

# Decaying dark matter: cosmological constraints and implications for the $S_8$ tension

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Based on:

arXiv:2102.12498 (PRD in press)

arXiv:2008.09615 (PRD in press)

with Riccardo Murgia, Vivian Poulin and Julien Lavalle



# Tensions in cosmology

With the era of precision cosmology, several discrepancies have emerged

- $S_8$  with weak-lensing data  
[KiDS-1000 2007.15632](#)
- $H_0$  with local measurements  
[Riess++ 2012.08534](#)

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## *Unaccounted systematics?*



- Less exotic explanation ✓
- Difficult to account for all discrepancies ✗

## *Physics beyond $\Lambda$ CDM?*



- Reveal properties about the dark sector ✓
- Very challenging ✗

# Tensions in cosmology

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- $S_8$  with weak-lensing data  
KiDS-1000 2007.15632  This talk
- $H_0$  with local measurements  
Riess++ 2012.08534  My previous works:  
Murgia, GFA, Poulin 2009.10733  
Schöneberg, GFA++ 2107.10291

## *Unaccounted systematics?*

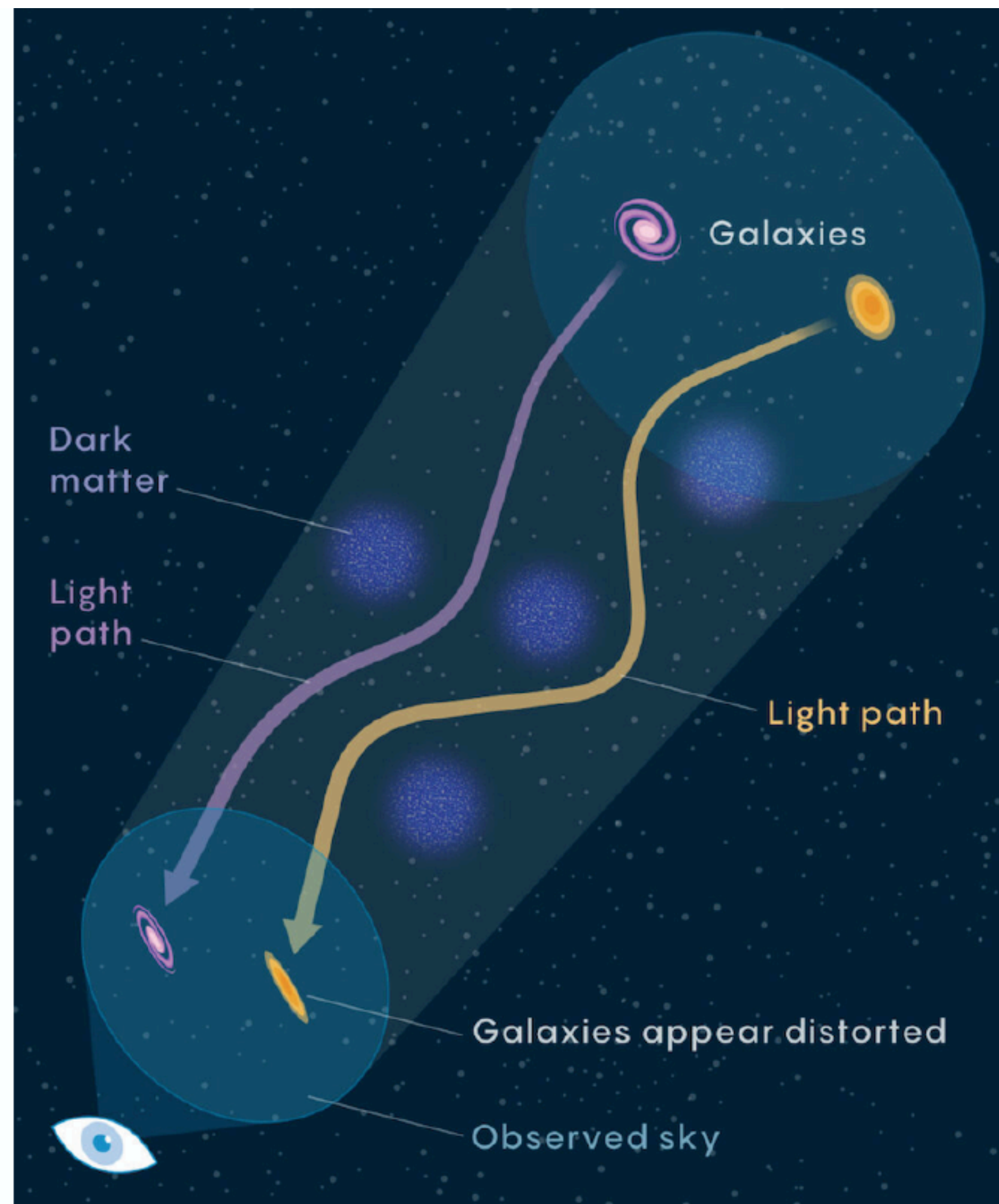
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## *Physics beyond $\Lambda$ CDM?*

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- Very challenging 

# The $S_8$ tension

Weak-lensing surveys are mainly sensible to  $S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$



KiDS+BOSS+2dfLenS\*:

$$S_8 = 0.766^{+0.020}_{-0.014}$$

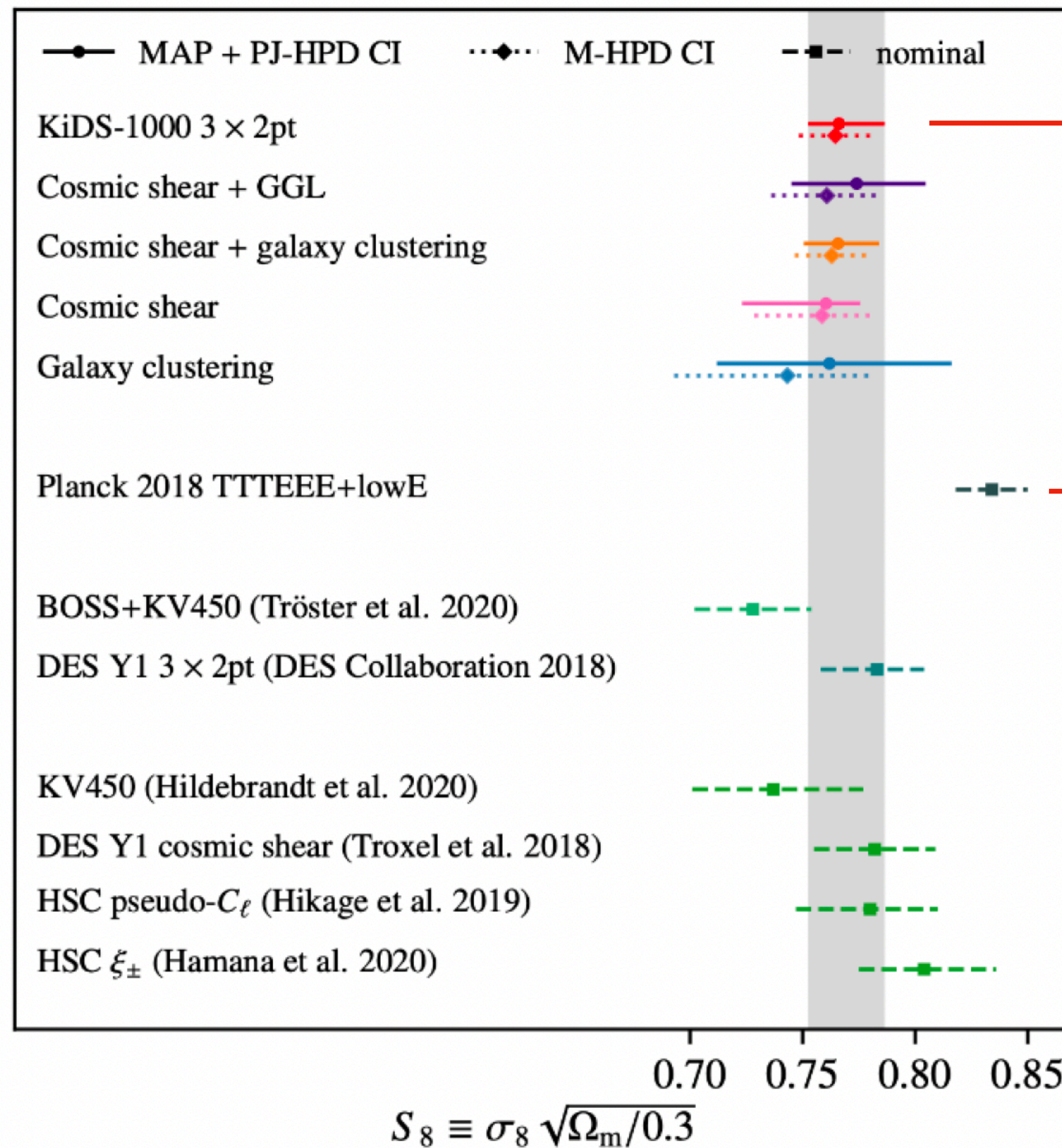
Planck (*under*  $\Lambda$ CDM):

$$S_8 = 0.830 \pm 0.013$$

→  $\sim 2 - 3\sigma$  tension

\*Other surveys such as DES, CFHTLenS or HSC yield similar results

# The $S_8$ tension



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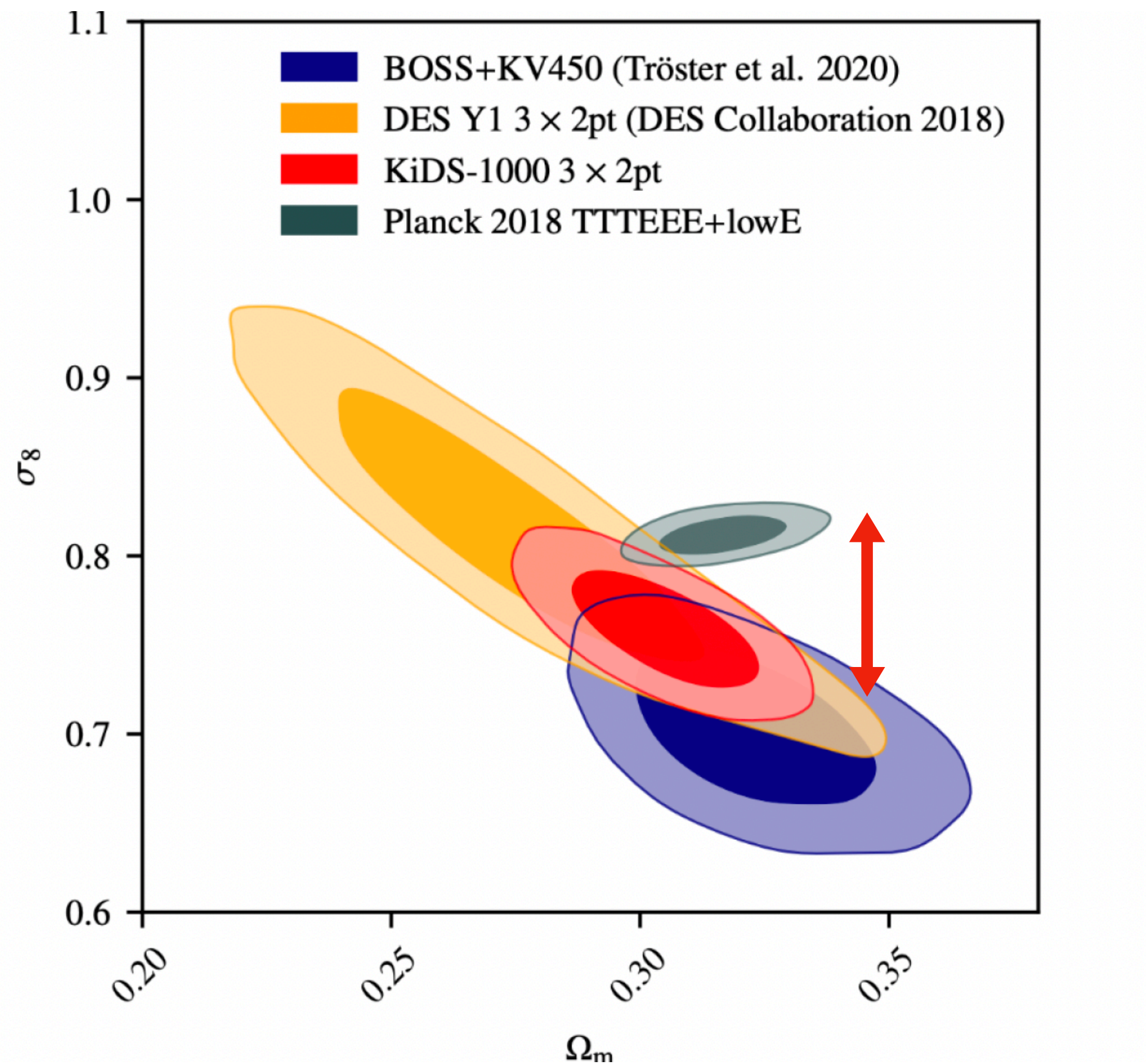
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# What is needed to resolve the $S_8$ tension?

Di Valentino++ 2008.11285

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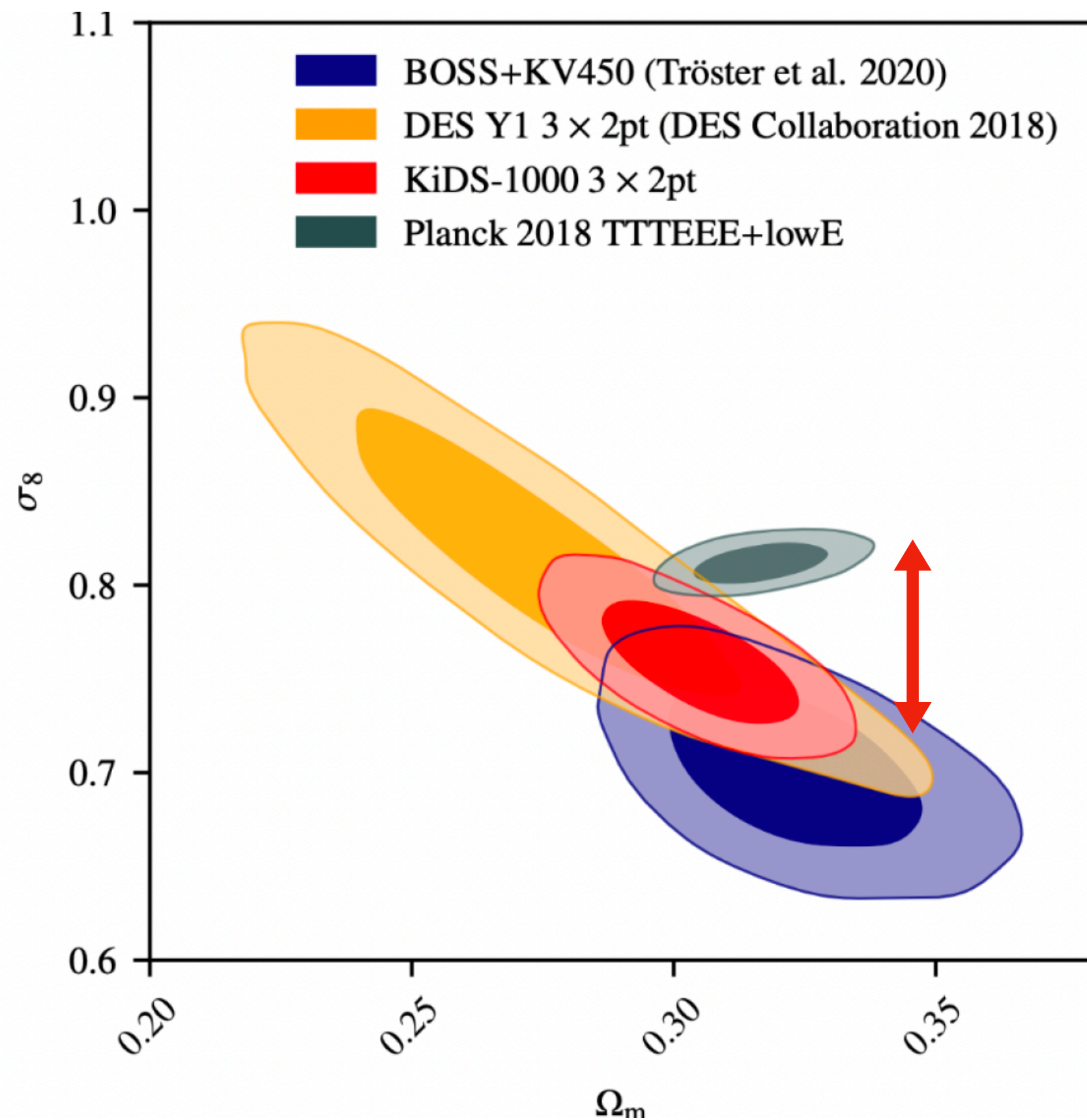
$\Omega_m$  should be left unchanged

$$\sigma_8 = \int P_m(k, z=0) W_R^2(k) d \ln k$$

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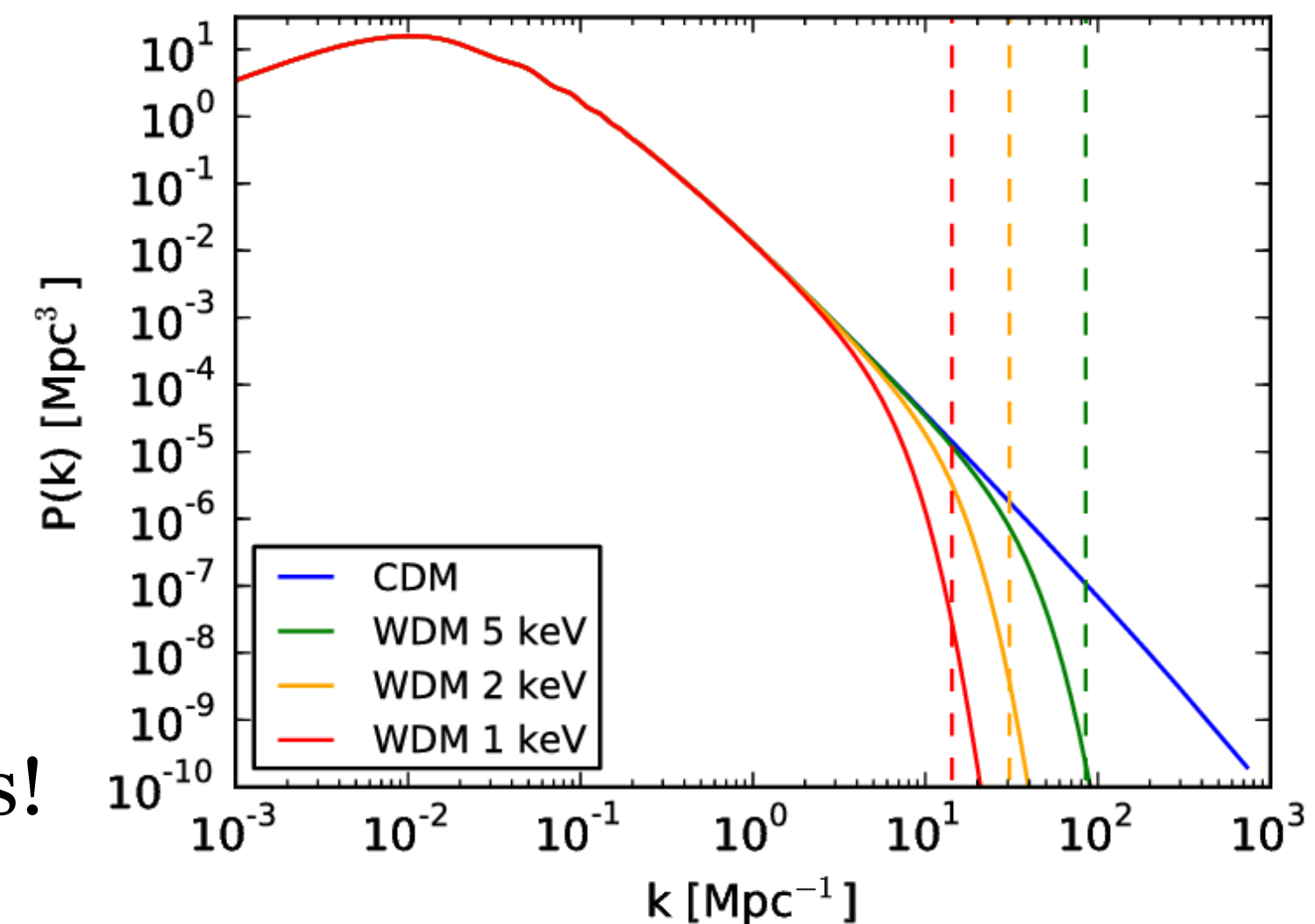


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Need to **suppress power** at scales  $k \sim 0.1 - 1 h/\text{Mpc}$



**Ex:** Warm Dark Matter  
Very constrained by many probes!

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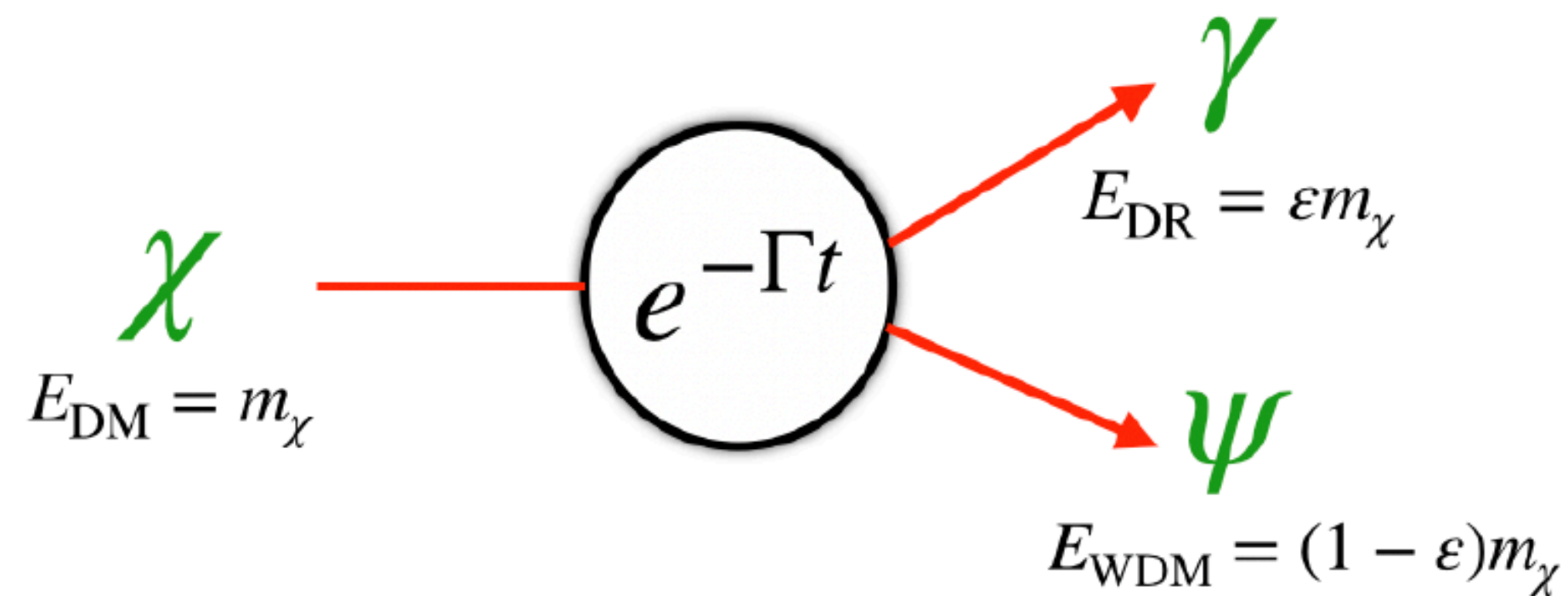
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- What about massive products?

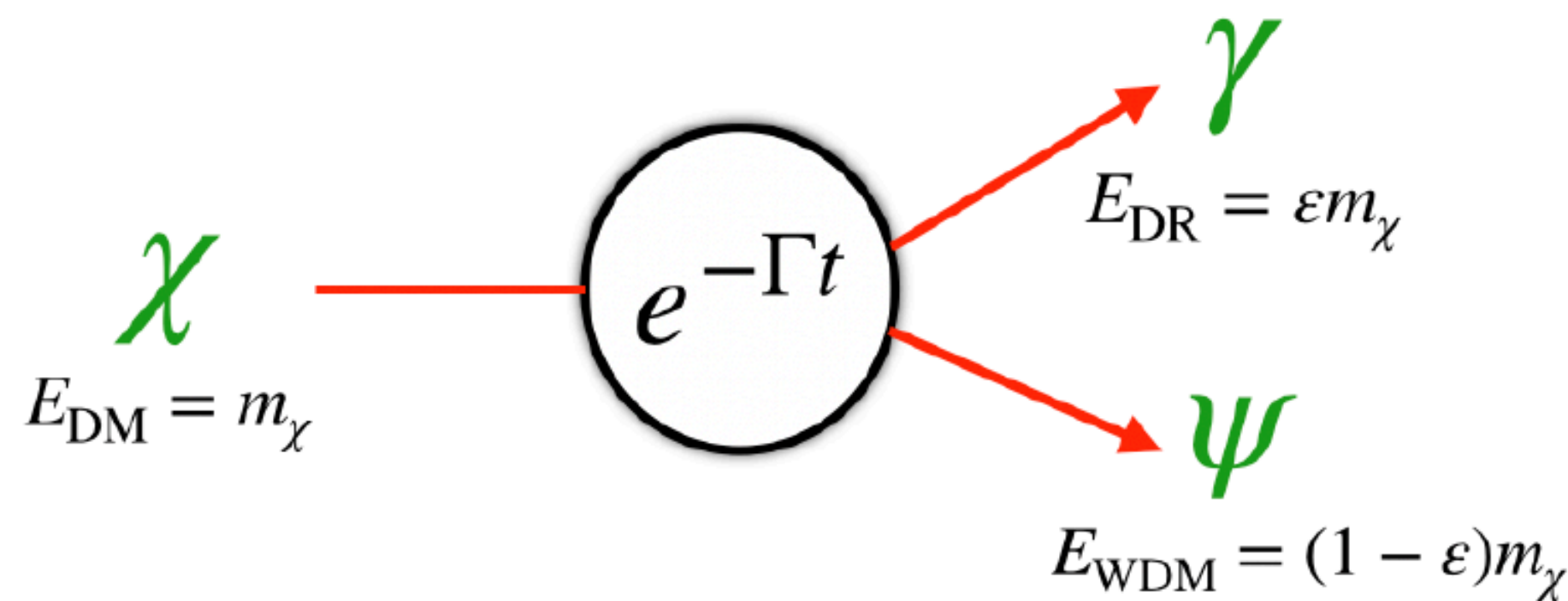
# 2-body Dark Matter decay

We explore DM decays to massless (**Dark Radiation**) and massive (**Warm Dark Matter**) particles,  $\chi(\text{DM}) \rightarrow \gamma(\text{DR}) + \psi(\text{WDM})$



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The model is fully specified by:

$$\{\Gamma, \varepsilon\} \quad \text{where} \quad \varepsilon = \frac{1}{2} \left( 1 - \frac{\mathbf{m}_\psi^2}{\mathbf{m}_\chi^2} \right) \begin{cases} = \mathbf{0} & \text{for } \Lambda\text{CDM} \\ = \mathbf{1/2} & \text{for DM} \rightarrow \text{DR} \end{cases}$$

# 2-body Dark Matter decay

Aoyama++ 1402.2972	→	Full treatment of perts.	No parameter scan
Vattis++ 1903.06220	→	Resolution to $H_0$ tension ?	No perturbations
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**Our goal:** Perform parameter scan by including full treatment of linear perts, in order to assess the impact on the  $S_8$  tension

# Evolution of perturbations: full treatment

- Effects on  $P_m(k)$  and  $C_\ell$  ? Track **linear perts.** for the particles species involved in the decay:  $\delta_i$ ,  $\theta_i$  and  $\sigma_i$  for  $i = \text{dm, dr, wdm}$

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- Boltzmann hierarchy of eqs. Dictate the evolution of the **p.s.d. multipoles**  $\Delta f_\ell(q, k, \tau)$ 
  - ◆ **DM and DR treatments are easy**, momentum d.o.f. are integrated out
  - ◆ **For WDM**, one needs to follow the evolution of the full p.s.d.  
Computationally expensive  $\longrightarrow \mathcal{O}(10^8)$  ODEs to solve!

# Evolution of perturbations: fluid equations

New fluid eqs.\*, based on previous approximation for massive neutrinos

Lesgourgues & Tram, 1104.2935

$$\dot{\delta}_{\text{wdm}} = -3aH(c_{\text{syn}}^2 - w)\delta_{\text{wdm}} - (1 + w)\left(\theta_{\text{wdm}} + \frac{\dot{h}}{2}\right) + a\Gamma(1 - \varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}(\delta_{\text{dm}} - \delta_{\text{wdm}})$$

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where

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and

$$c_{\text{syn}}^2(k, \tau) = c_a^2(\tau)\left[1 + (1 - 2\varepsilon)T(k/k_{\text{fs}})\right]$$

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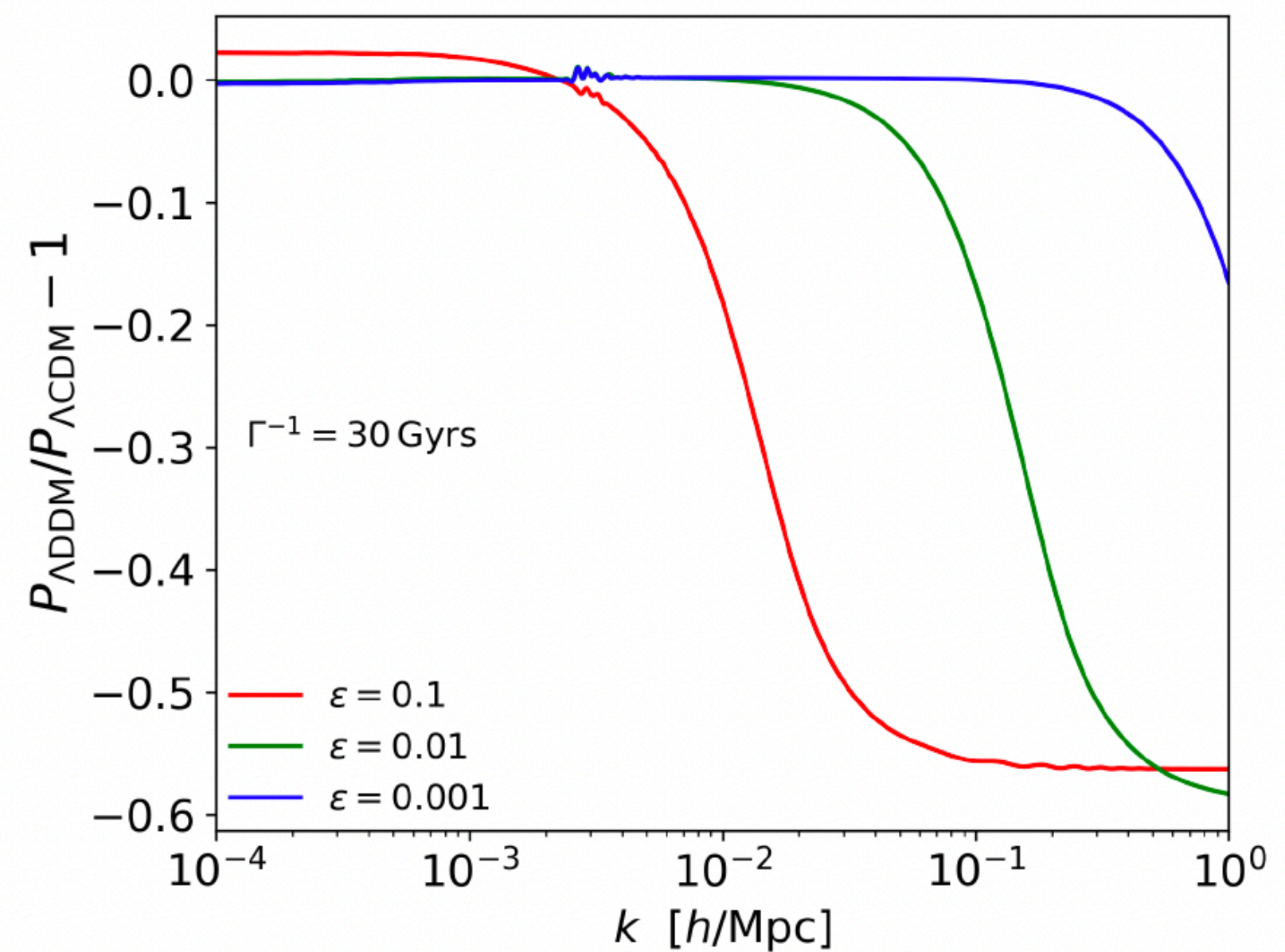
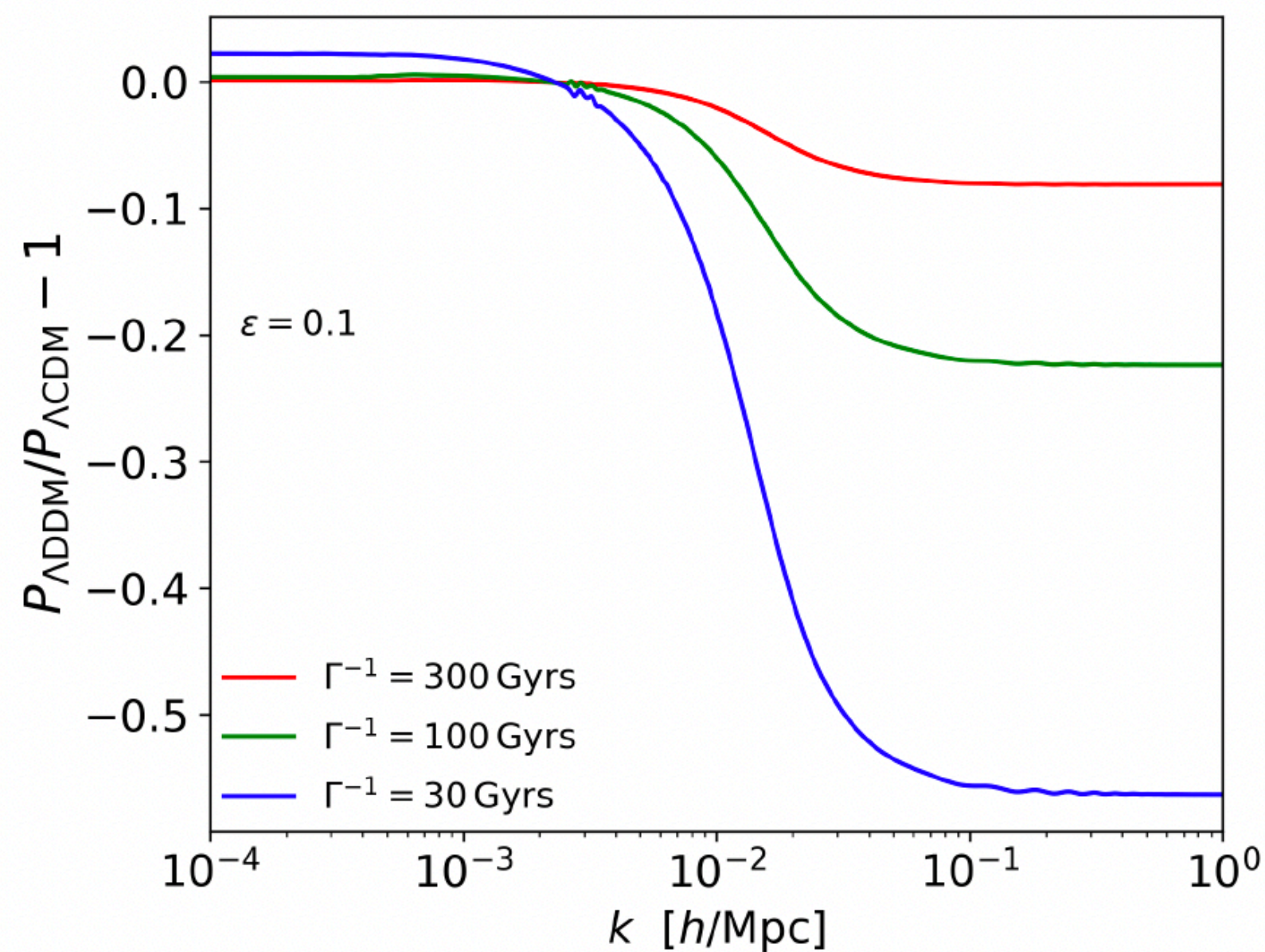
$$c_{\text{syn}}^2(k, \tau) = c_a^2(\tau)\left[1 + (1 - 2\varepsilon)T(k/k_{\text{fs}})\right]$$

***CPU time reduced from  $\sim 1$  day to  $\sim 1$  minute!***

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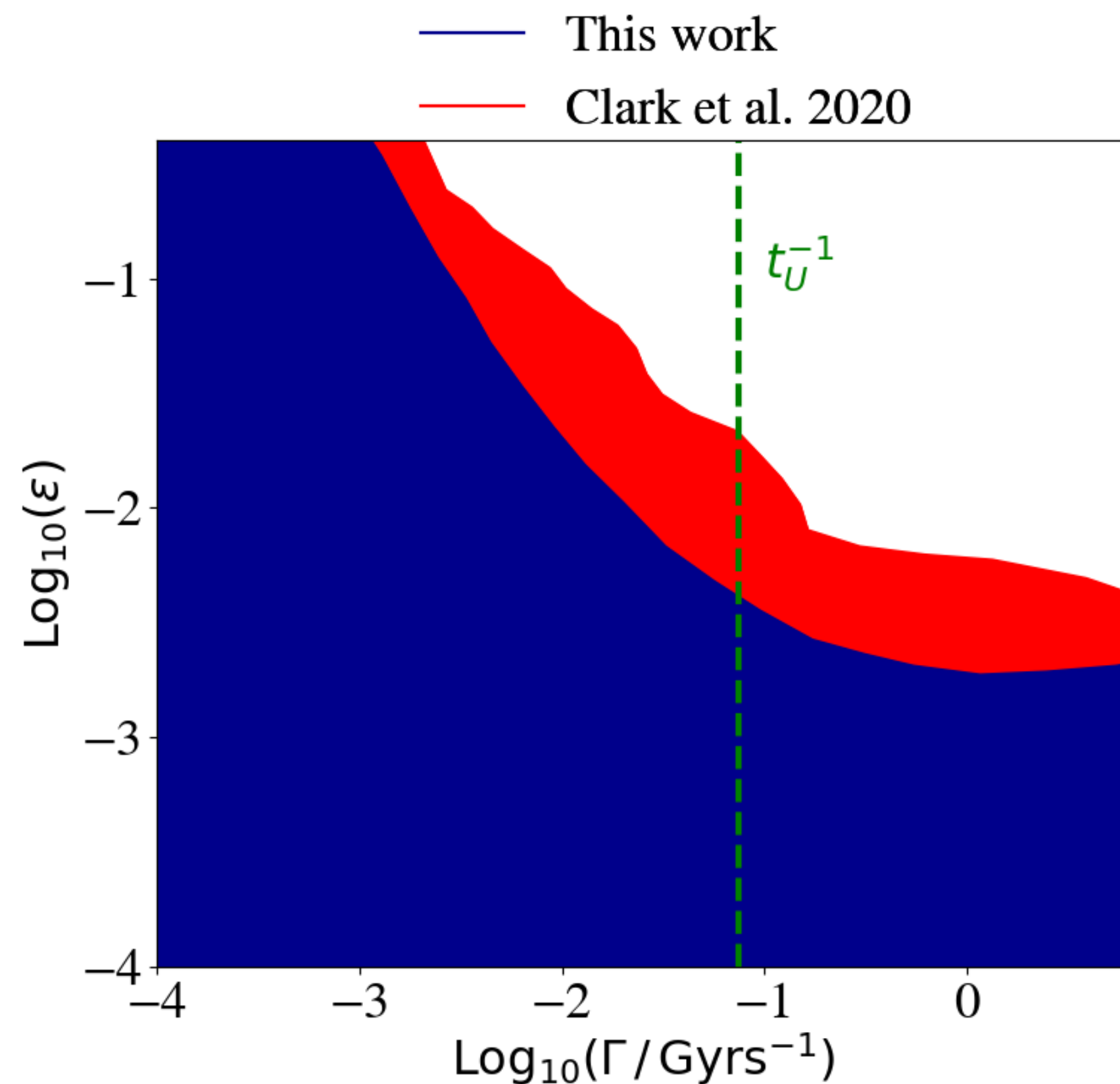
# Impact of decaying DM on the matter spectrum

The WDM daughter leads to a power suppression in  $P_m(k)$  at small scales  $k > k_{\text{fs}}$ , where  $k_{\text{fs}} \sim aH/c_a$



# General constraints on the 2-body DM decay

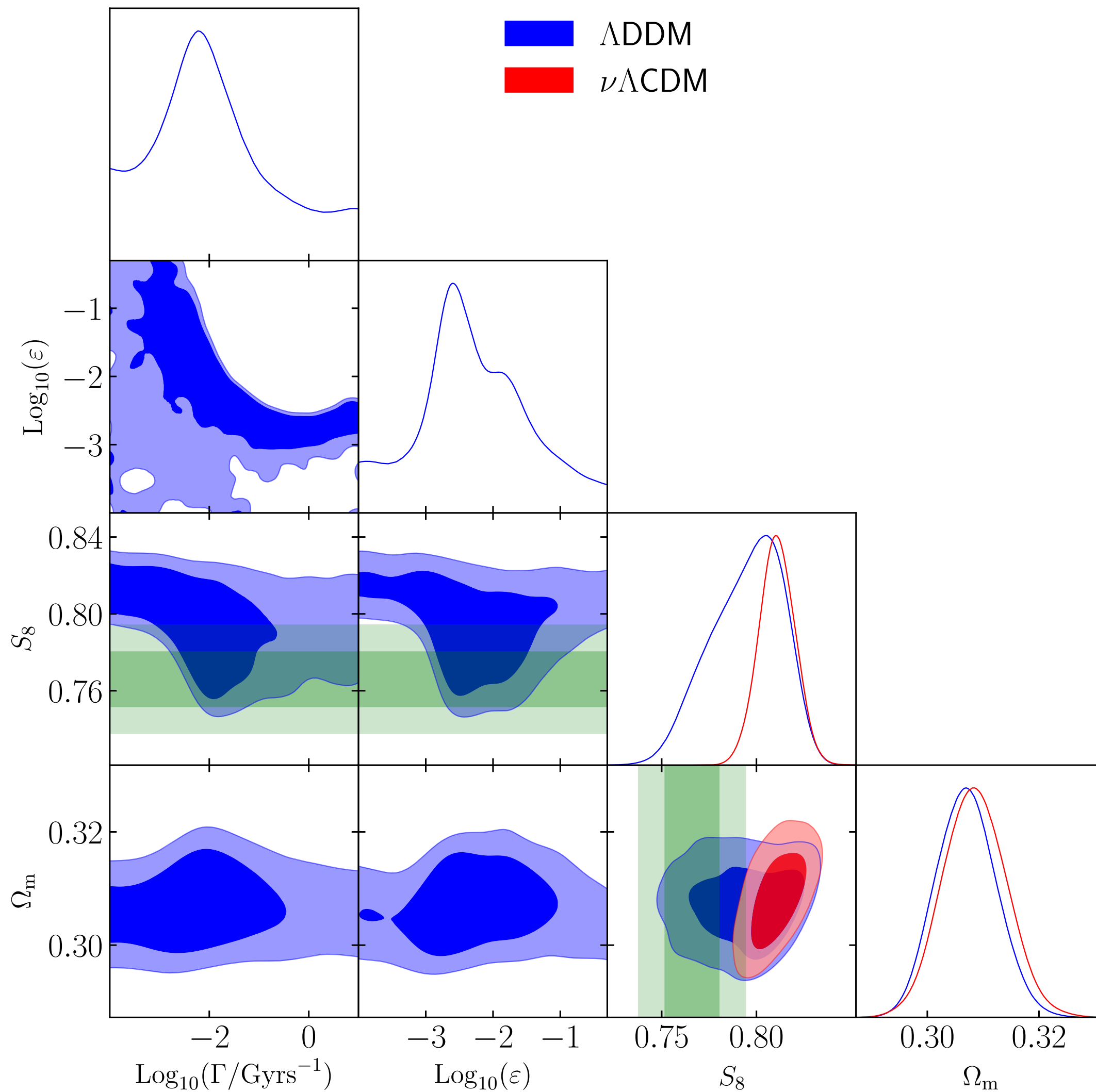
Planck+BAO+SNIa analysis



Strong **negative correlation** between  $\epsilon$  and  $\Gamma$

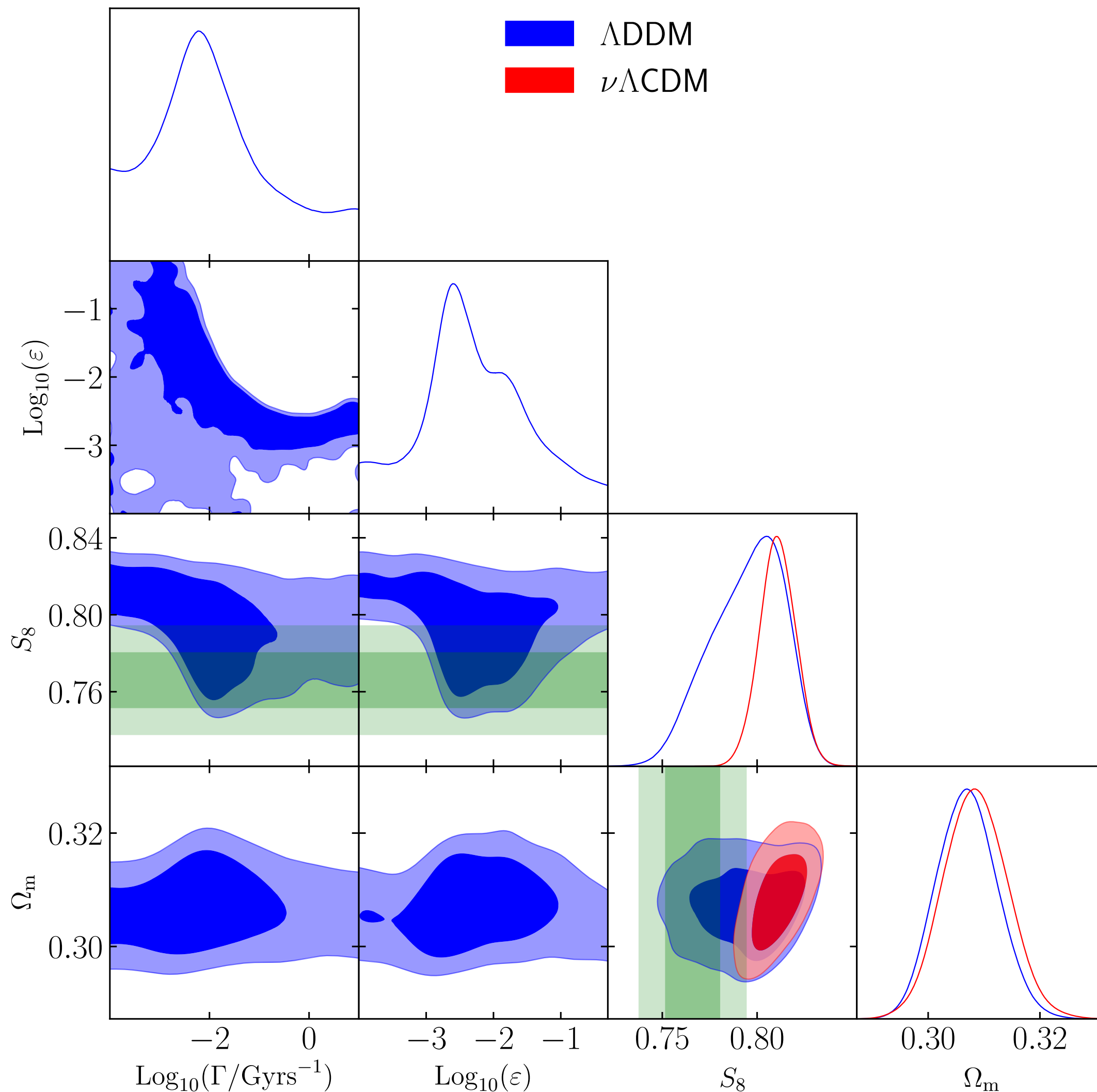
Constraints up to **1 order of magnitude stronger** than previous literature

# Explaining the $S_8$ tension



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- MCMC analysis using Planck+BAO+SNIa+prior on  $S_8$  from KIDS+BOSS+2dfLenS
- Reconstructed  $S_8$  values are in excellent agreement with WL data!

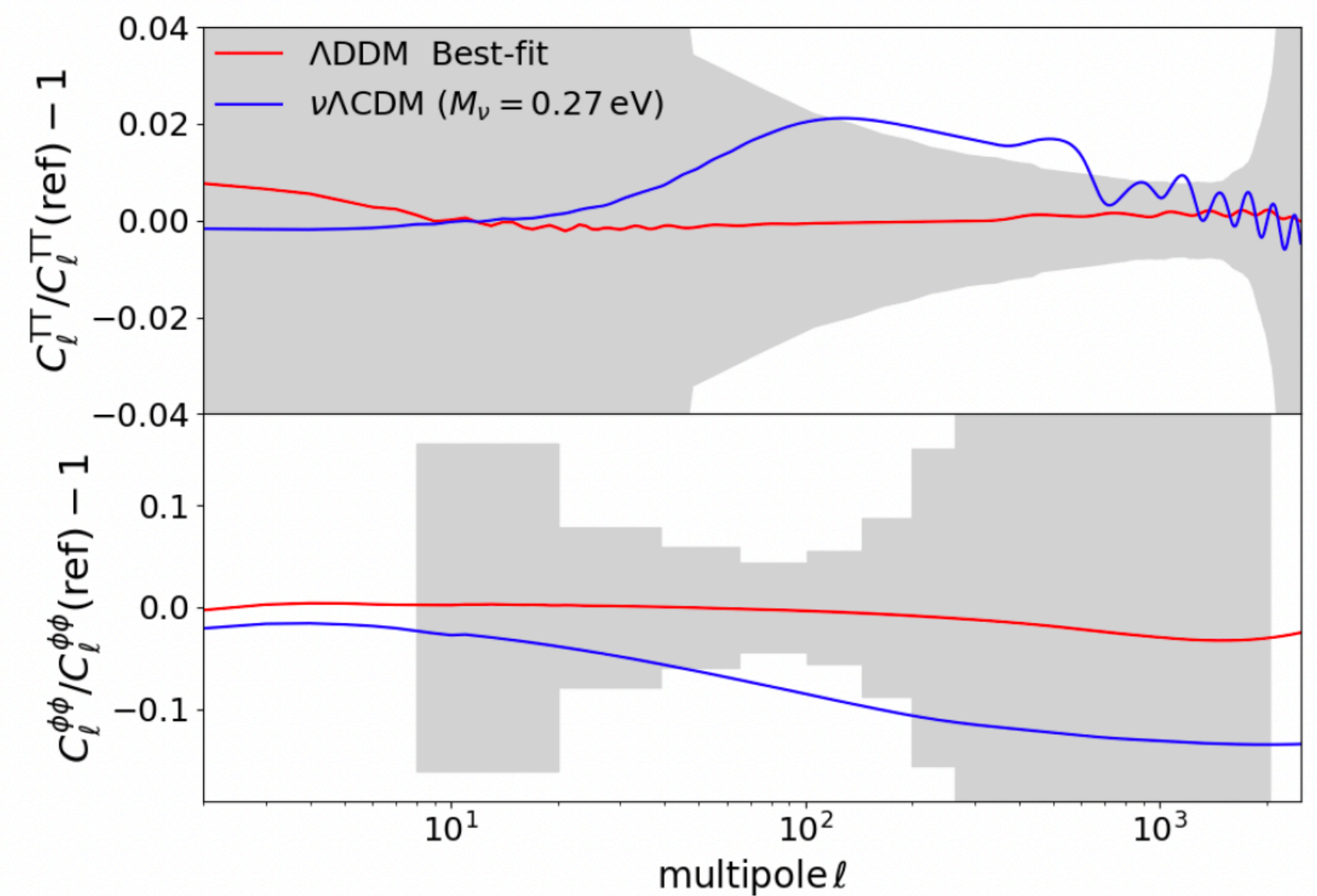
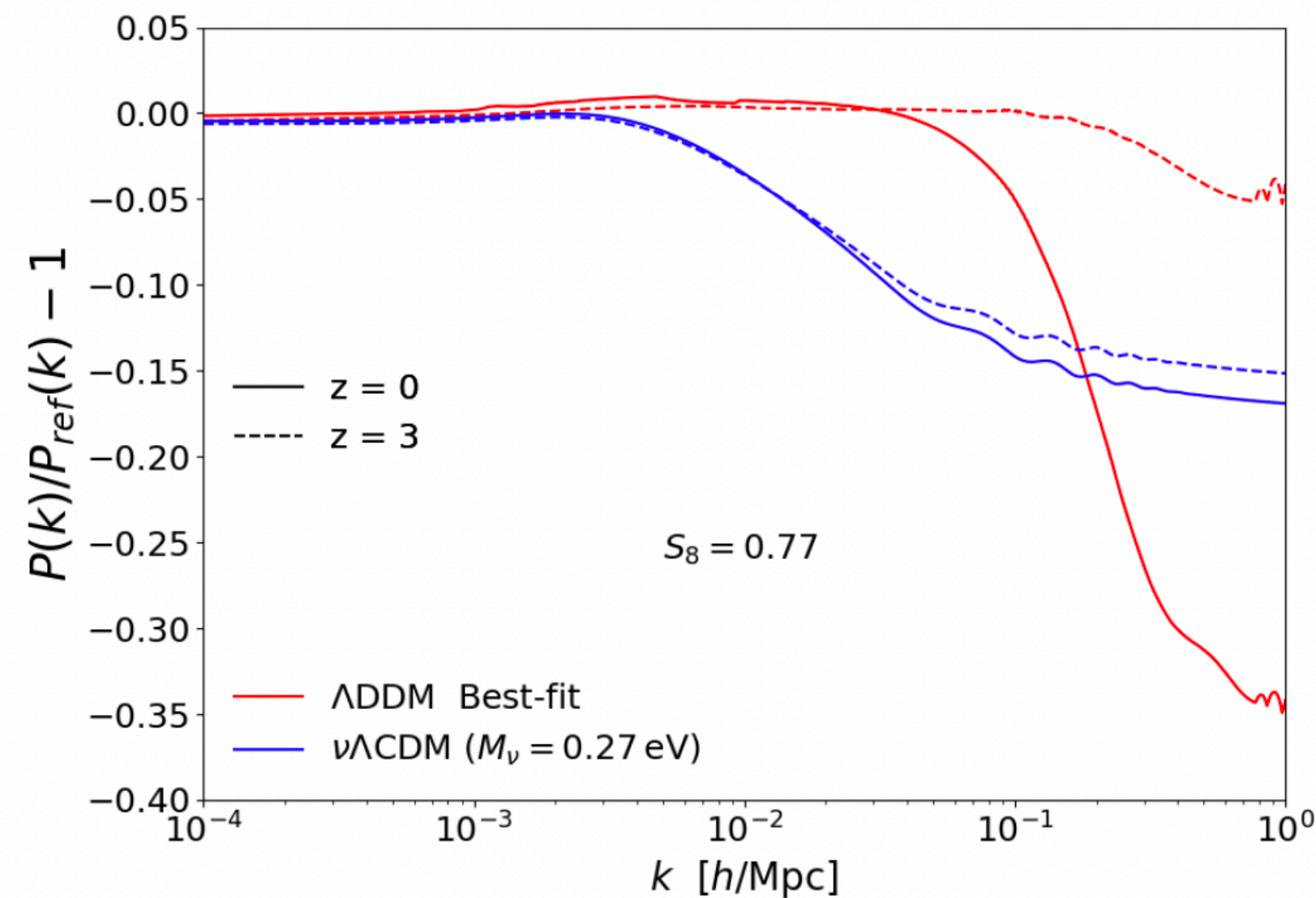
	$\nu\Lambda$ CDM	$\Lambda$ DDM
$\chi^2_{\text{CMB}}$	1015.9	1015.2
$\chi^2_{S_8}$	5.64	0.002

$$\longrightarrow \Delta\chi^2_{\min} \simeq -5.5$$

$$\Gamma^{-1} \simeq 55 (\varepsilon/0.007)^{1.4} \text{ Gyr}$$

# Why does the 2-body DM decay work better than massive neutrinos?

The 2-body decay gives a better fit thanks to the **time-dependence of the power suppression** and the cut-off scale



# Interesting implications

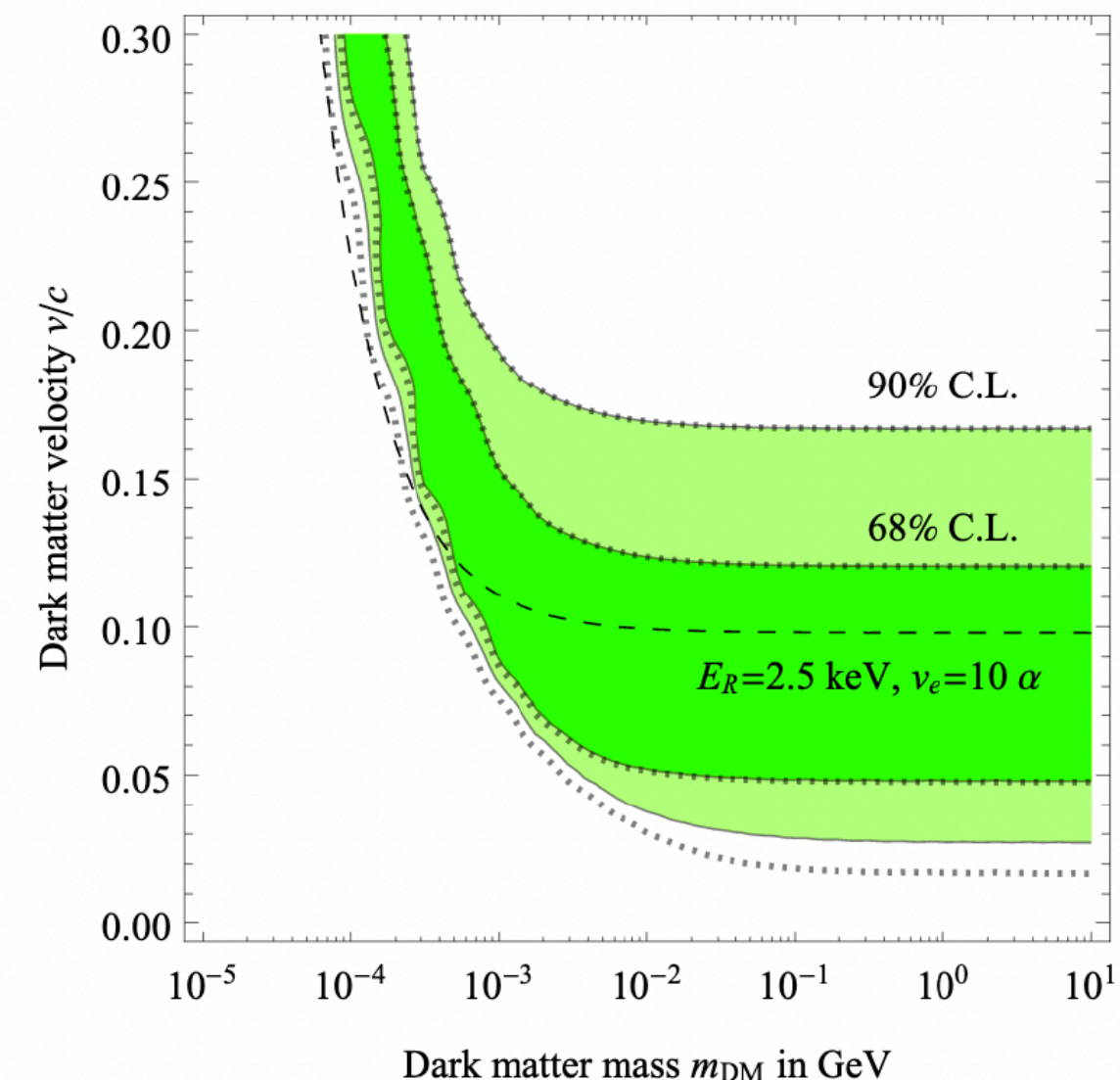
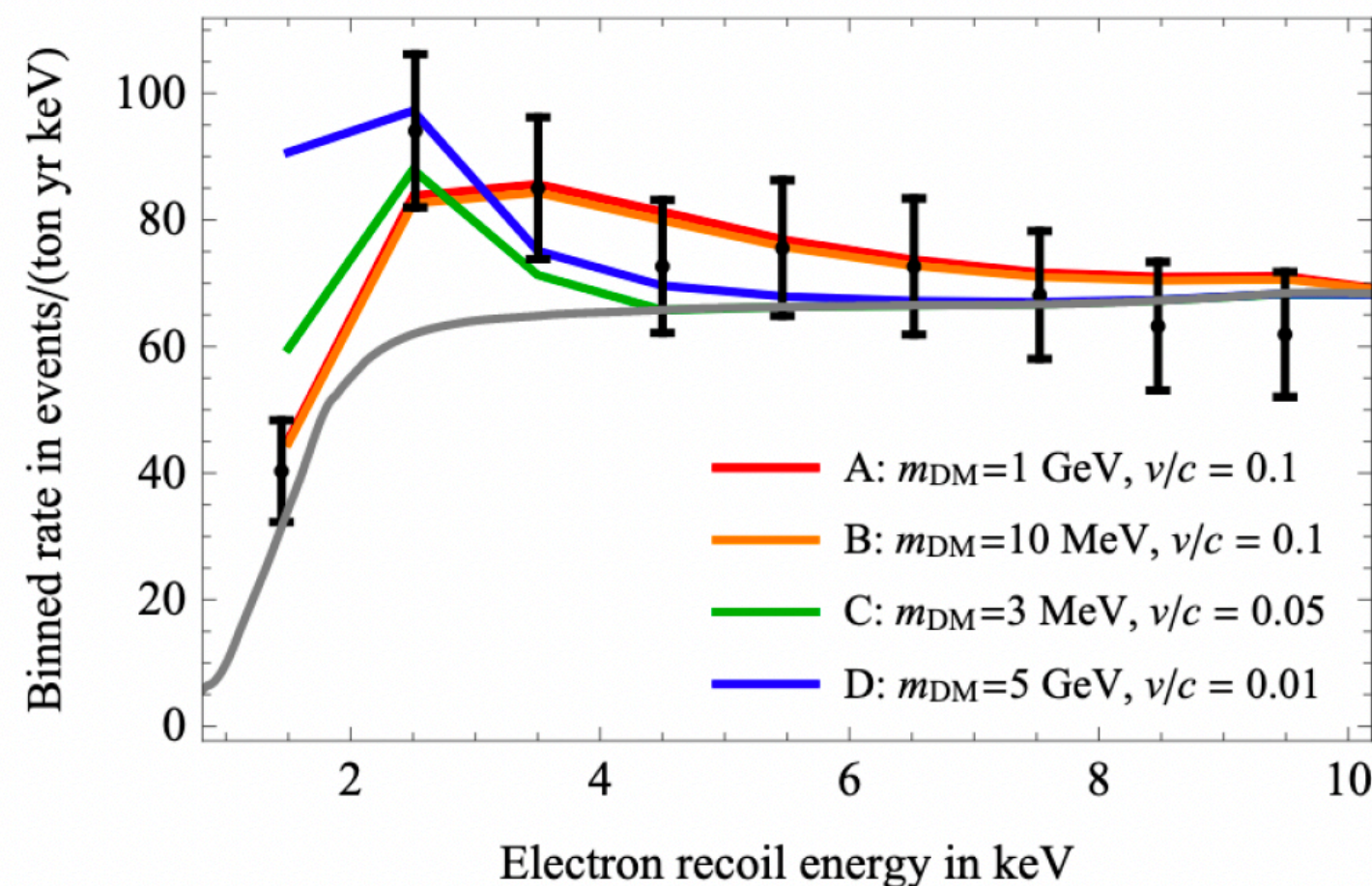
- **Model building:** Why  $\varepsilon \ll 1/2$ , i.e.  $m_{\text{wdm}} \sim m_{\text{dm}}$  ?  
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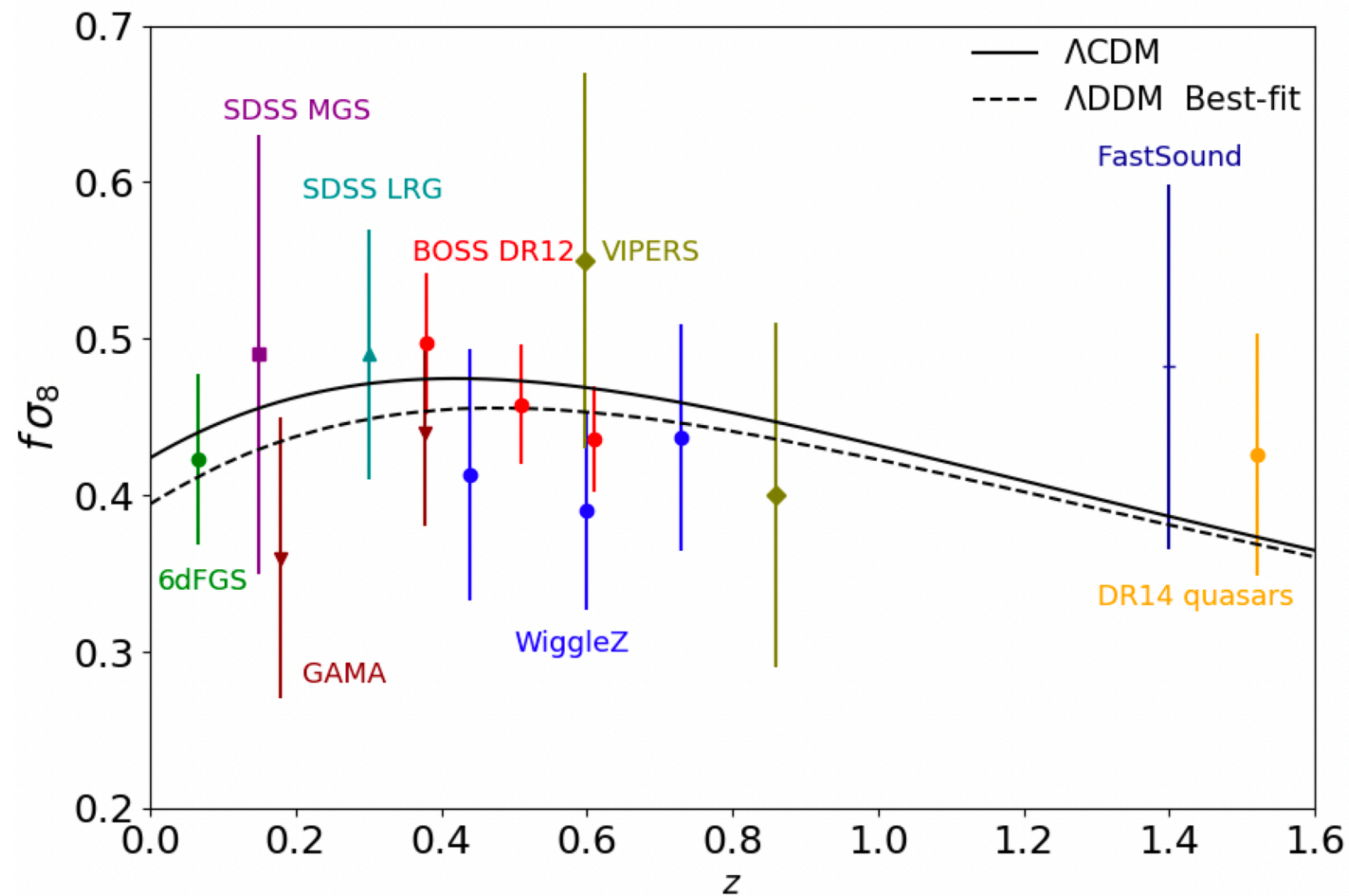
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- **Small-scale crisis of  $\Lambda$ CDM:** Reduction in the abundance of subhalos and their concentrations [Wang++ 1406.0527](#)
- **Xenon-1T excess:** It could be explained by a fast DM component, such as the WDM, with  $v/c \simeq \varepsilon$  [Kannike++ 2006.10735](#)



# Prospects for the 2-body DM decay



Accurate measurements of  $f\sigma_8$  at  $0 \lesssim z \lesssim 1$  will further test the 2-body decay

**Next goal:** Predict non-linear matter power spectrum (using either N-body simulations or EFT of LSS)

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- It could have interesting implications for model building, the small-scale crisis, and the recent **Xenon-1T excess**
- **Future growth factor** measurements can further test this scenario

# **BACK-UP SLIDES**

# The full Boltzmann hierarchy

$$f(q, k, \mu, \tau) = \bar{f}(q, \tau) + \delta f(q, k, \mu, \tau)$$

Expand  $\delta f$  in multipoles. The Boltzmann eq. leads to the following **hierarchy** (in *synchronous* gauge comoving with the mother)

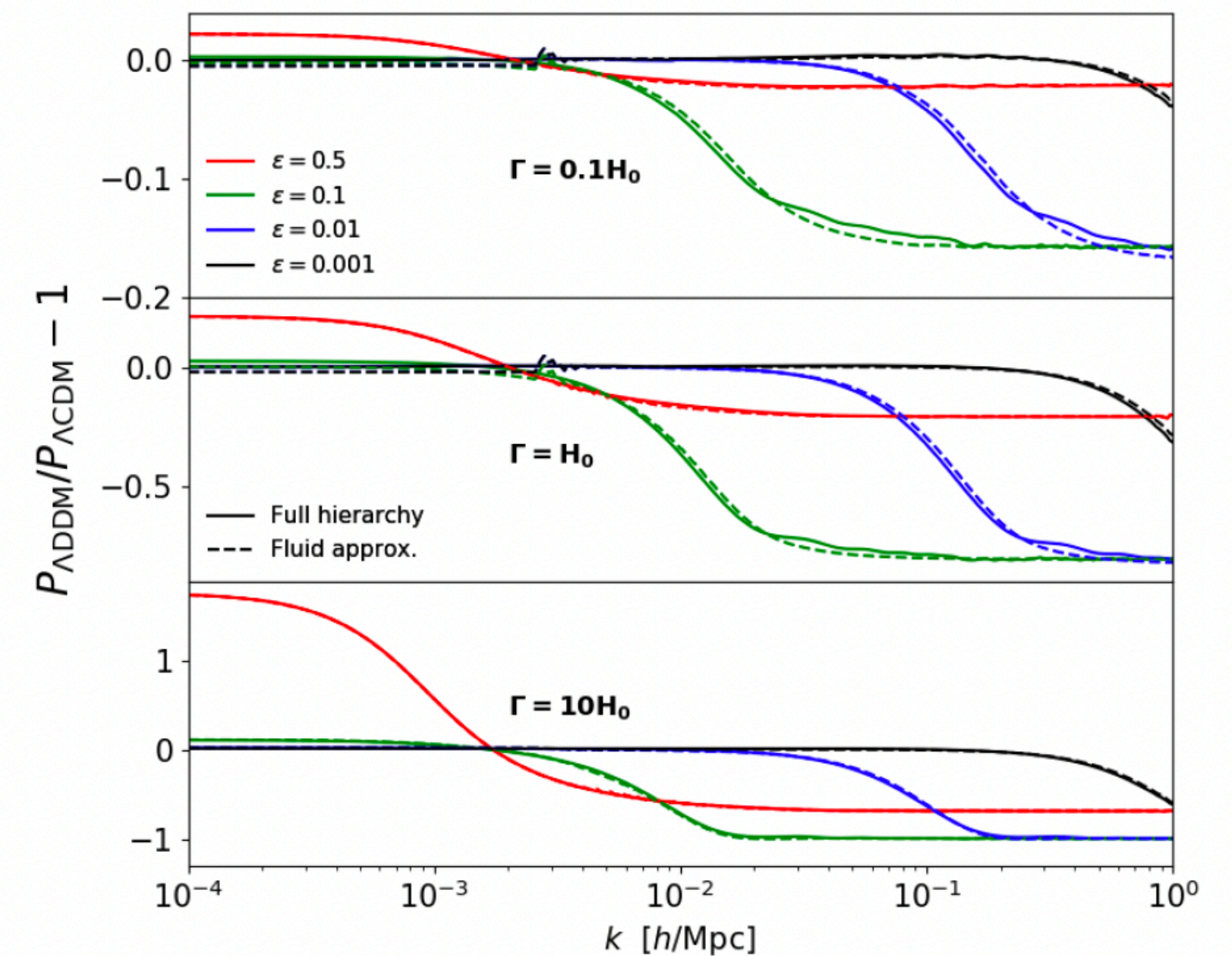
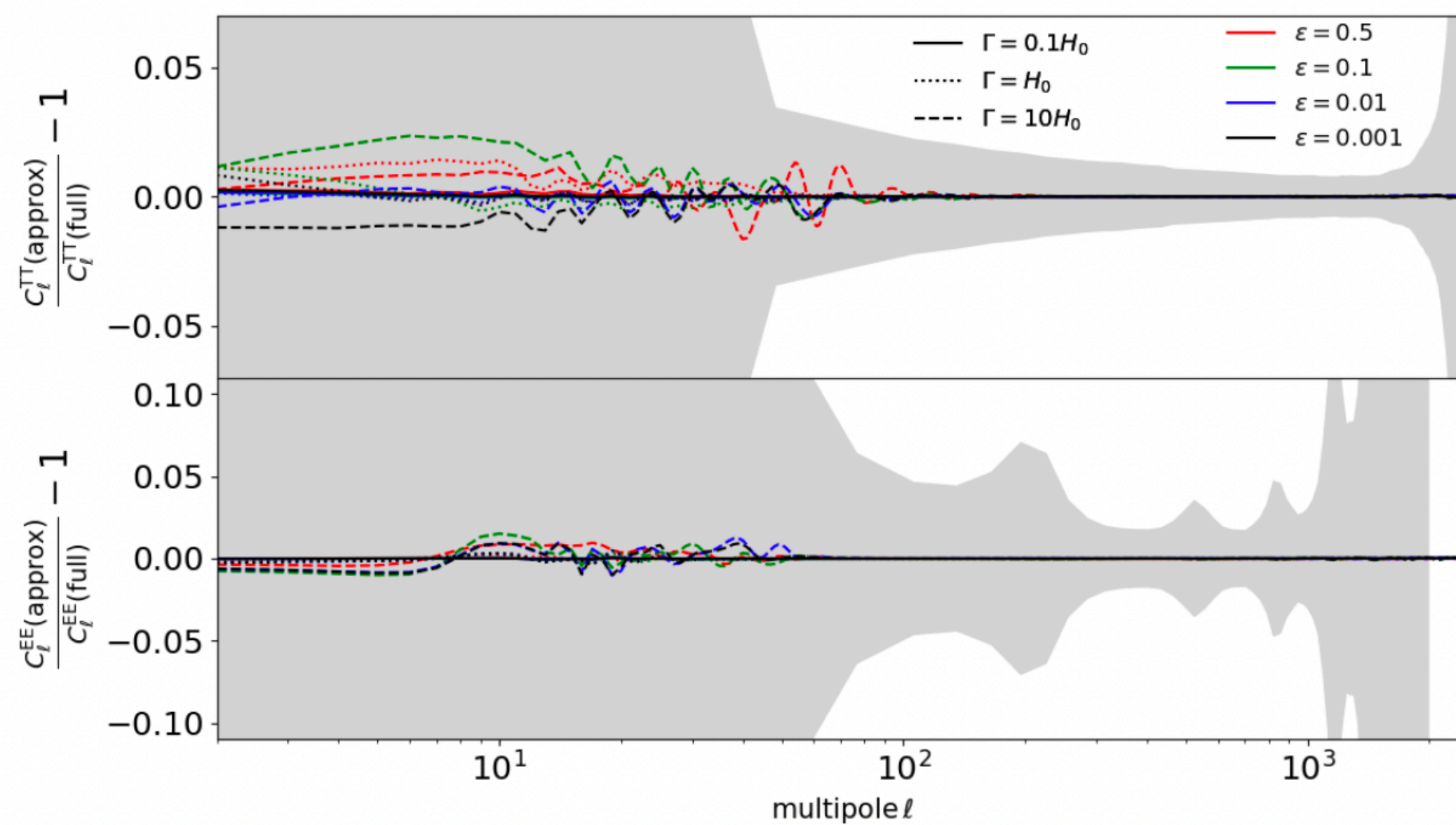
$$\begin{aligned} \frac{\partial}{\partial \tau} (\delta f_0) &= -\frac{\mathbf{q}k}{a\mathbf{E}} \delta f_1 + q \frac{\partial \bar{f}}{\partial q} \frac{\dot{h}}{6} + \frac{\Gamma \bar{N}_{\text{dm}}(\tau)}{4\pi q^3 H} \delta(\tau - \tau_q) \delta_{\text{dm}}, \\ \frac{\partial}{\partial \tau} (\delta f_1) &= \frac{\mathbf{q}k}{3a\mathbf{E}} [\delta f_0 - 2\delta f_2], \\ \frac{\partial}{\partial \tau} (\delta f_2) &= \frac{\mathbf{q}k}{5a\mathbf{E}} [2\delta f_1 - 3\delta f_3] - q \frac{\partial \bar{f}}{\partial q} \frac{(\dot{h} + 6\dot{\eta})}{15}, \\ \frac{\partial}{\partial \tau} (\delta f_\ell) &= \frac{\mathbf{q}k}{(2\ell + 1)a\mathbf{E}} [\ell \delta f_{\ell-1} - (\ell + 1) \delta f_{\ell+1}] \quad (\text{for } \ell \geq 3). \end{aligned}$$

where  $q = a(\tau_q) p_{\text{max}}$ . In the relat. limit  $\mathbf{q}/a\mathbf{E} = 1$ , so one can take

$$F_\ell \equiv \frac{4\pi}{\rho_c} \int dq \, q^3 \delta f_\ell \quad \text{and integrate out the dependency on } \mathbf{q}$$

# Checking the accuracy of the WDM fluid approx.

We compare the full Boltzmann hierarchy calculation with the WDM fluid approx.

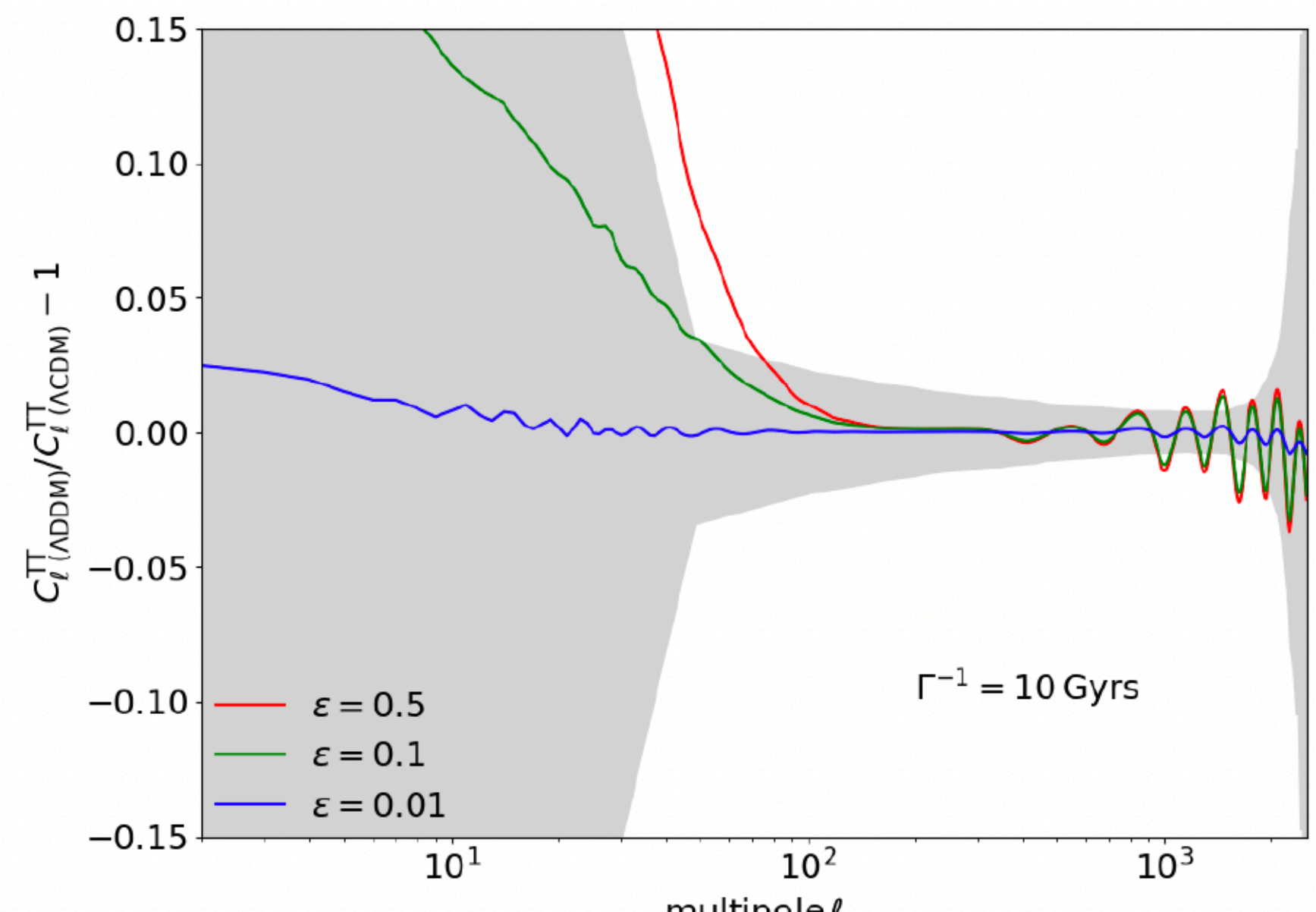
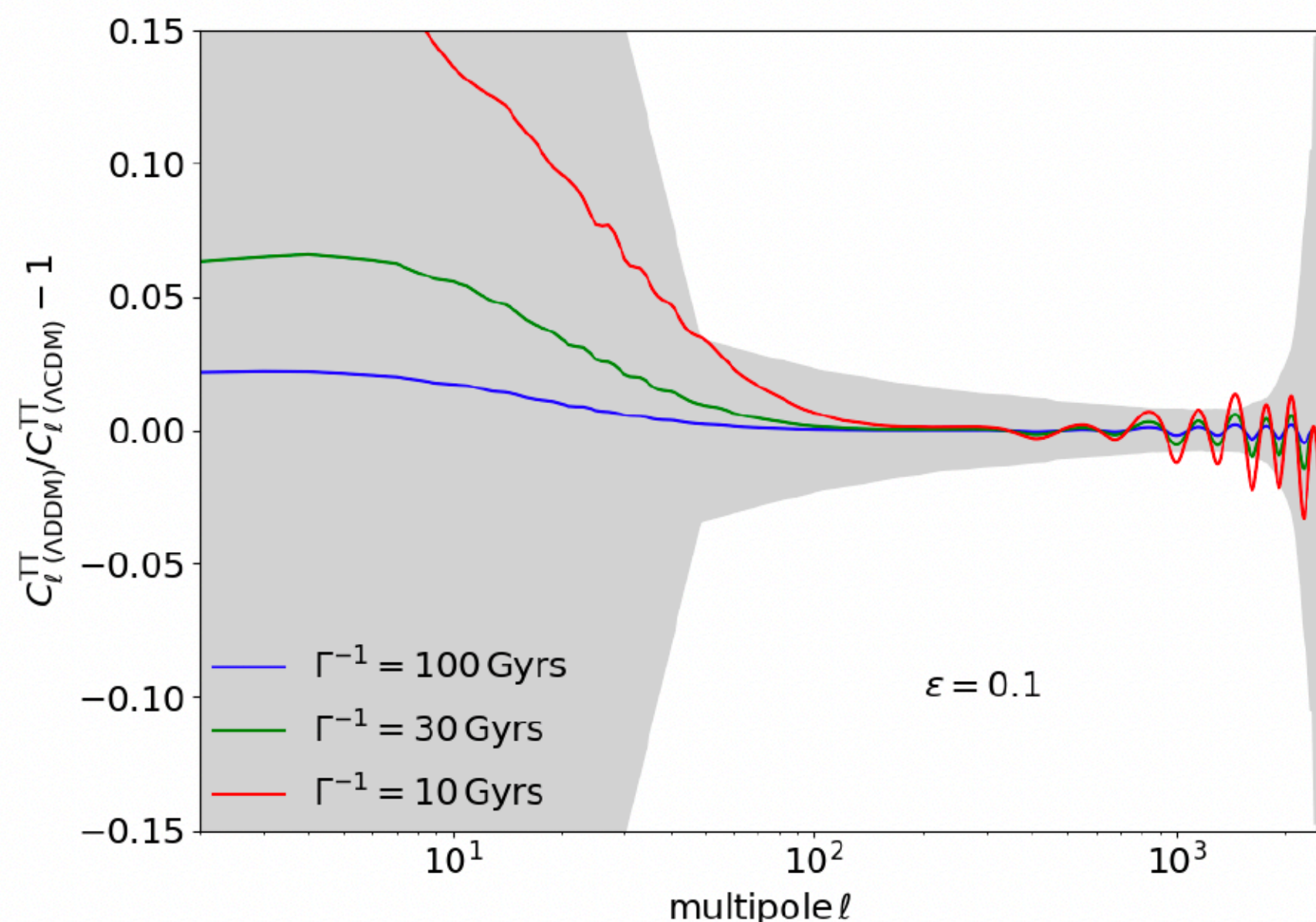


The max. error on  $S_8$  is  $\sim 0.65\%$ , smaller than the  $\sim 1.8\%$  error of the measurement from BOSS+KiDS+2dfLenS

# Impact on the CMB temperature spectrum

Low- $\ell$  : **enhanced** Late Integrated Sachs Wolfe (**LISW**) effect

High- $\ell$  : **suppressed** lensing (higher contrast between peaks)

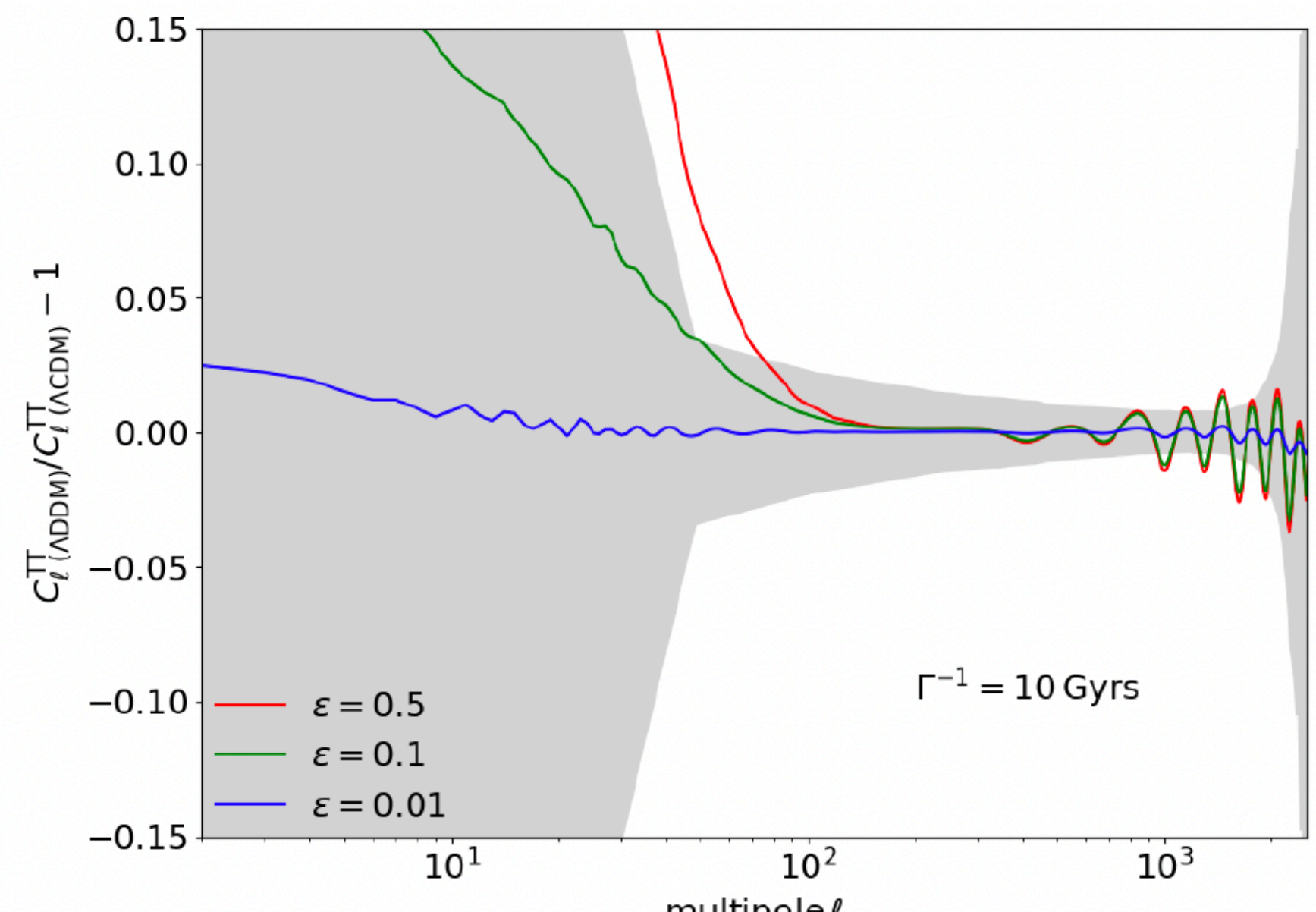
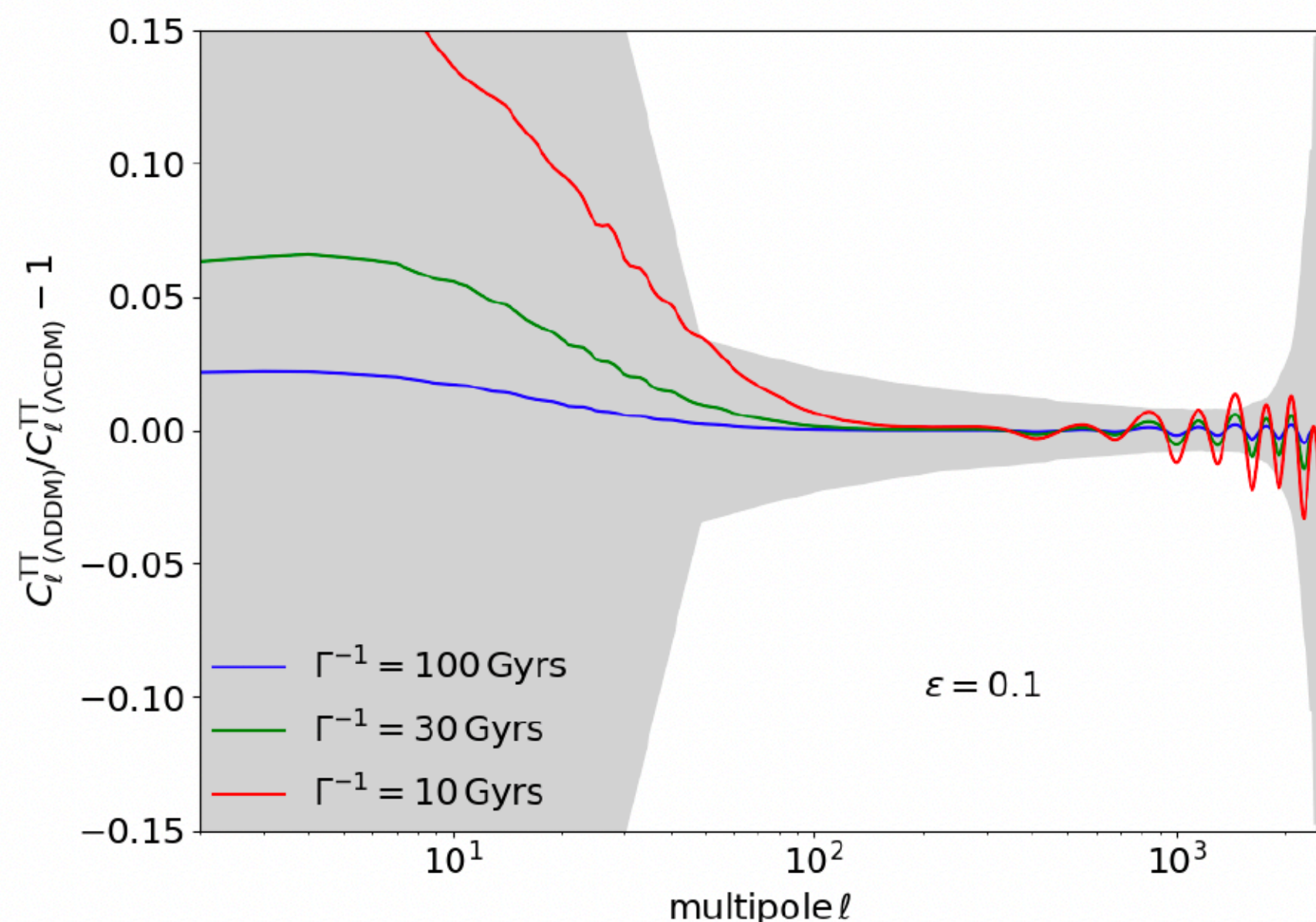


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# Resolution to the $S_8$ tension

The level of detection depends on the level of tension with  $\Lambda$ CDM

