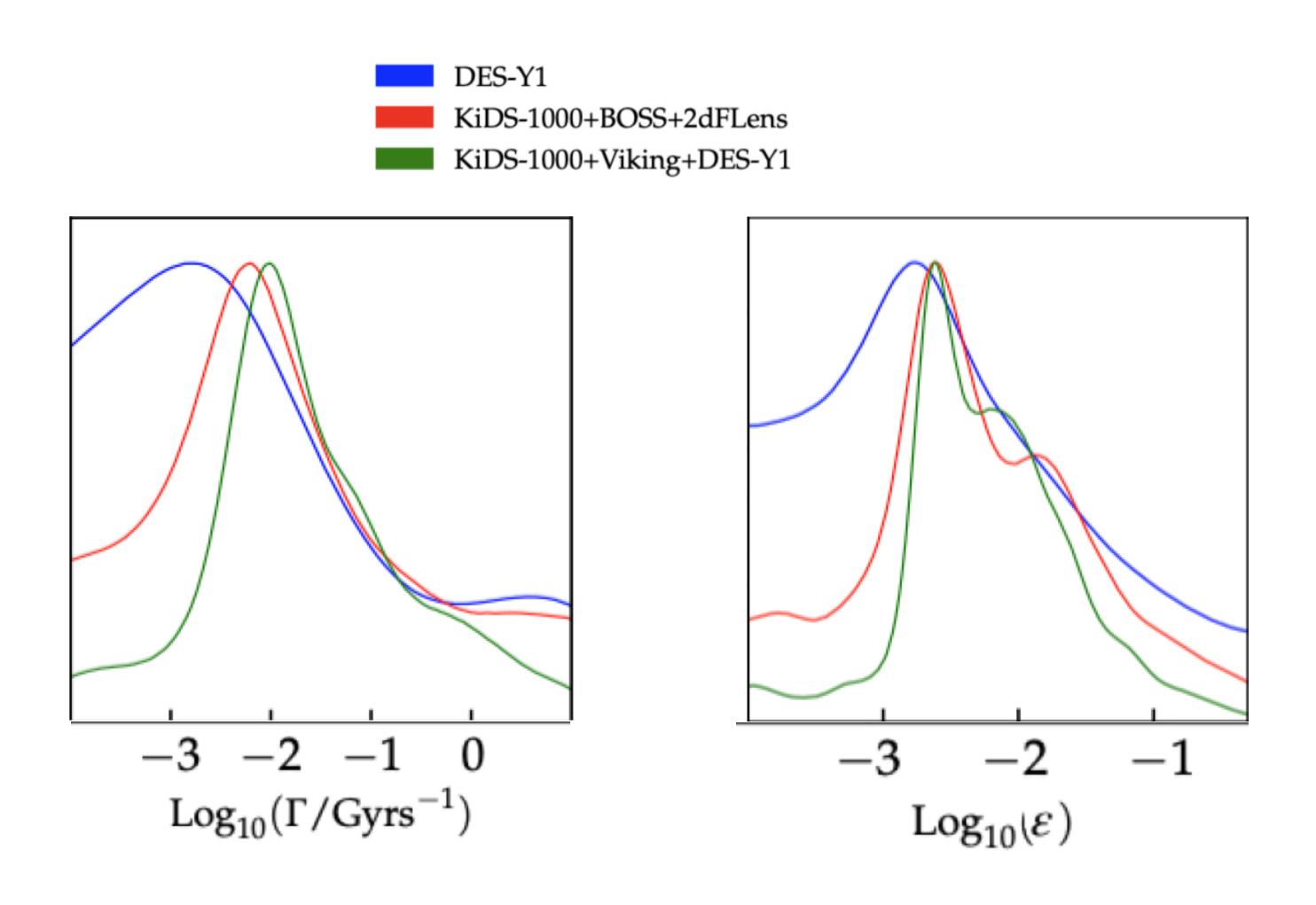
Low-*e*: enhanced Late Integrated Sachs Wolfe (LISW) effect

High-*e*: **suppressed** lensing (higher contrast between peaks)

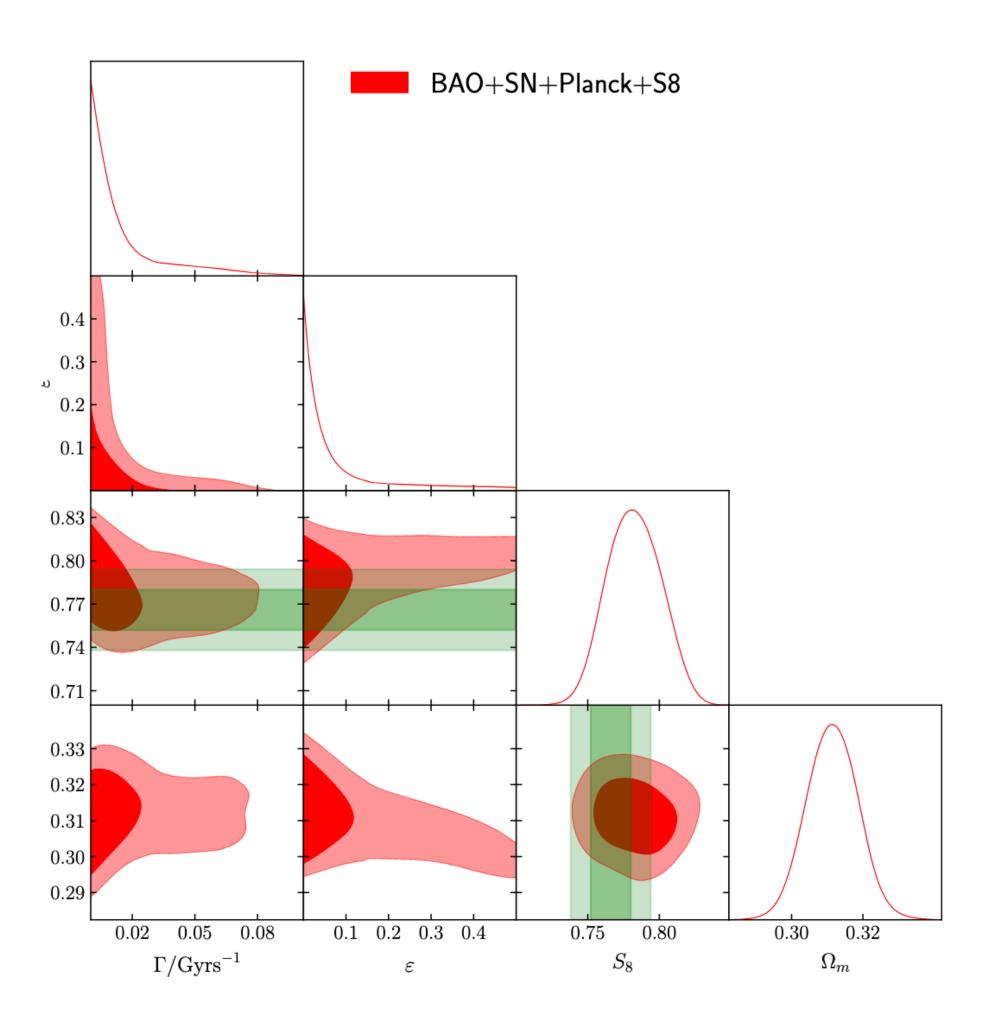


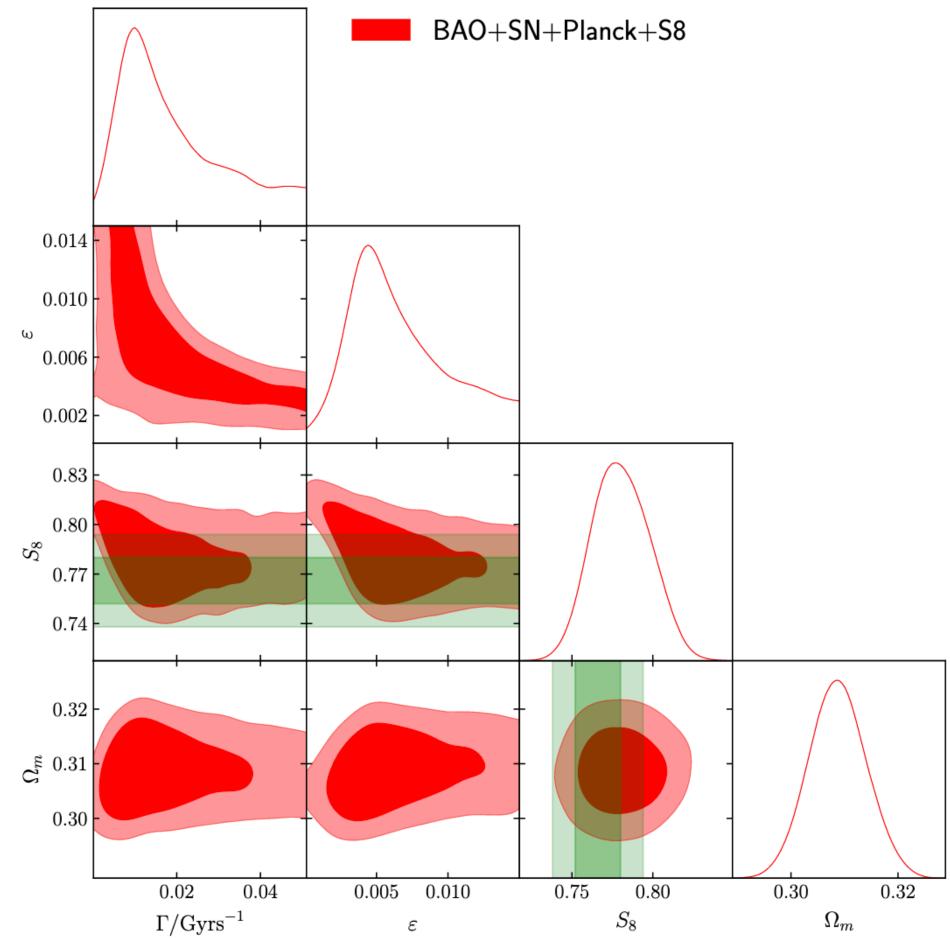
### DDM resolution to the S<sub>8</sub> tension

The level of detection depends on the level of tension with  $\Lambda CDM$ 

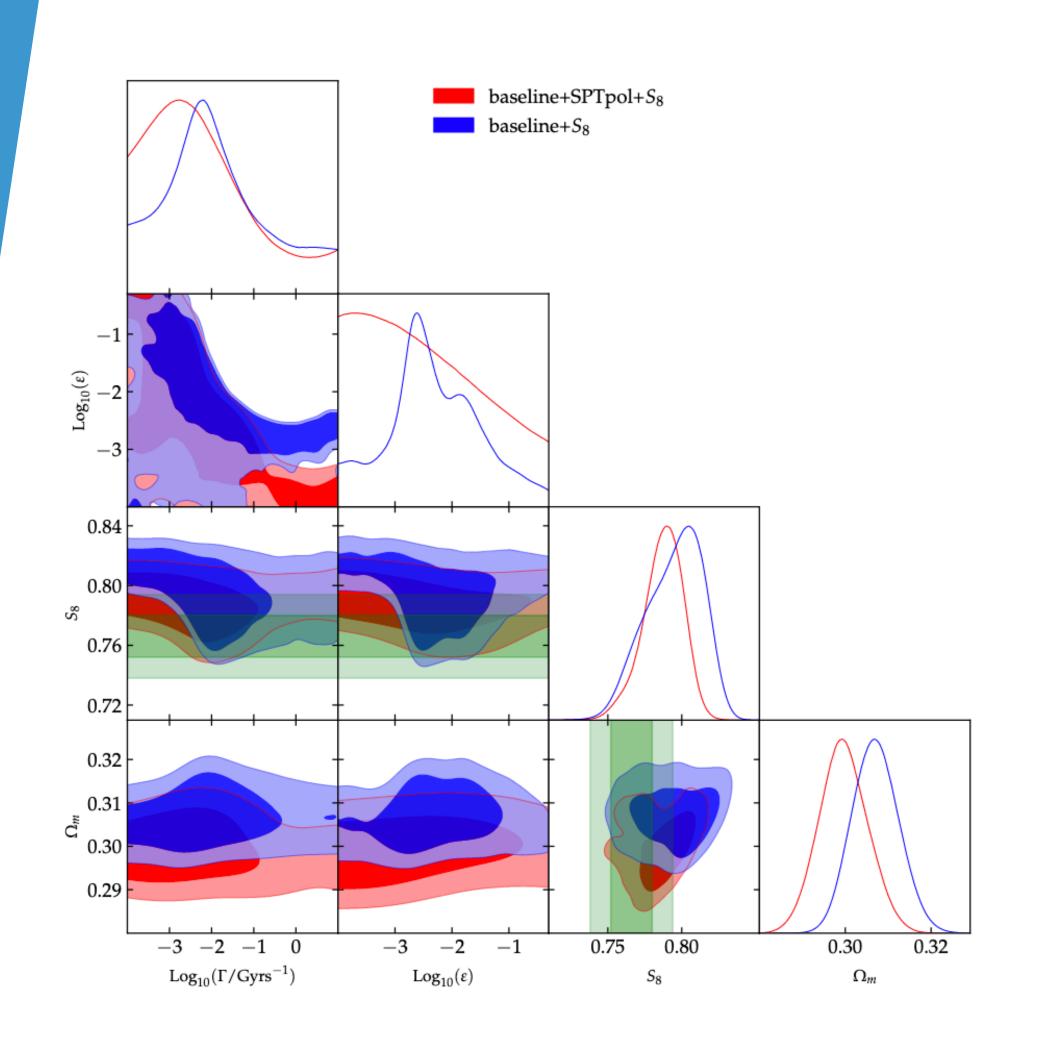


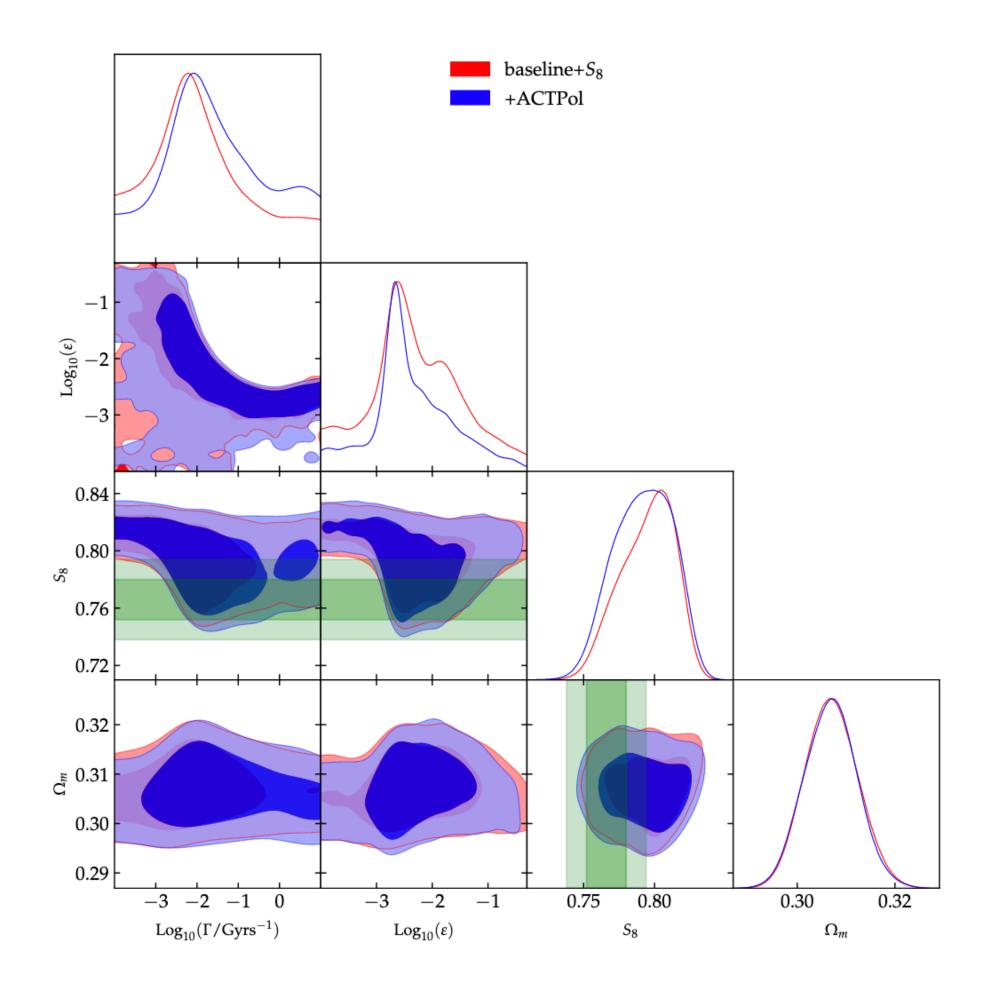
# DDM results with linear priors





### DDM results with SPTPol and ACT datasets





# DDM results marginalizing over lensing information

