Constraining decaying dark matter with the effective field theory of large-scale structures

Guillermo Franco Abellán

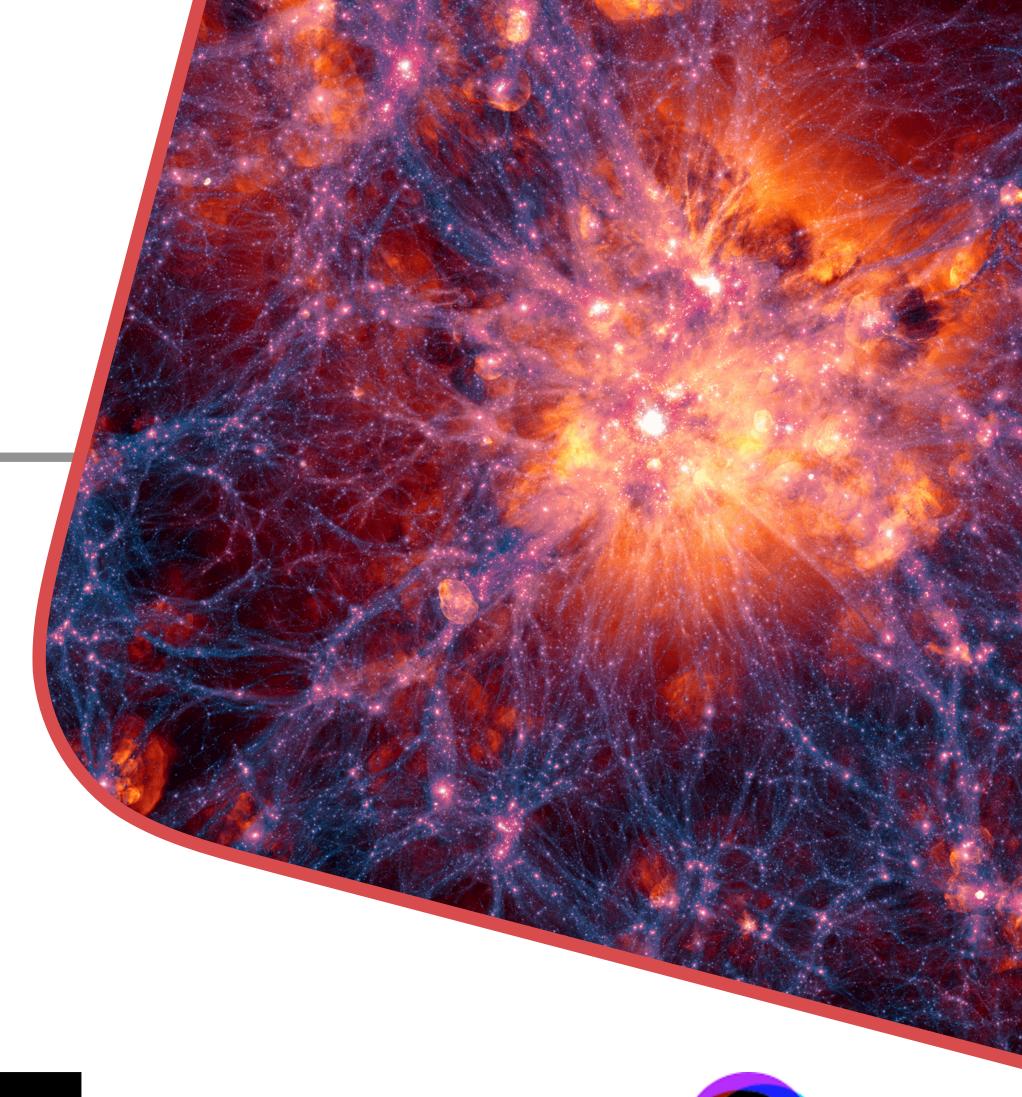
Based on:

PRD 106 (2022), [arXiv:2203.07440]

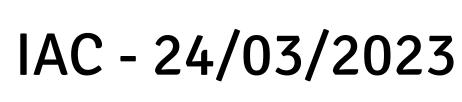
PRD 104 (2021), [arXiv:2102.1249]

PRD 105 (2022), [arXiv:2008.09615]

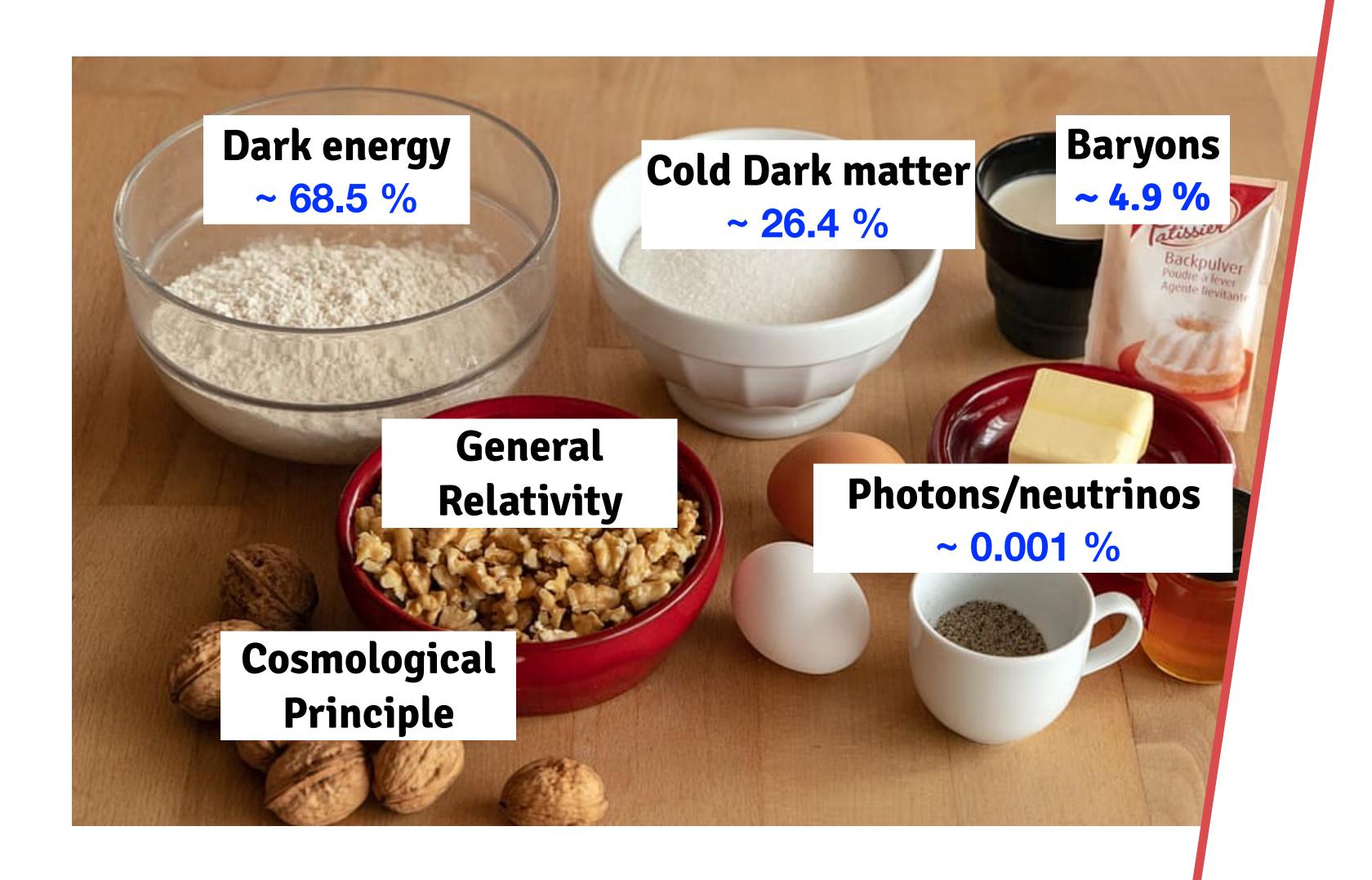
with T. Simon, V. Poulin, P. Du, Y. Tsai, R. Murgia, J. Lavalle



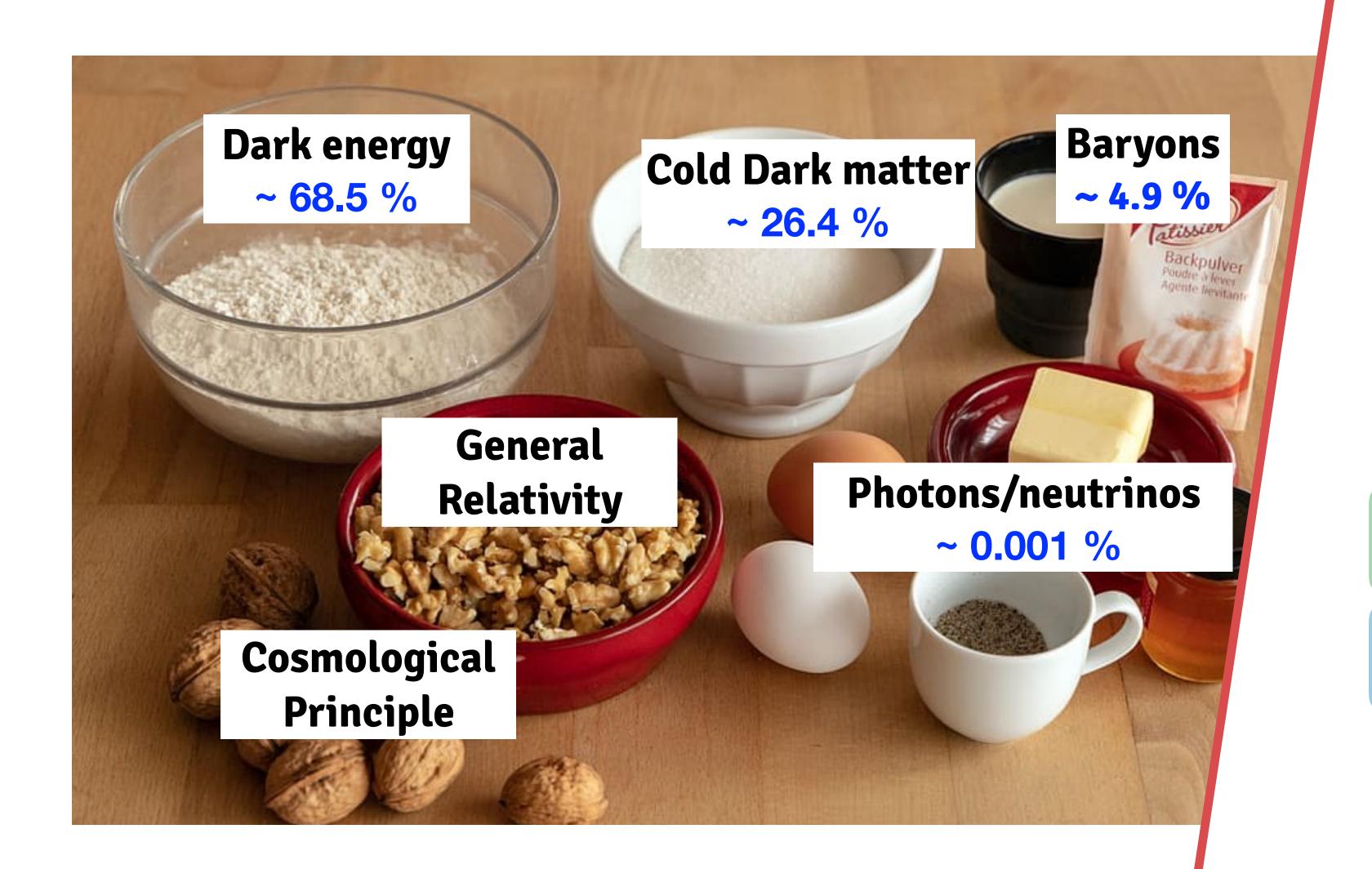
GRAPPA *



Concordance **\CDM** model of cosmology:



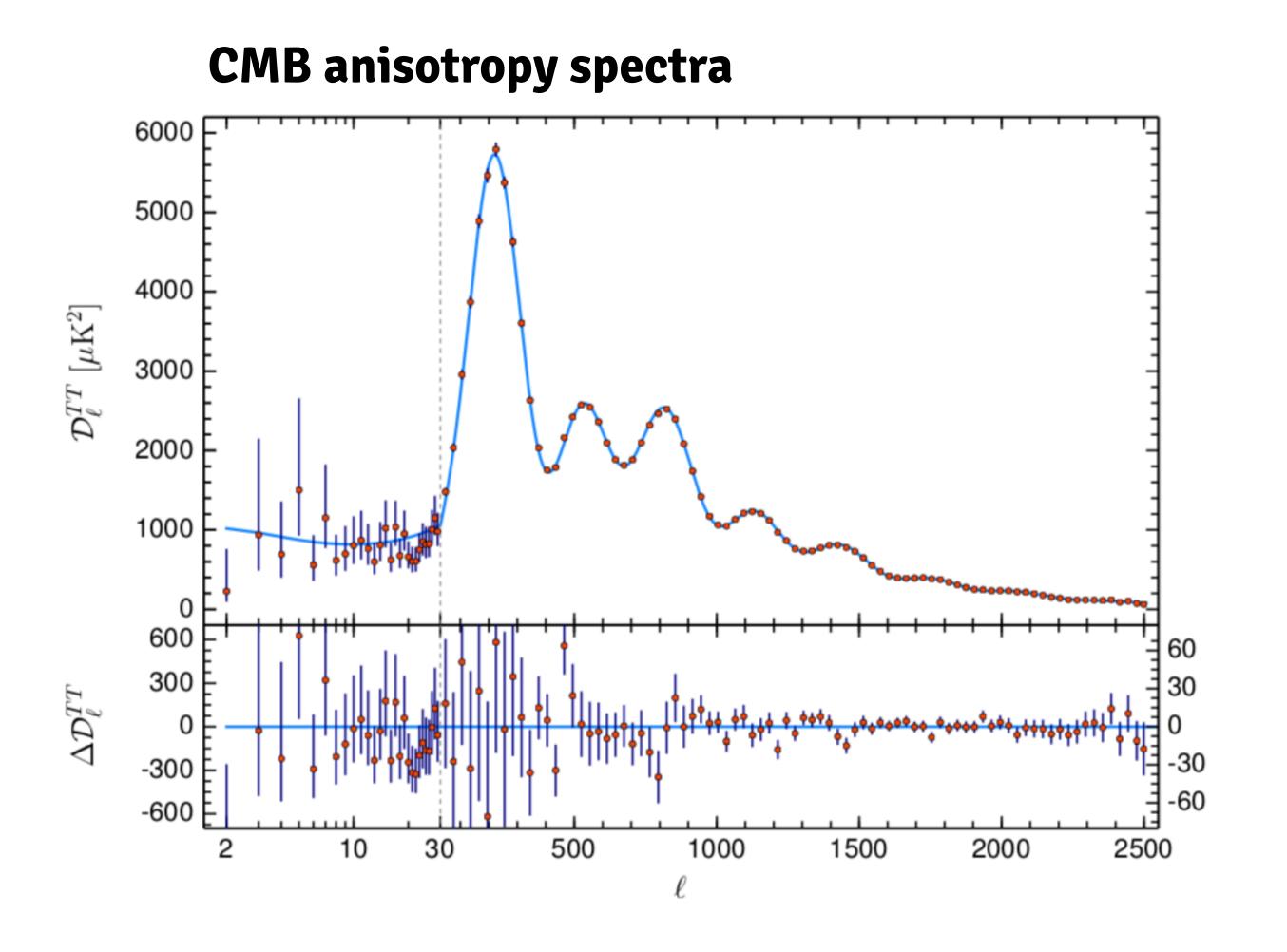
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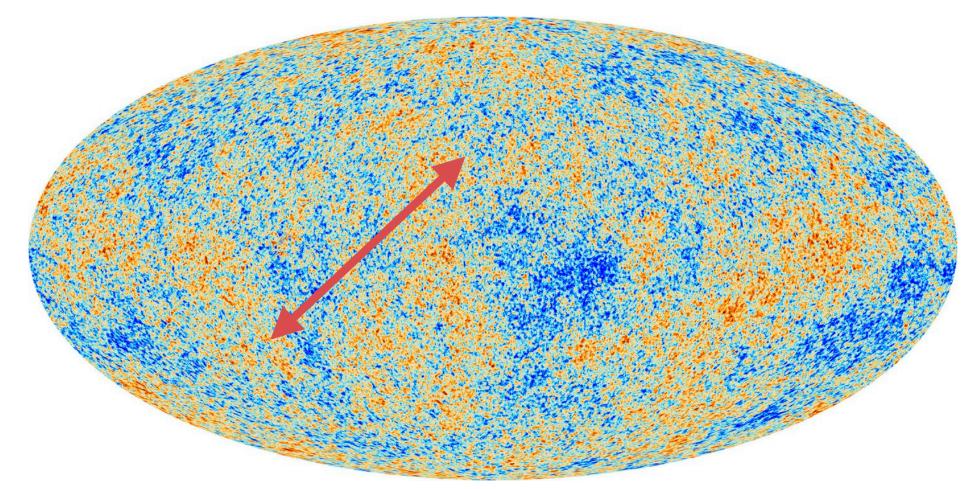


Only 6 free parameters:

$$\omega_{\rm c}$$
 $\omega_{\rm b}$ H_0
 A_s n_s $\tau_{\rm reio}$

Excellent agreement with a wide variety of observations

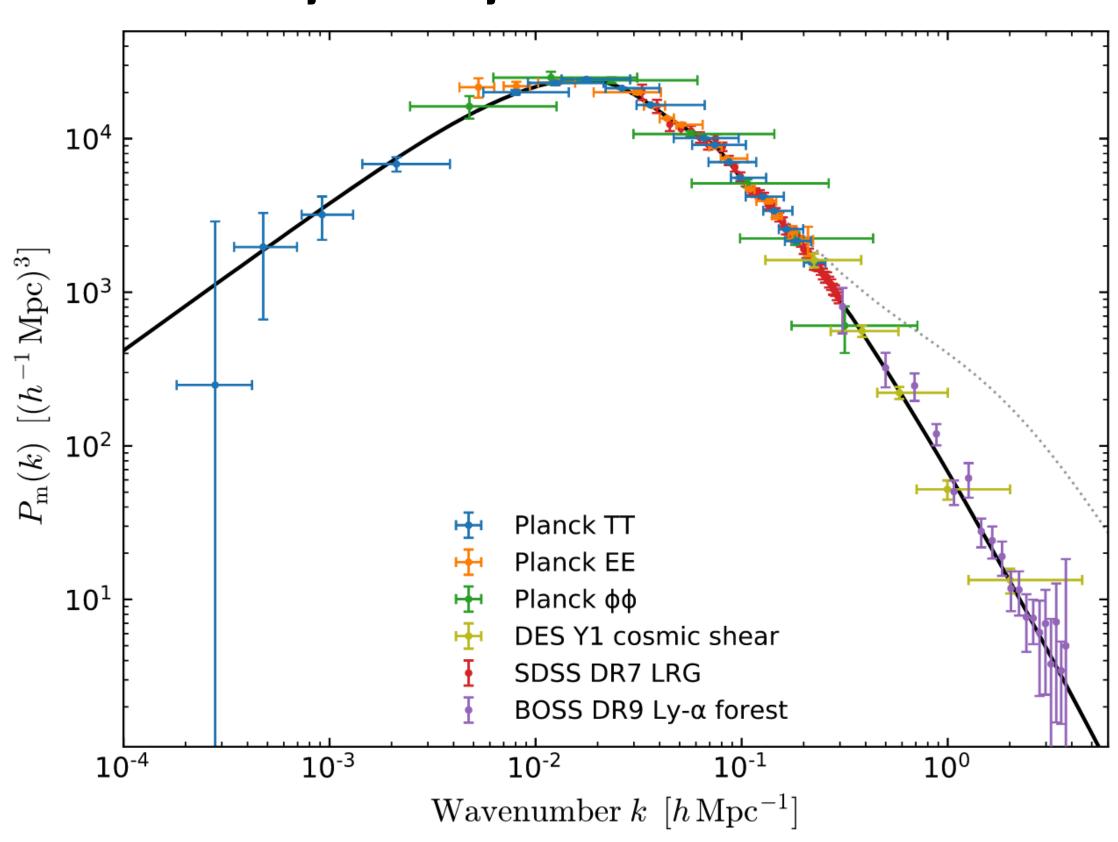


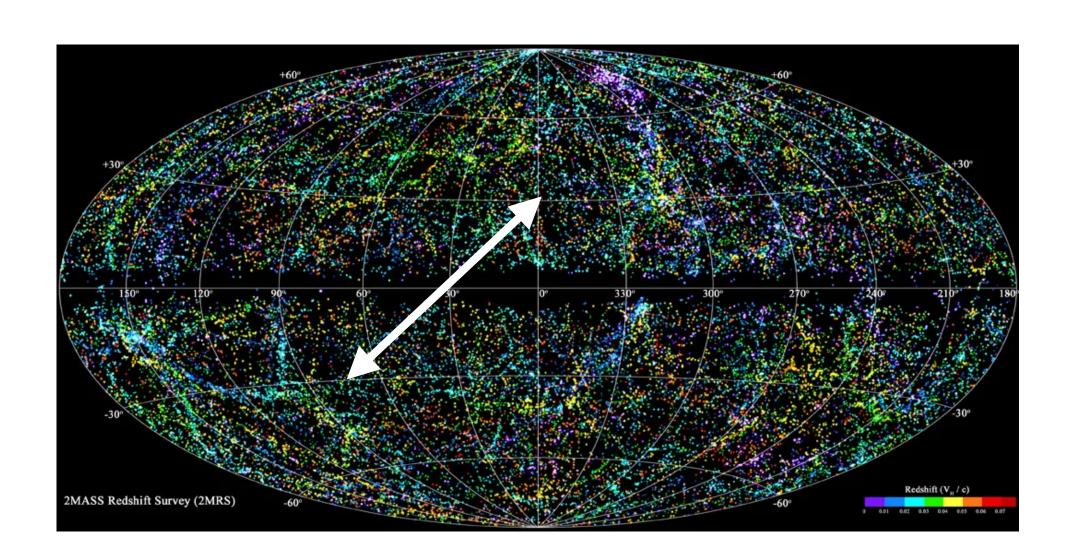


$$C_{\ell}^{\mathrm{TT}} \sim \left\langle \left(\frac{\delta T}{T} \right)^2 \right\rangle \qquad \ell \sim \frac{1}{\theta}$$

Excellent agreement with a wide variety of observations

Matter power spectrum

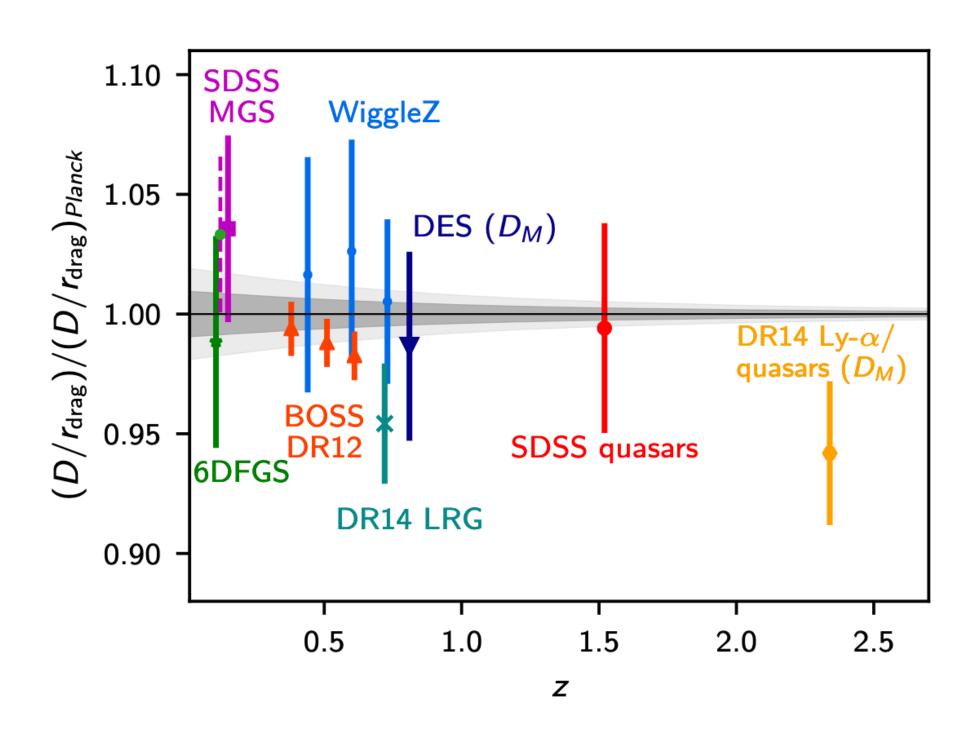




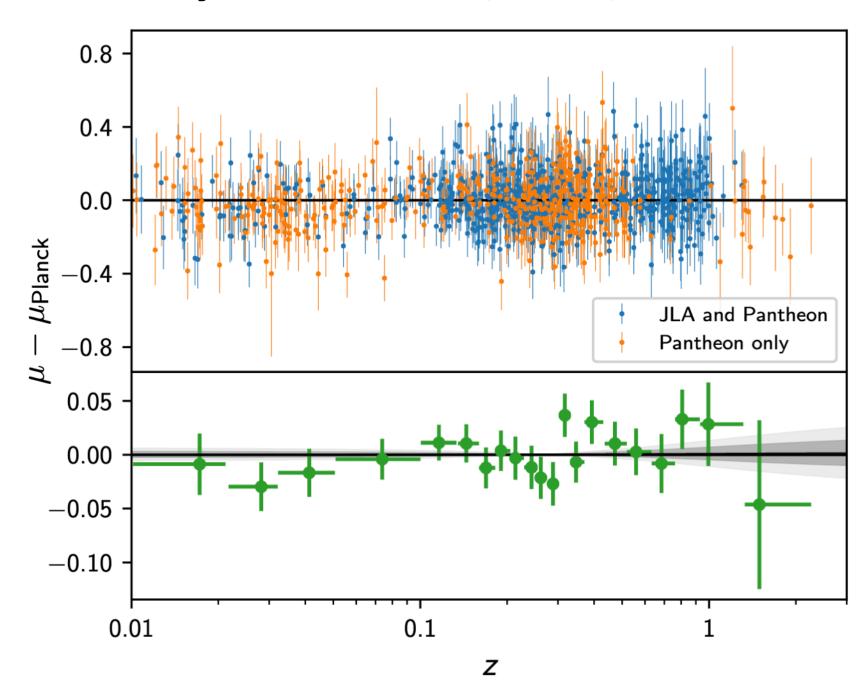
$$P(k) \sim \left\langle \left(\frac{\delta \rho}{\rho}\right)^2 \right\rangle \qquad k = \frac{2\pi}{\lambda}$$

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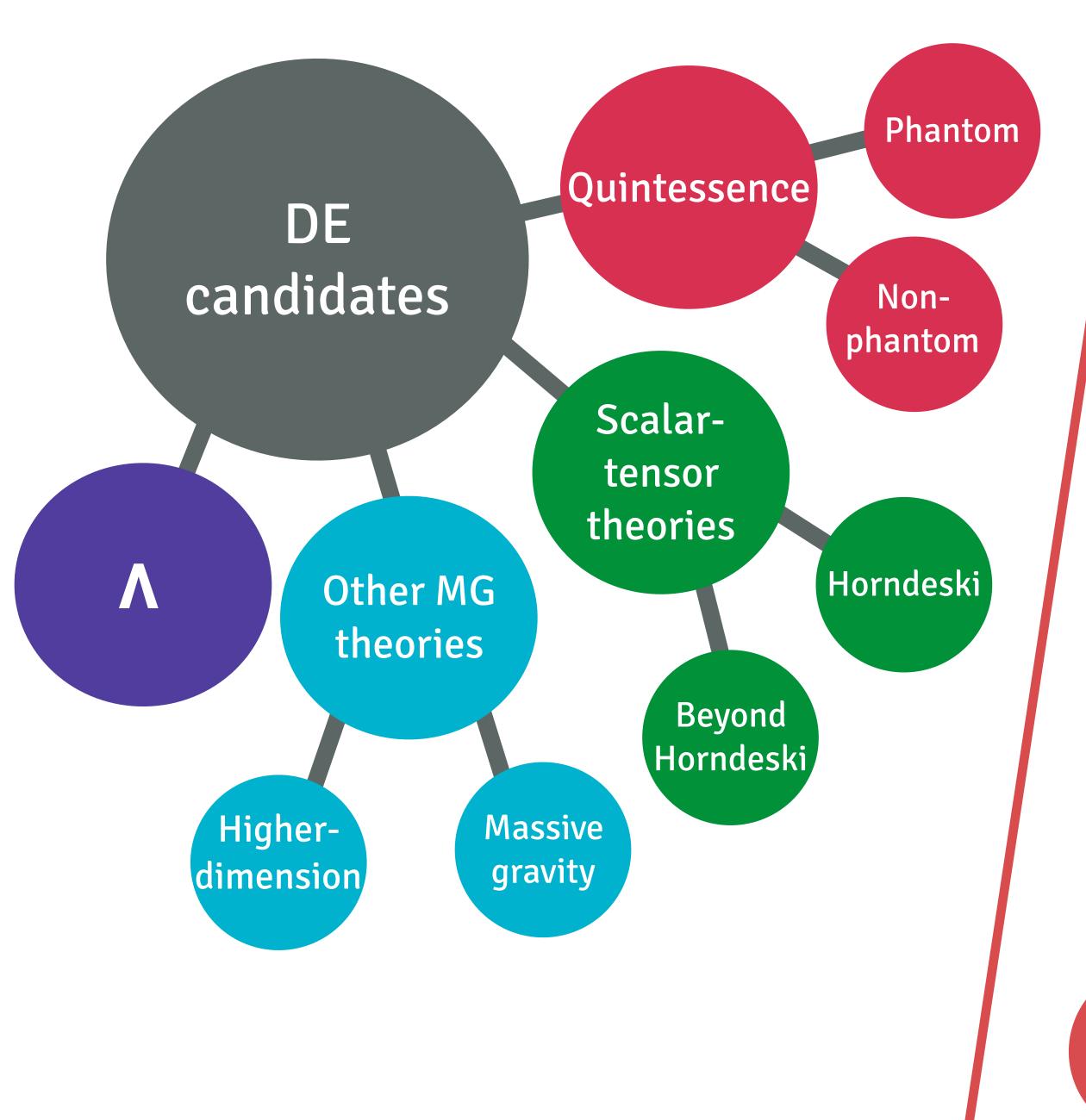
Baryon acoustic oscillations (BAO)



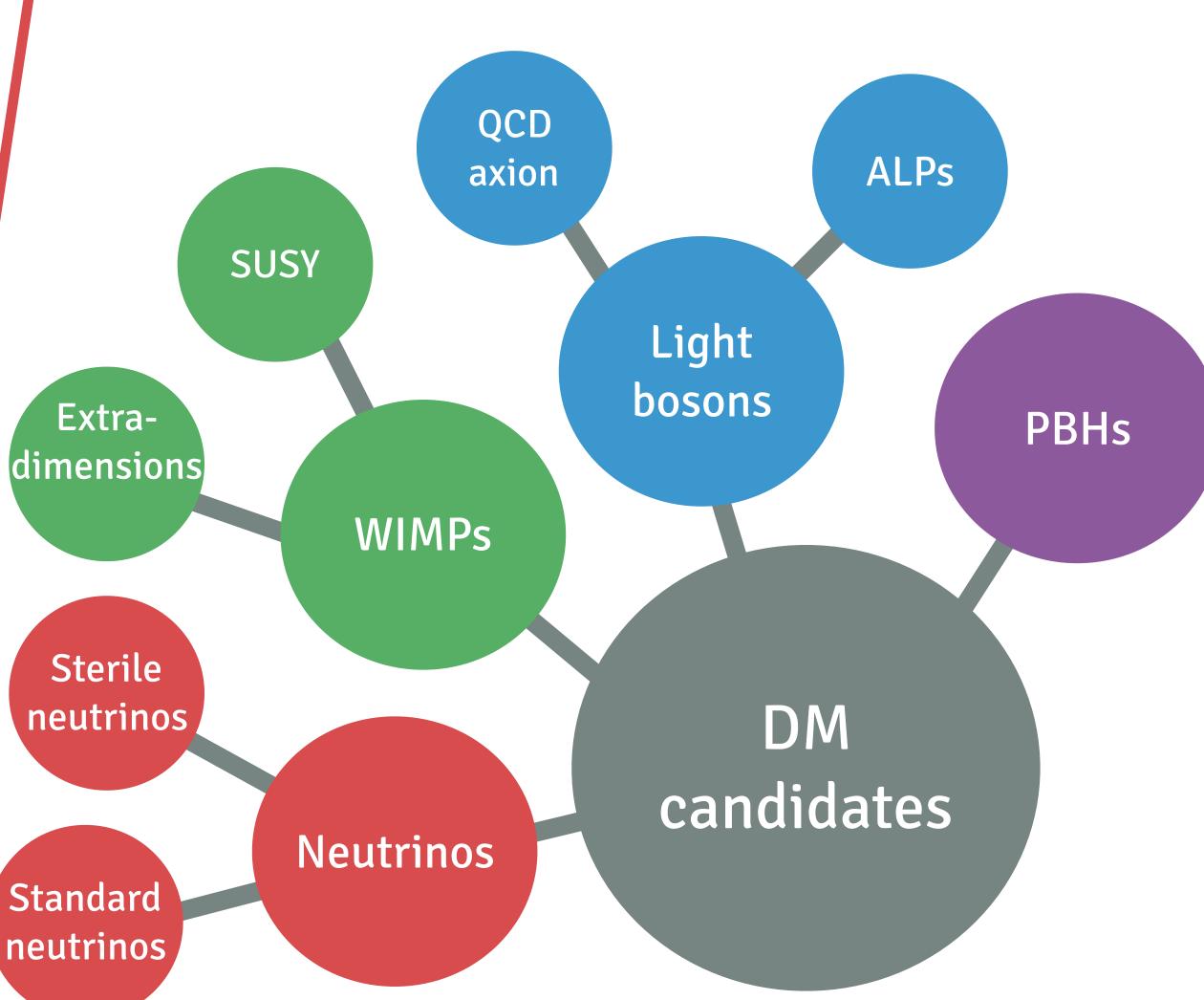
Supernovae la (SNIa)



$$D(z) \propto \int_0^z \frac{dz'}{H(z')}$$



However, the nature of the dark sector remains a mystery



In addition, several discrepancies have emerged in recent years

 H_0 tension (5 σ)

[Riess+ 21] [Planck 18]

 S_8 tension (2-3 σ)

[KiDS 20] [DES 21] [Planck 18]

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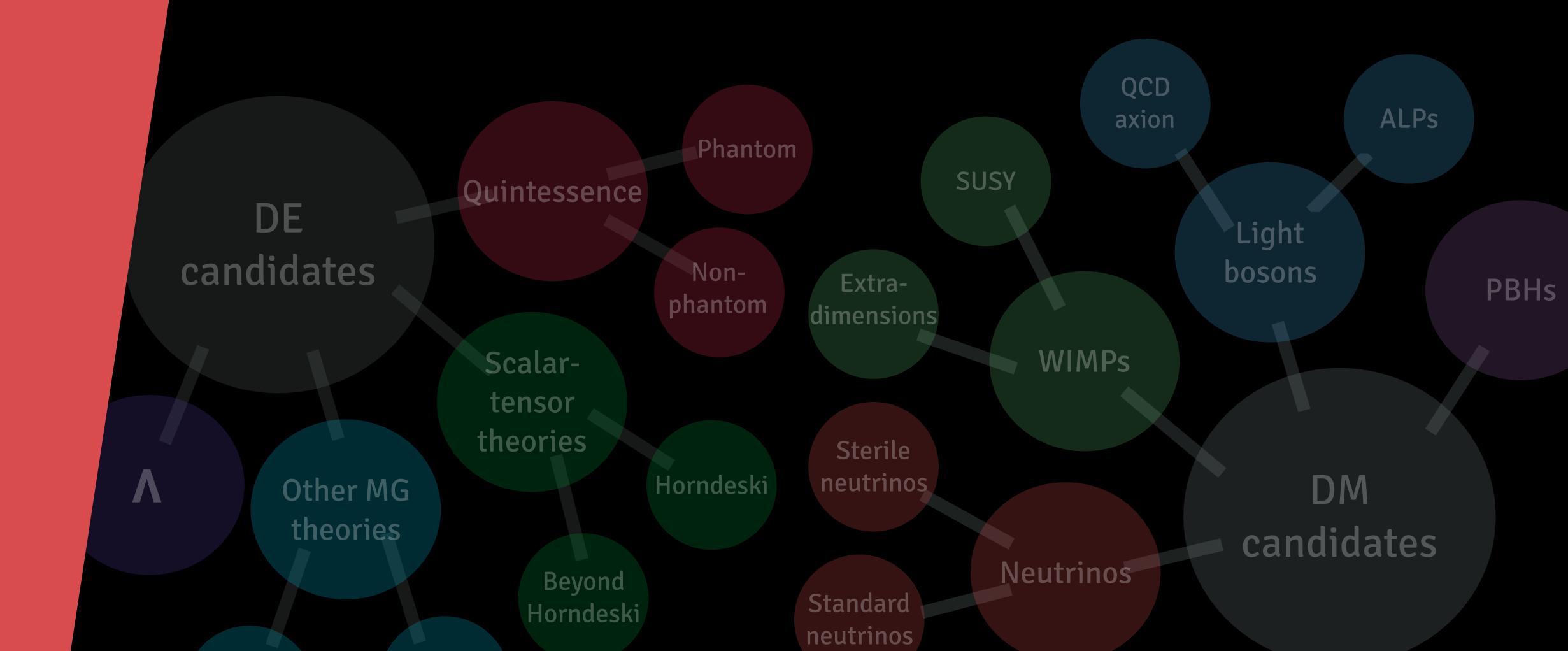
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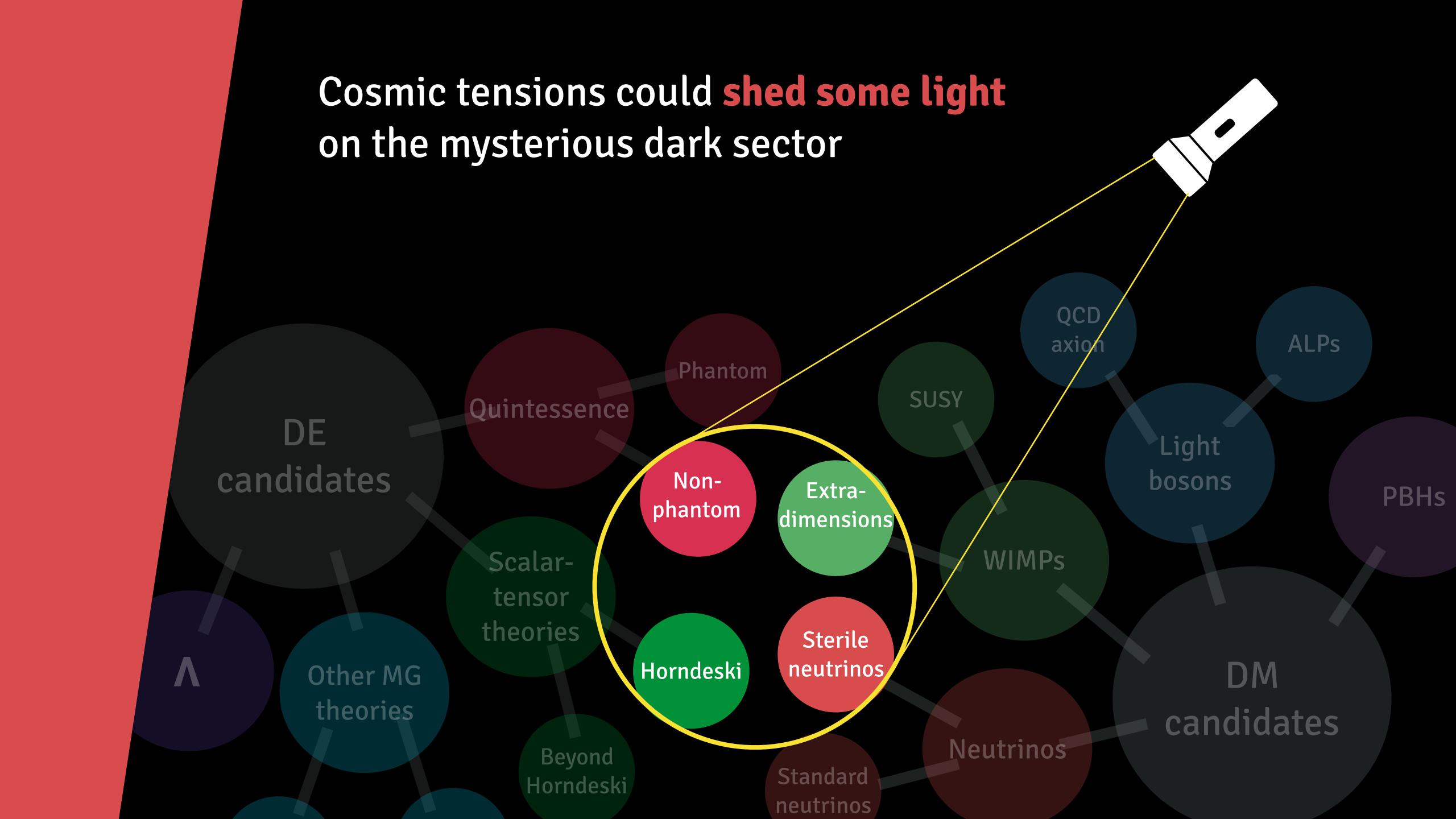
[KiDS 20] [DES 21] [Planck 18]

Systematics?

New physics?

Cosmic tensions could shed some light on the mysterious dark sector





Part I:

DDM AND THE
S₈ TENSION

Part II:

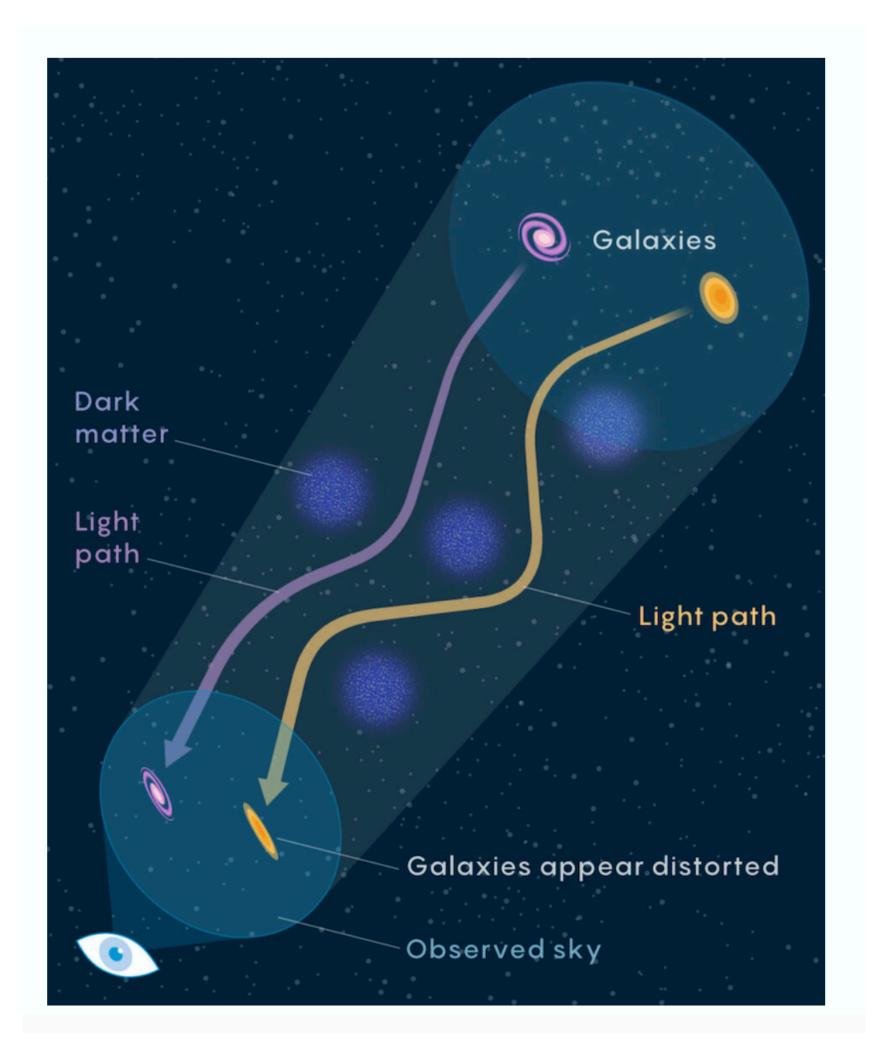
INTRODUCTION
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Part III:

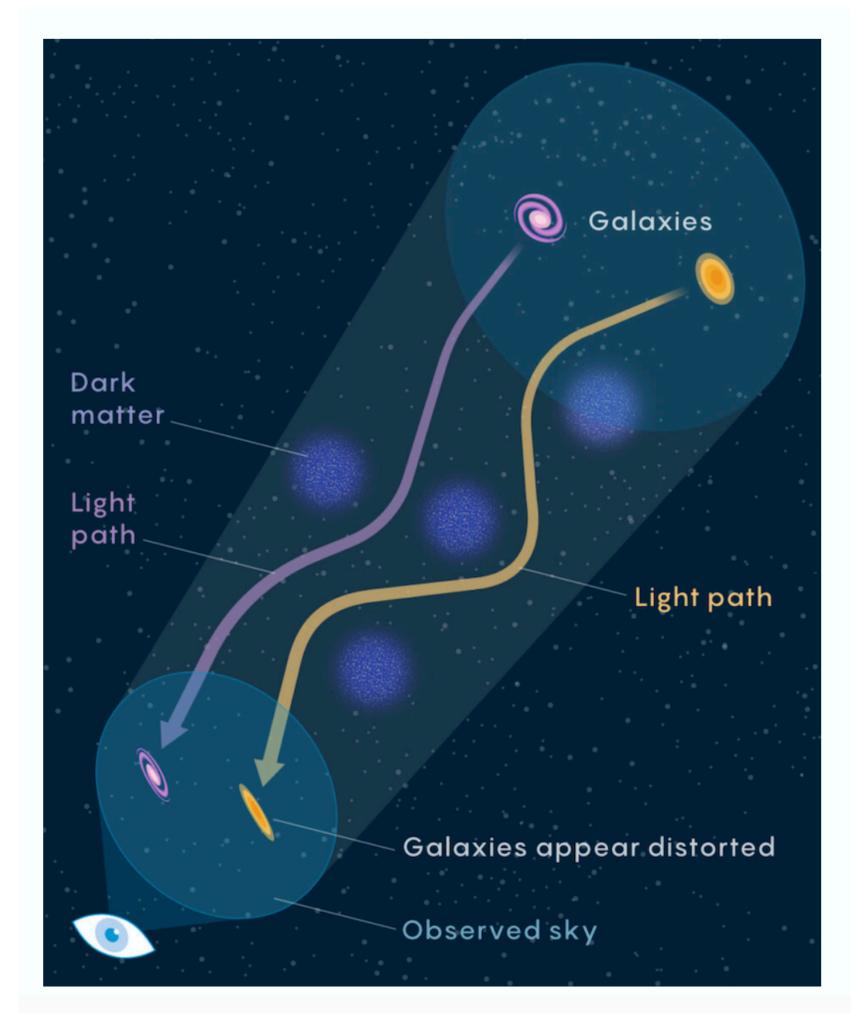
CONSTRAINTS ON DDM FROM EFTofLSS

Part I:

DDM AND THE S₈ TENSION

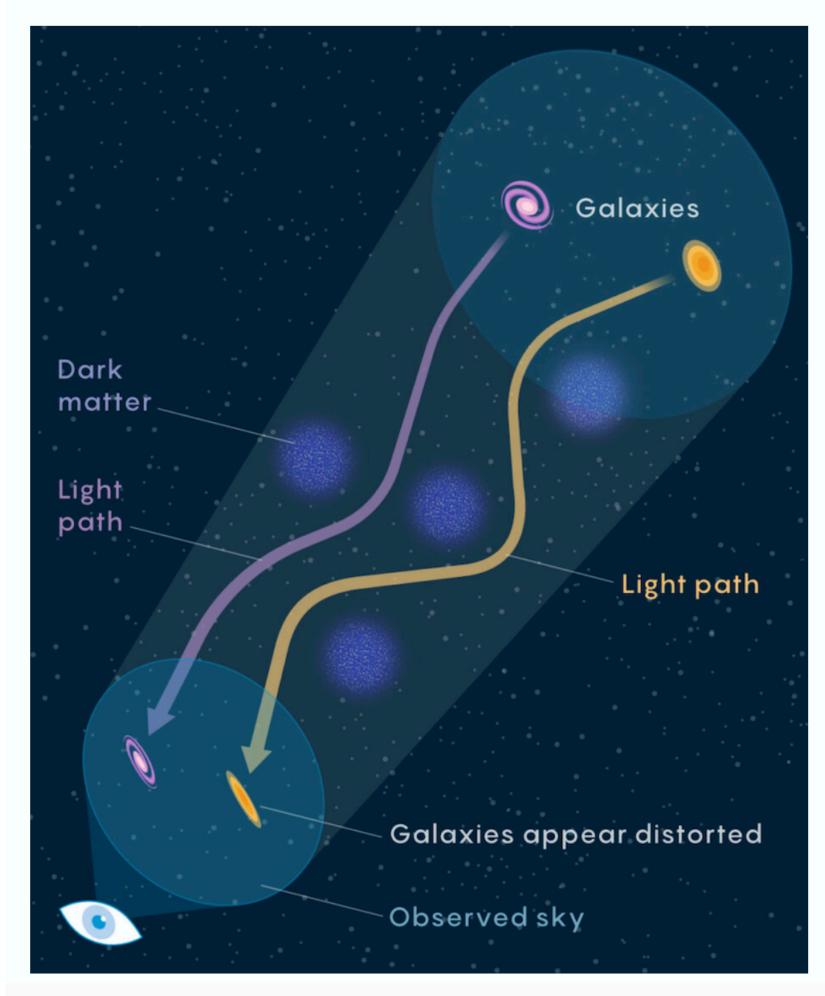


$$S_8 = \sigma_8 \sqrt{\Omega_m/0.3}$$



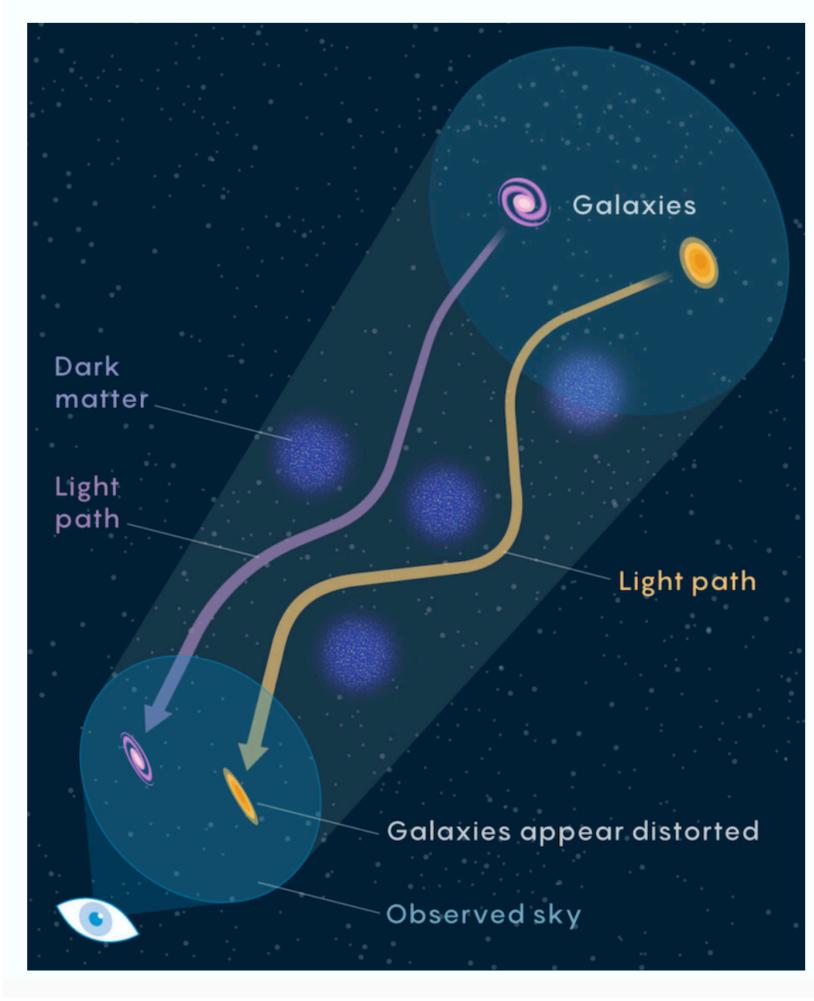
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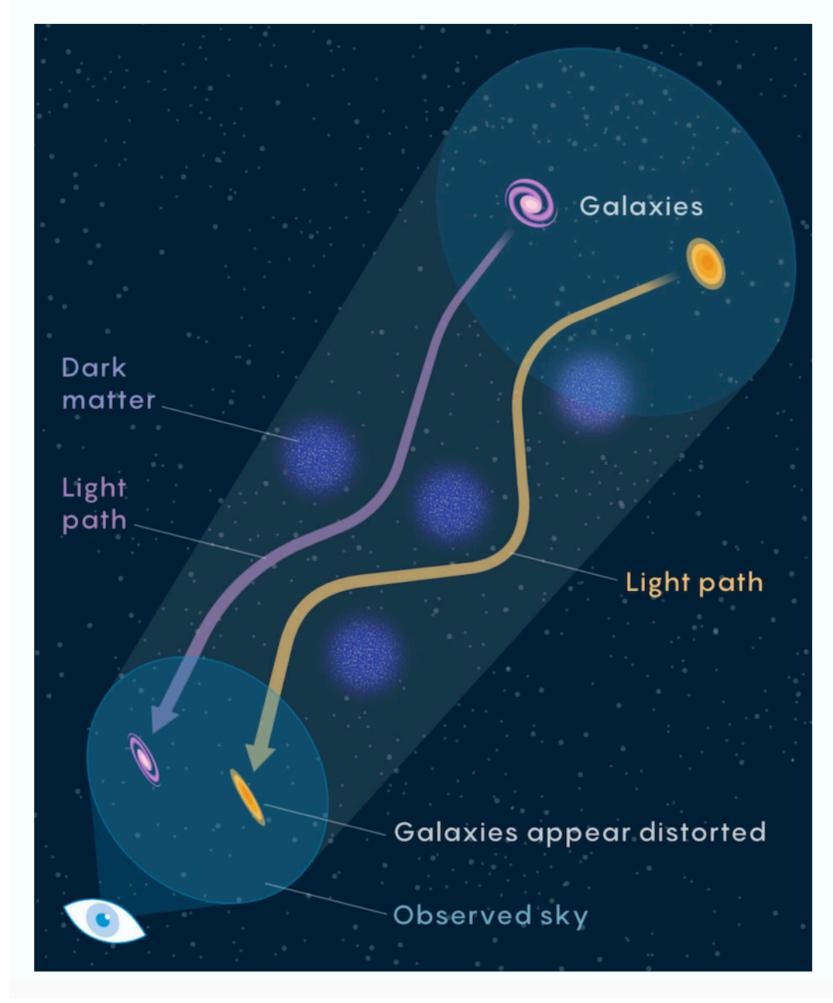
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 \sim 3 σ tension



KiDS galaxy weak-lensing $S_8 = 0.766 \pm 0.016$ Planck CMB $S_8 = 0.830 \pm 0.013$



Other weak-lensing surveys and LSS probes see also tension at 2-3σ level

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Ω_m should be left unchanged (well constrained by SNIa & galaxy clustering)

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$$\sigma_8^2 = \int P_m(k, z = 0) W_8^2(k) d\ln k$$

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Suppress matter power at scales $k \sim 0.1 - 1 \ h/{\rm Mpc}$

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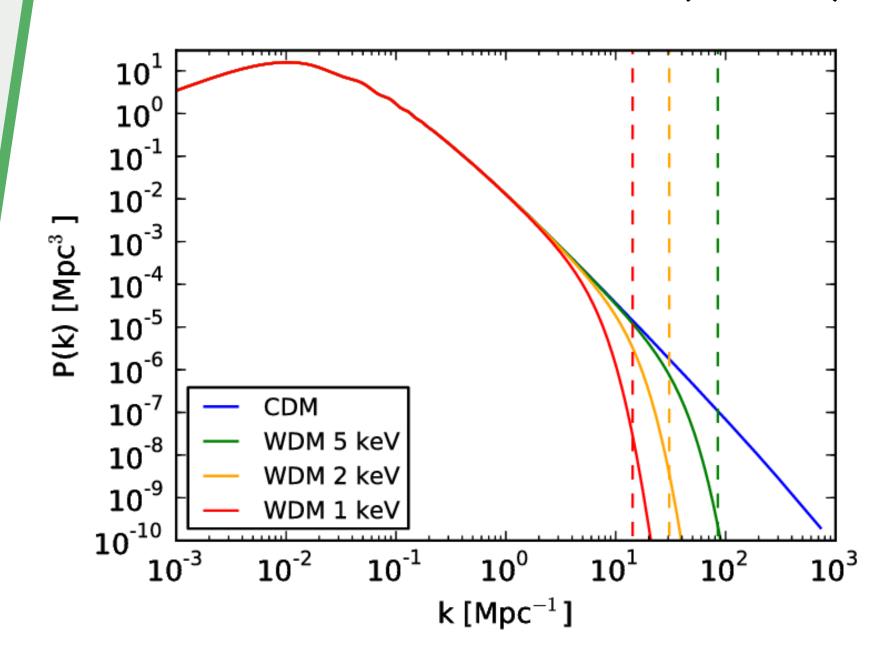
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Ex: Warm dark matter (WDM)



Very constrained by Ly-a! [Iršič+ 17]

Decaying Dark Matter (DDM)

Well motivated theoretically (ex: R-parity violation, hidden U(1) symmetries, superWIMPs,...)

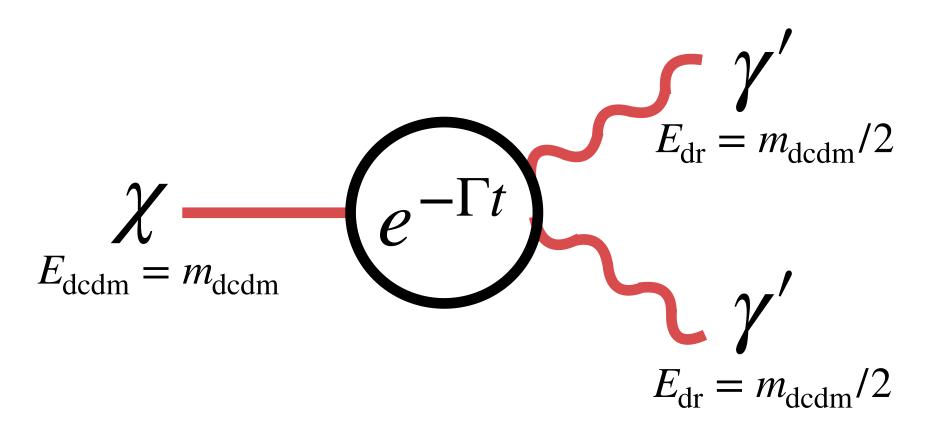
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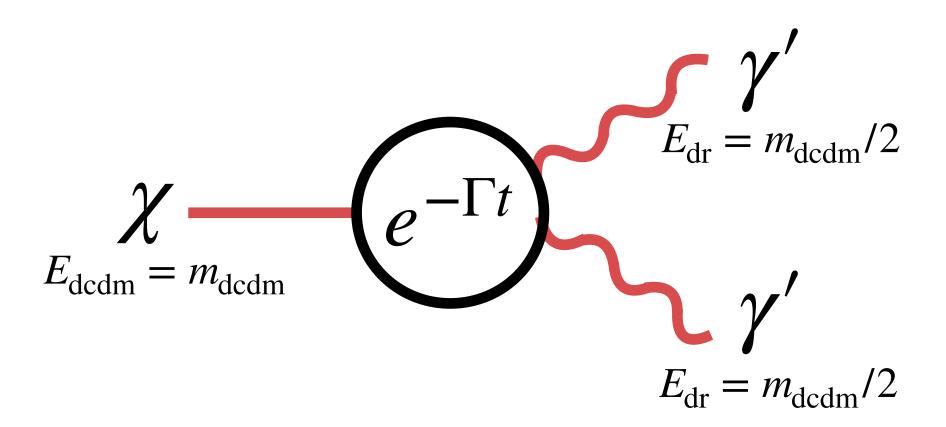
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- Decay products?
 - To Standard Model particles Model-dependent, strongly constrained $\Gamma^{-1} \gtrsim 10^7 10^{10}~t_U$ [Blanco+ 18]
 - To "invisible" particles, i.e., dark radiation (DR) Model-independent, much less constrained $\Gamma^{-1}\gtrsim 10~t_U$ [Nygaard+ 20]

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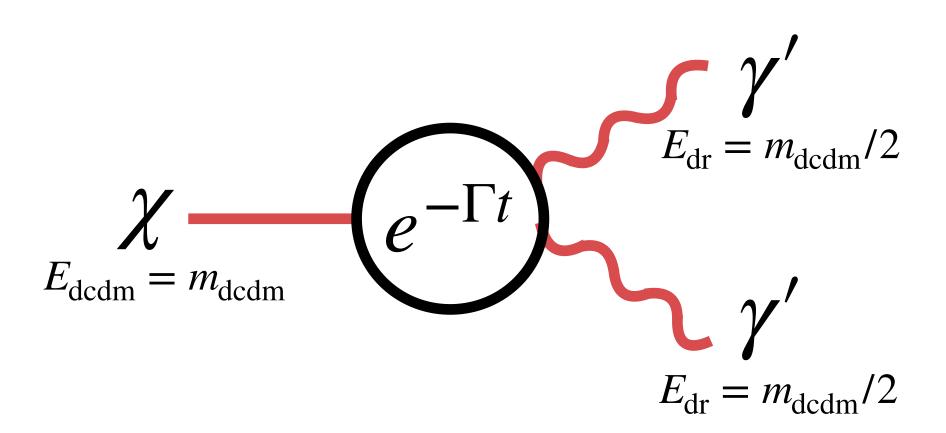






Decay rate Γ DCDM fraction $f_{
m dcdm}$

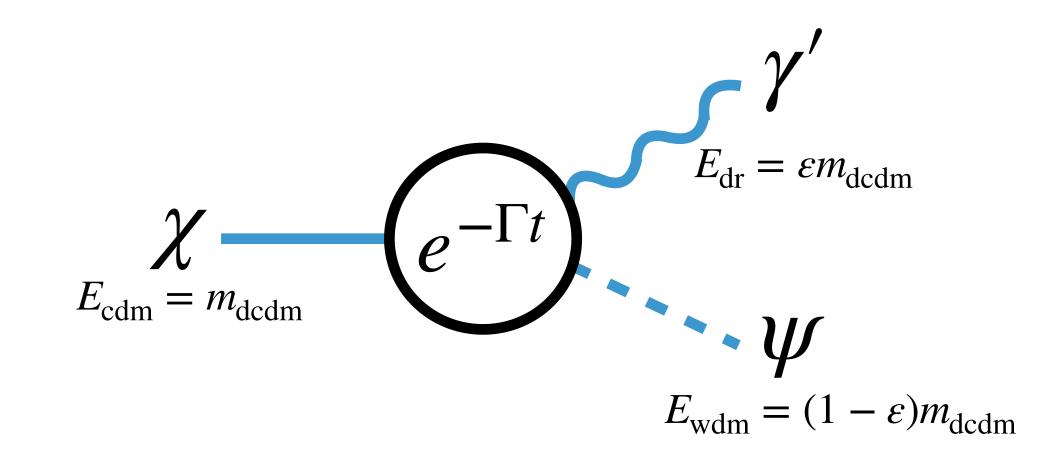
$$f_{\text{dcdm}} \equiv \frac{\Omega_{\text{dcdm}}^{\text{ini}}}{\Omega_{\text{dcdm}}^{\text{ini}} + \Omega_{\text{cdm}}} \in [0, 1]$$

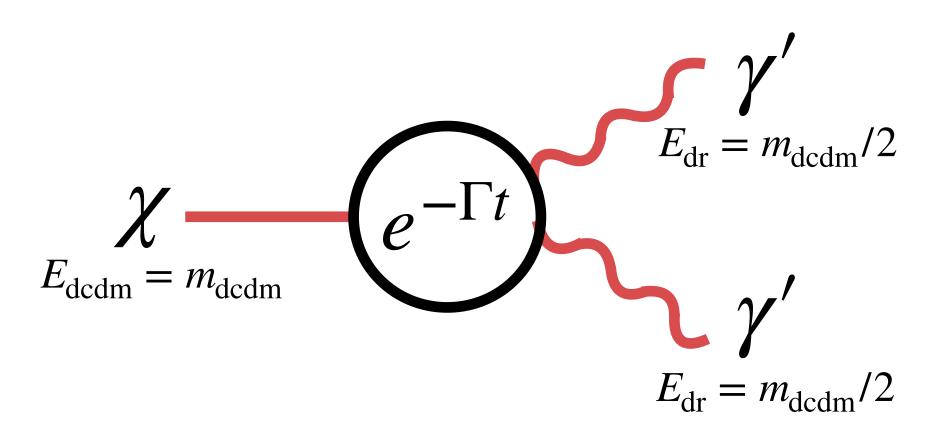


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CDM decaying to DR + WDM

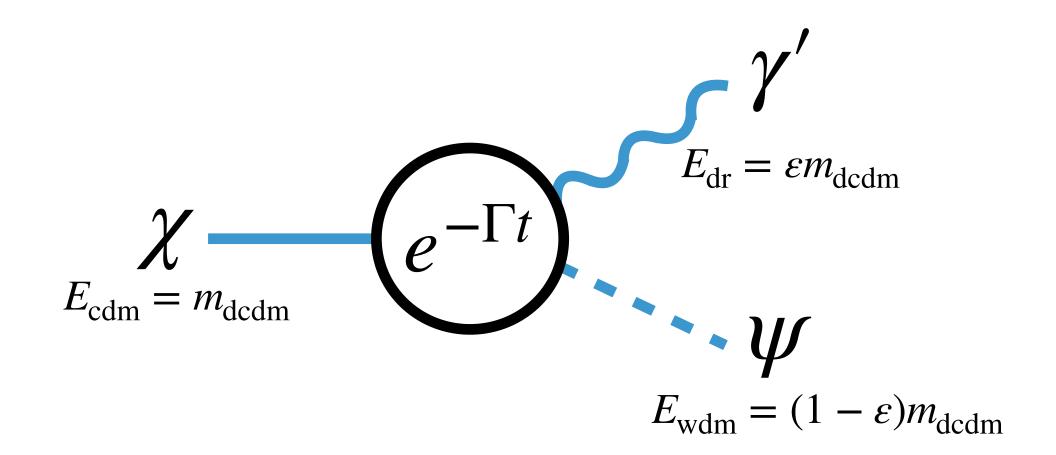




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CDM decaying to DR + WDM



Decay rate Γ DR energy fraction \mathcal{E}

$$\varepsilon = \frac{1}{2} \left(1 - \frac{m_{\text{wdm}}^2}{m_{\text{dcdm}}^2} \right) \in [0, 1/2]$$

$$\dot{\rho}_{dcdm} + 3H\rho_{dcdm} = -\Gamma \rho_{dcdm}$$

$$\dot{\rho}_{dr} + 4H\rho_{dr} = +\Gamma \rho_{dcdm}$$

+ linear perturbed eqs. for DCDM and DR

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What's the impact on cosmological observables?

Expansion history H(z)

Impacted by CDM \rightarrow DR, not much by CDM \rightarrow DR+WDM ($\rho_{wdm} \sim \rho_{cdm} \sim a^{-3}$)

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CMB anisotropy spectra C_ℓ^{TT,EE}

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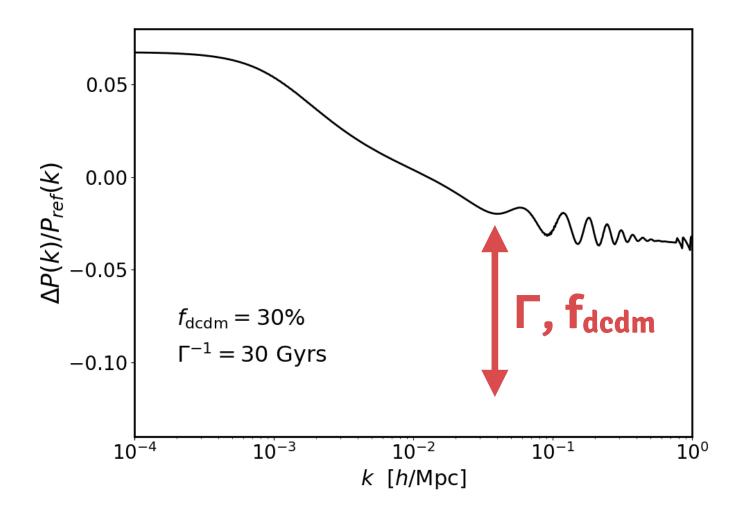
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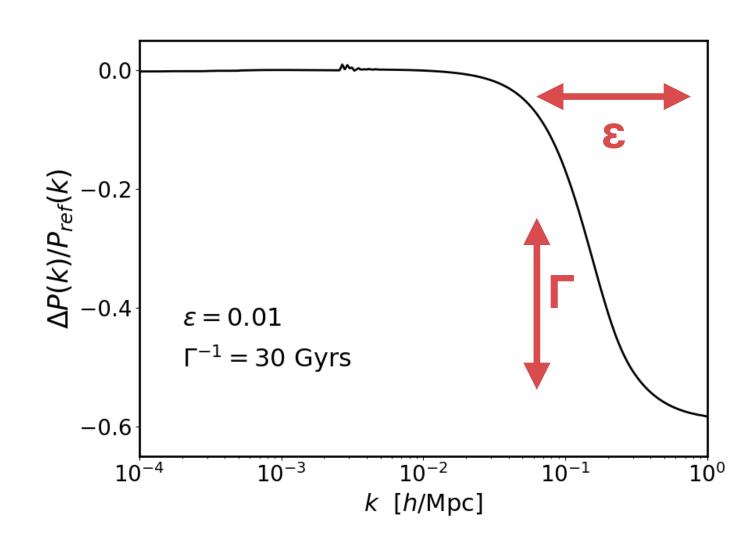
Impact even for late decays, both models affect LISW and CMB lensing

Linear matter power spectrum P_m(k)

CDM → DR shifts position of the peak



CDM \rightarrow DR+WDM suppresses power at k > k_{fs}



CDM \rightarrow DR has been shown to fail at explaining the S₈ tension [Poulin+16] [Schoneberg+21]

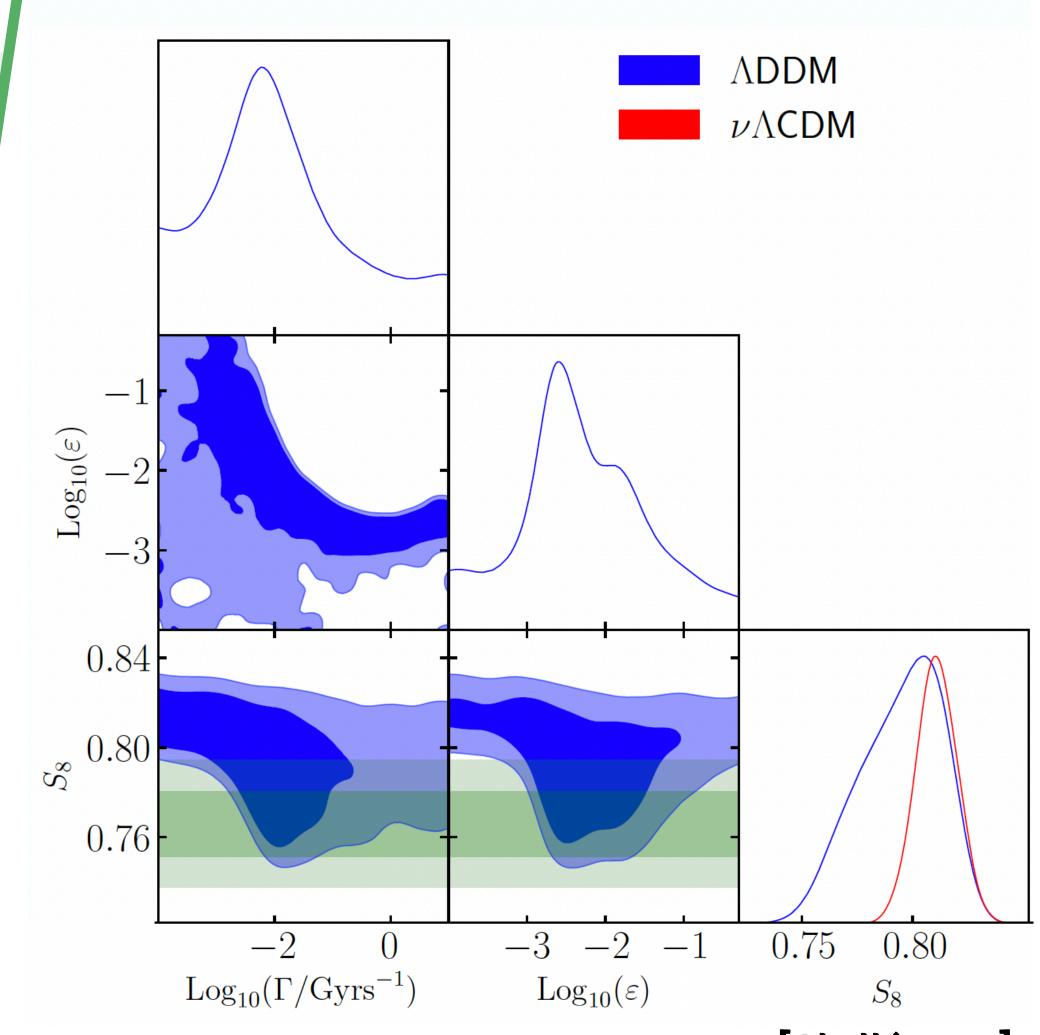
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What about the CDM → DR+WDM model?

CDM → DR+WDM solves the S₈ tension

Reconstructed S₈ values are in excellent agreement with WL data

Planck18 + BAO + SNIa +S₈ prior (KiDS+BOSS+2dfLenS):

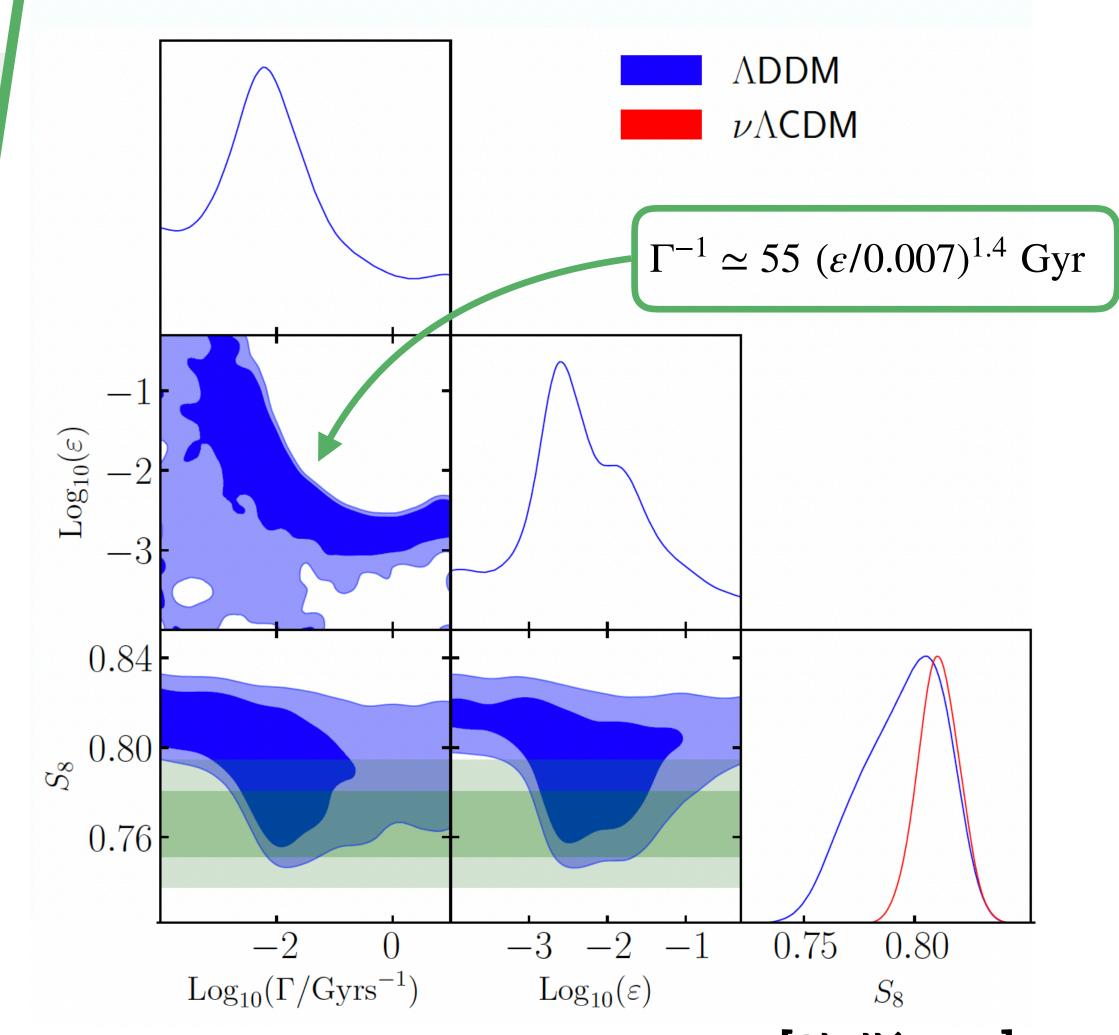


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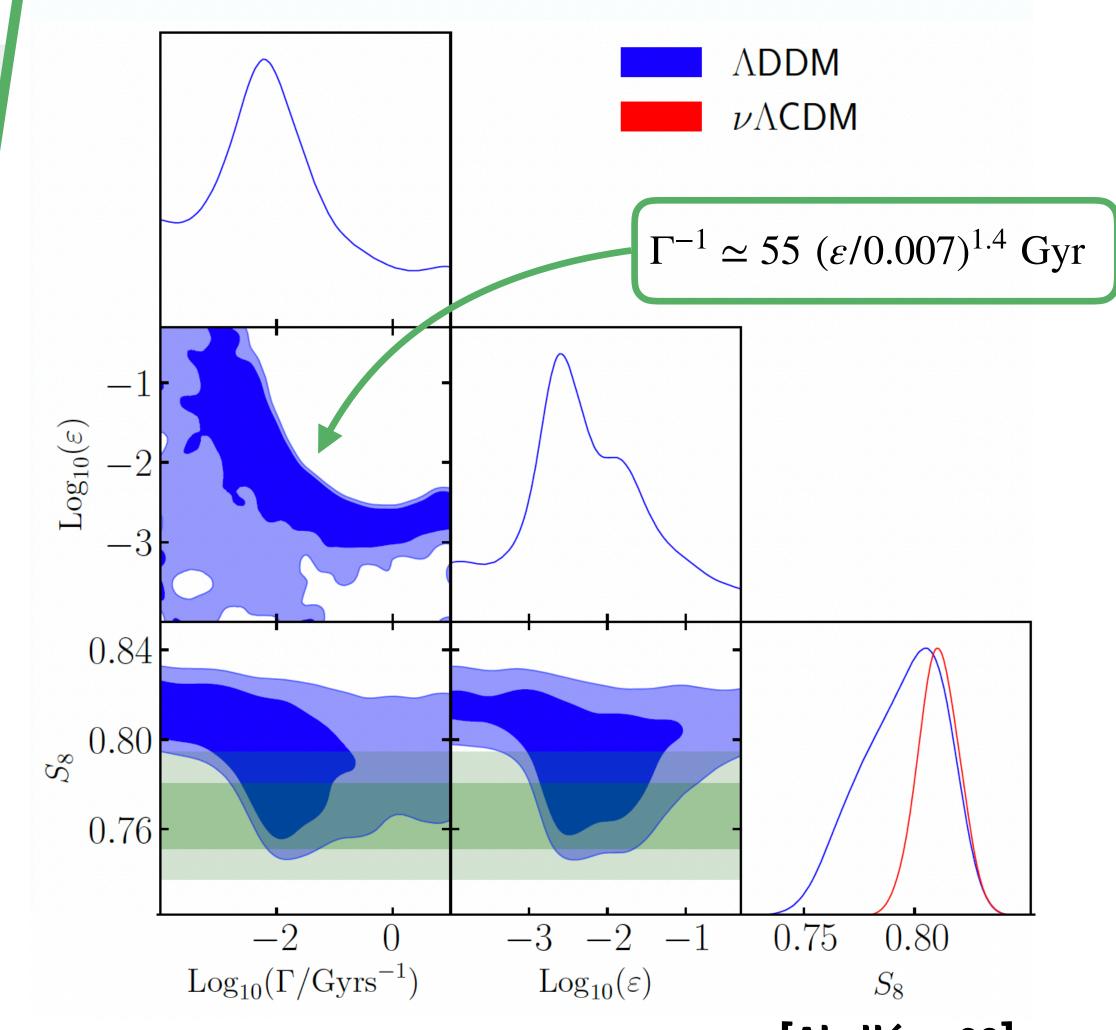
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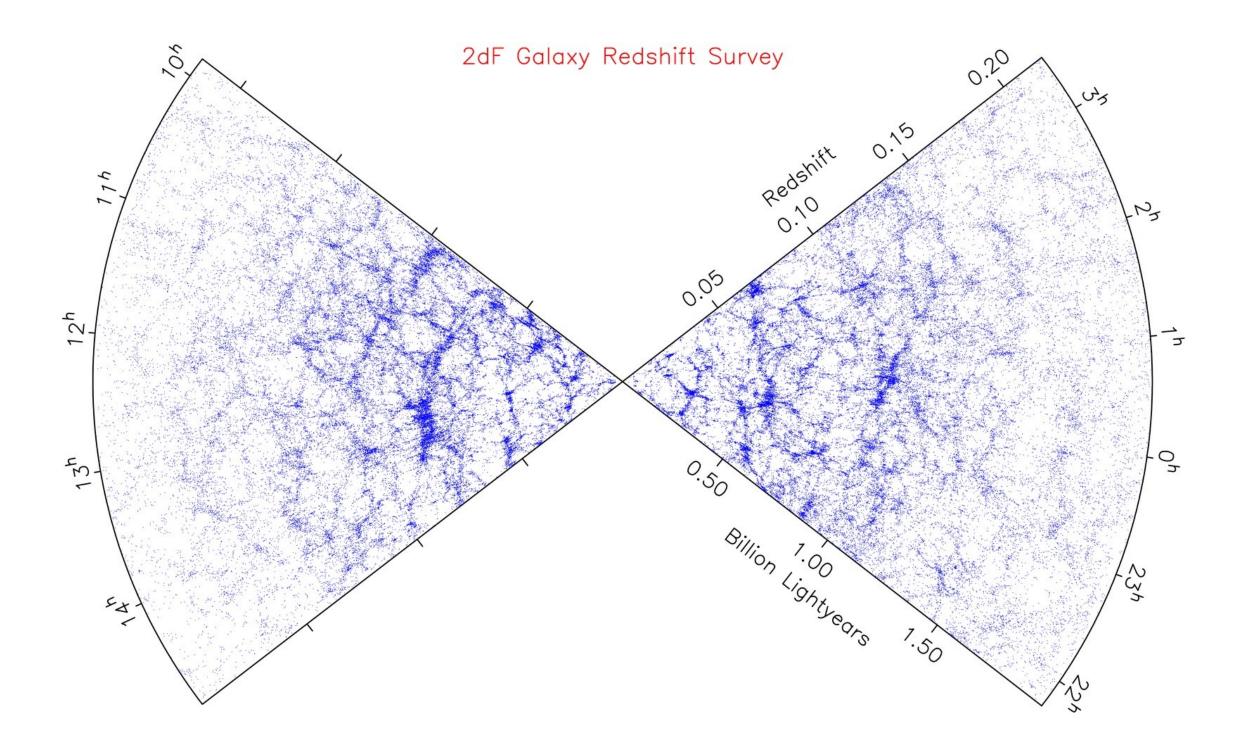
Can we test this with other probes?

Planck18 + BAO + SNIa +S₈ prior (KiDS+BOSS+2dfLenS):

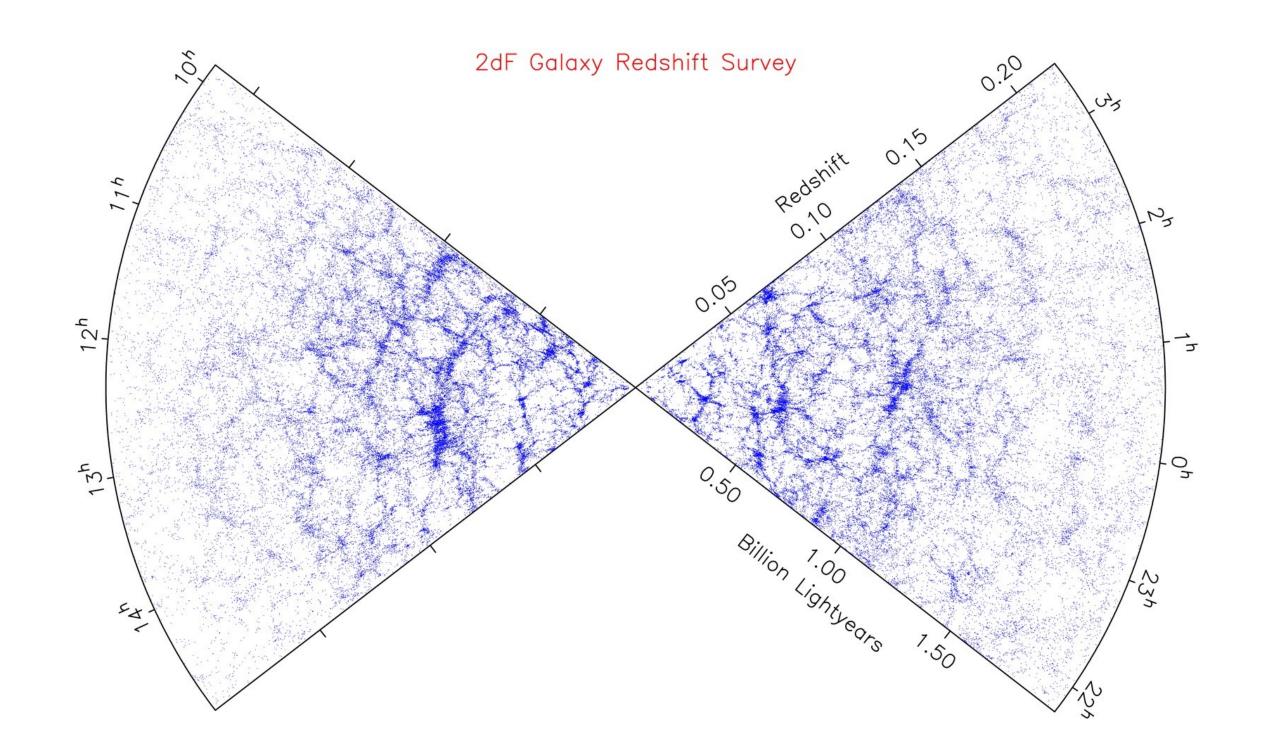


Part II:
INTRODUCTION
TO THE EFTofLSS

We would like to analyze the clustering of galaxies on the largest scales



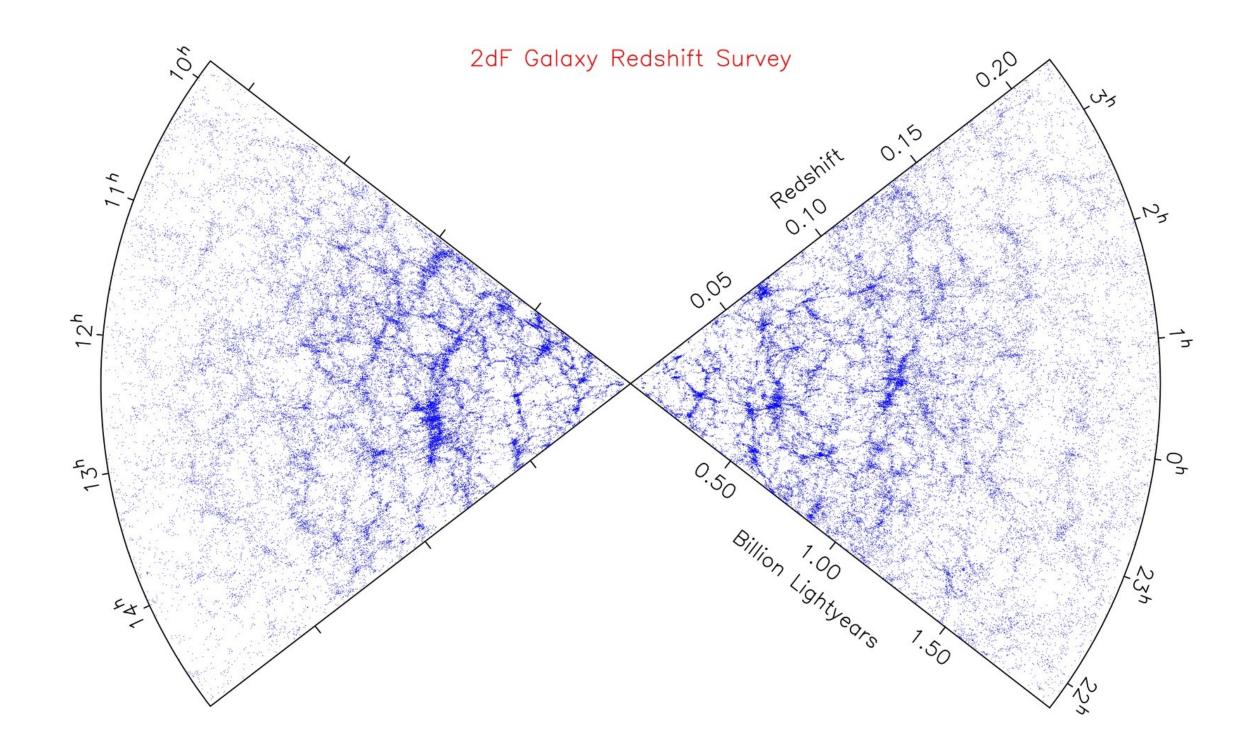
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Unfortunately, we cannot rely on linear theory anymore

$$\begin{split} \partial_{\tau}\delta + \partial_{i}[(1+\delta)v^{i}] &= 0 \\ \partial_{\tau}v^{i} + aHv^{i} + \partial^{i}\Phi + v^{j}\partial_{j}v^{i} &= -\frac{1}{a\rho}\partial_{j}\tau^{ij} \\ \Delta\Phi &= \frac{3}{2}a^{2}H^{2}\Omega_{m}\delta \end{split}$$

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800 600 growth 400 400 distance 200 0.05 0.20 0.10 0.15 Wavenumber k [h/Mpc]

Until now, galaxy clustering data was analyzed by focusing just on compressed features

$f\sigma_8/BAO$ tests

$$\frac{f\sigma_8(z)}{(f\sigma_8(z))_{\mathrm{fid}}} \frac{H(z)}{H(z)_{\mathrm{fid}}} \frac{D_A(z)}{D_A(z)_{\mathrm{fid}}}$$

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Can we use the full shape information?

N-body simulations









N-body simulations

- Unlimited range of scales
- Good for matter clustering
- Very time-consuming
- Galaxy formation uncertain

Effective Field Theory

- Limited range of scales (based on perturbation theory)
- More insight into data
- Fast & accurate predictions
- Marg. over uncertainties (free parameters that capture small-scale physics)

Two important ingredients

Galaxy bias

$$\delta_g = b_1 \delta + \dots$$

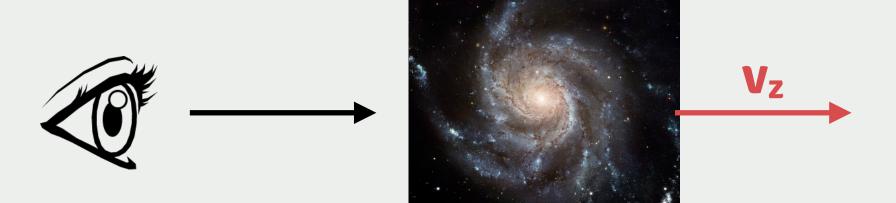
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$$z = Hr + v_z$$



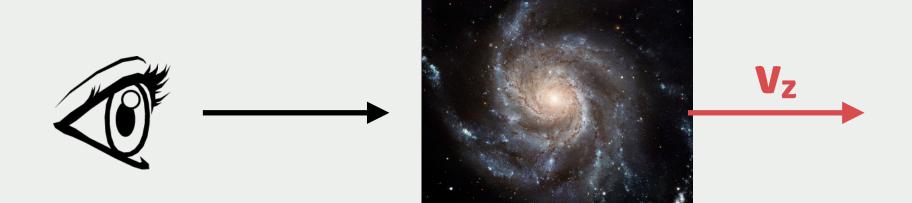
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Galaxy spectra are anisotropic

$$P(\mu, k)$$
 with $\mu \equiv \hat{\mathbf{k}} \cdot \hat{\mathbf{z}}$

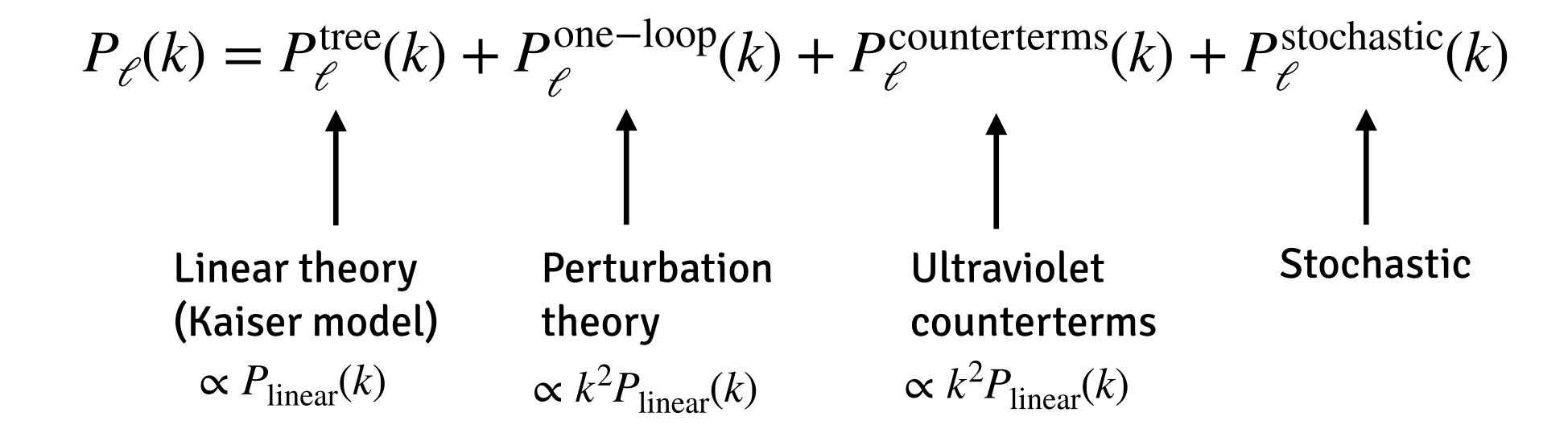
Expand in Legendre multipoles

$$P_{\ell}(k) = \frac{2\ell + 1}{2} \int_{-1}^{1} L_{\ell}(\mu) P(\mu, k) d\mu$$

$$\ell = 0$$
 (monopole), $\ell = 2$ (quadrupole), ...

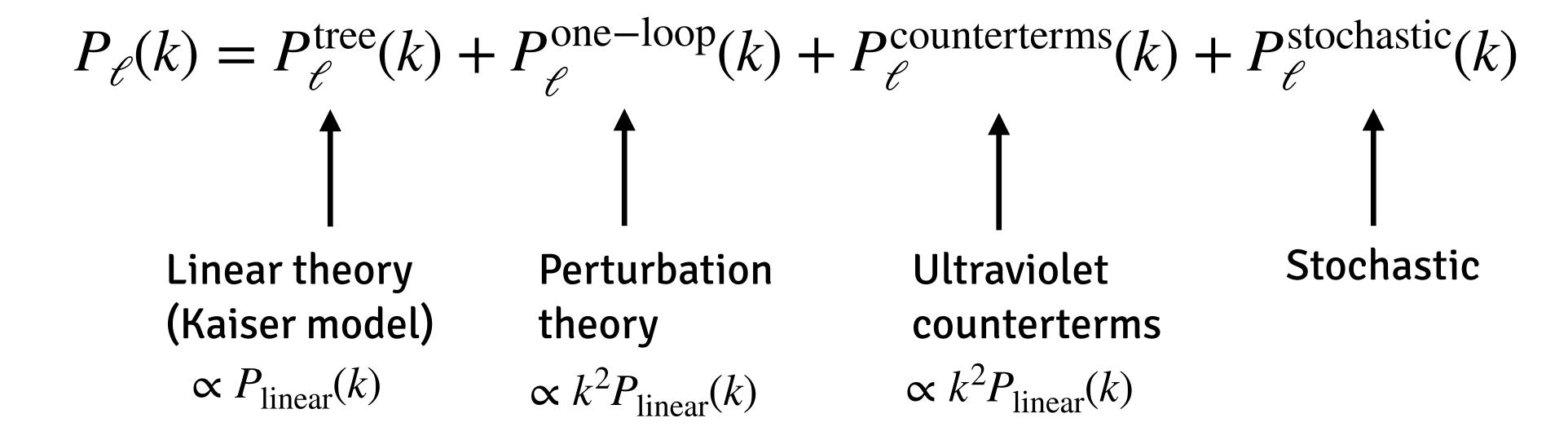
EFTofLSS approach:

[Baumann, Nicolis, Senatore, Zaldarriaga 10]



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Two main codes:



[Ivanov, Chudaykin, Philcox, Simonovic 20]



PyBird

[Zhang, D'Amico, Senatore 20]

CMASS NGC (monopole) $kP_{0,\Lambda CDM}(k) [h/{\rm Mpc}]^2$ Linear prediction (CLASS) EFT prediction at 1-loop (PyBird) 0.20 0.000.05 0.10 0.15 $k[h/\mathrm{Mpc}]$

[By Théo Simon]

The success of EFTofLSS

The EFTofLSS has been successfully applied to BOSS data

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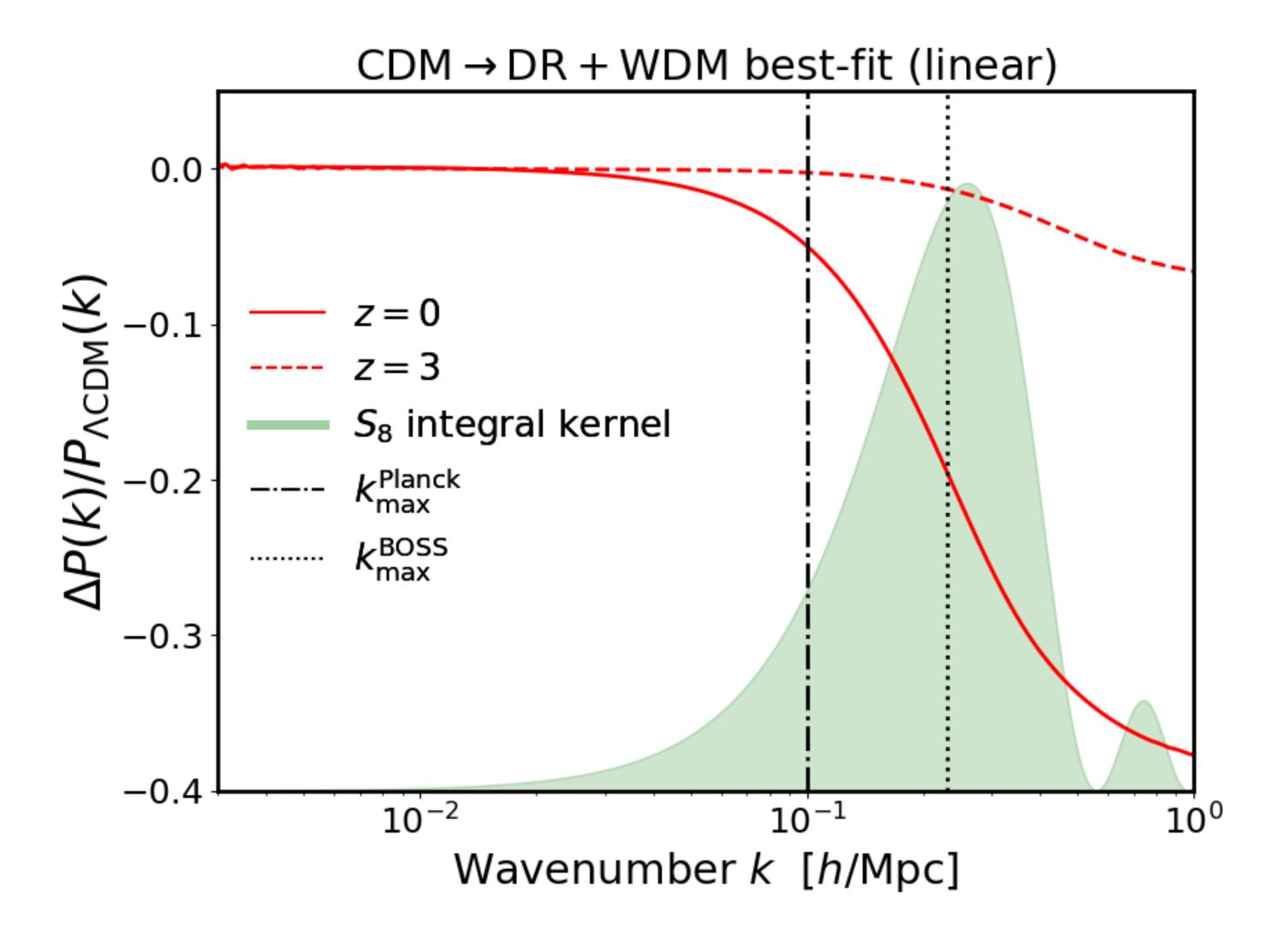
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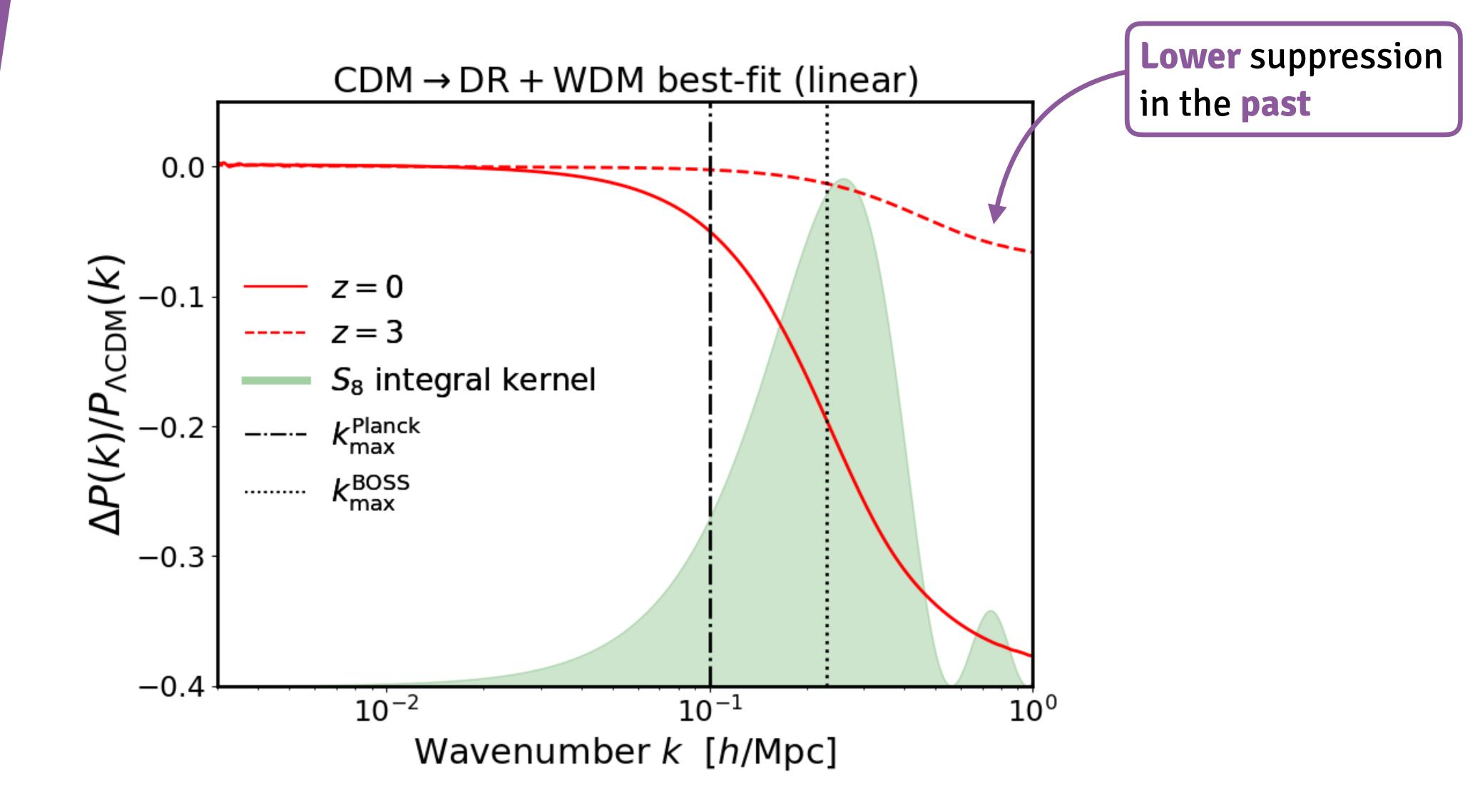
[Colas+ 19] [D'Amico+ 19]

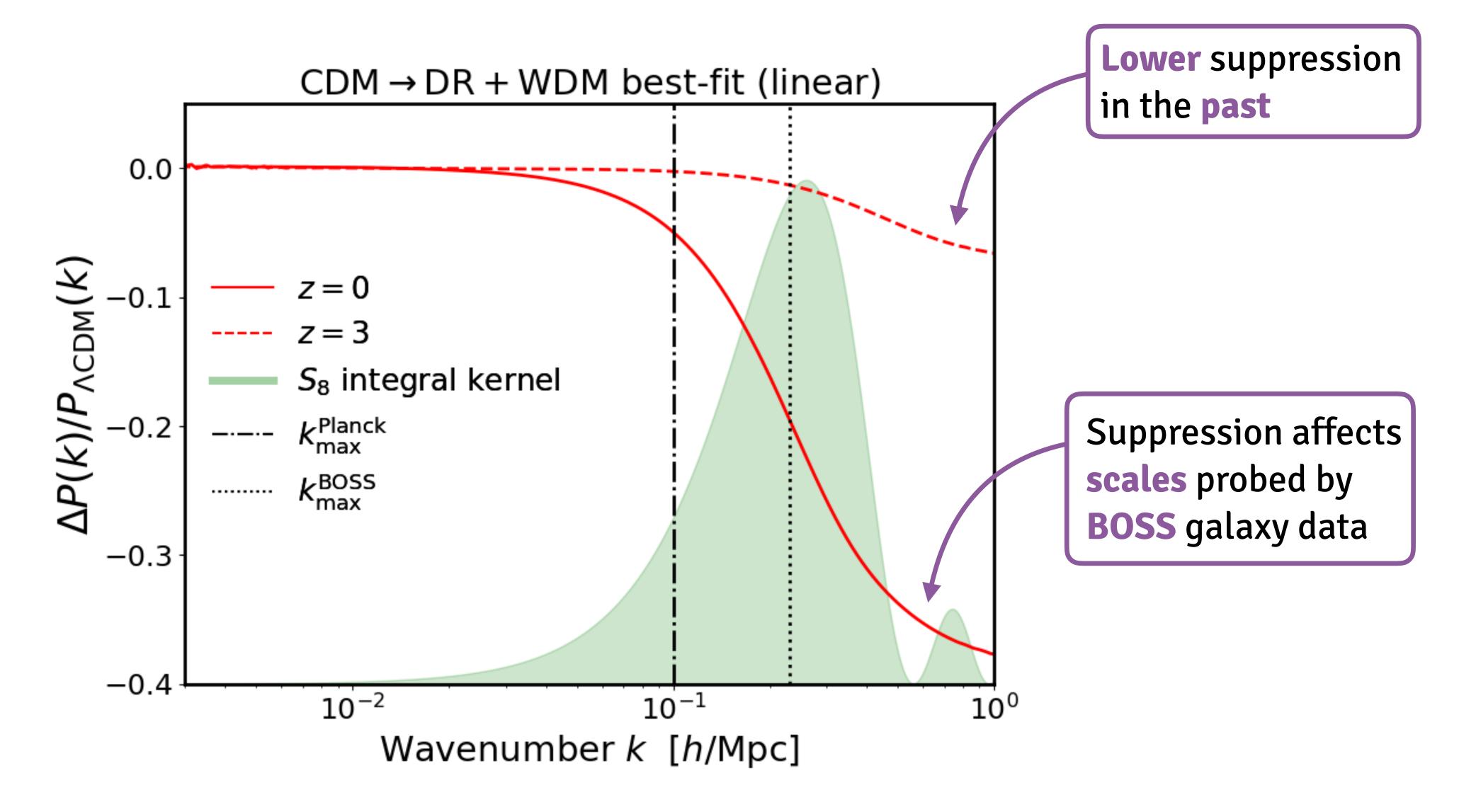
Can we use this to test DDM?

Part III:

CONSTRAINTS ON DDM FROM EFTofLSS





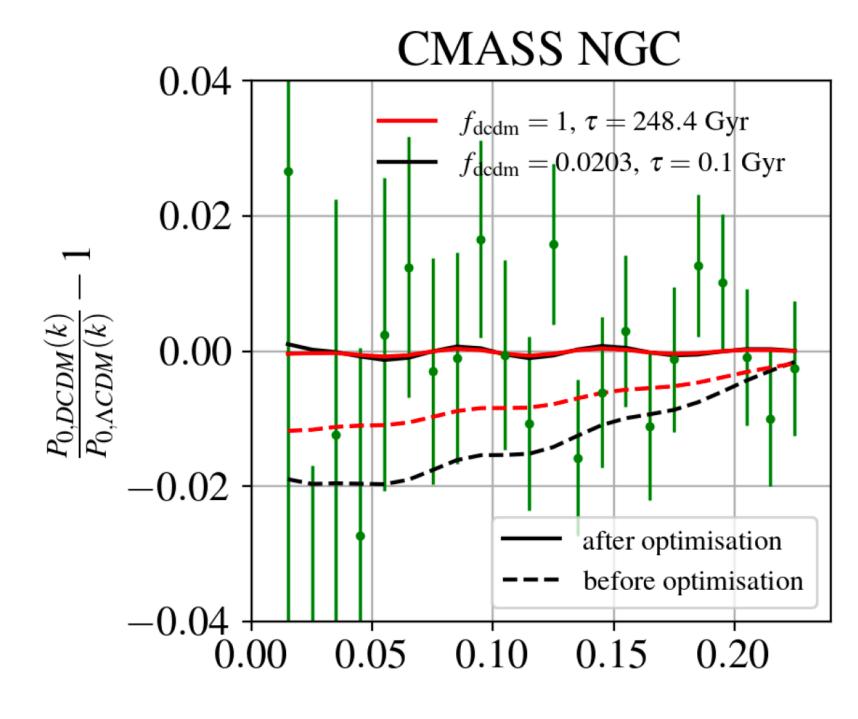


Test CDM → DR and CDM → DR + WDM w/ our modified CLASS version + PyBird

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Beware of nuisance EFT parameters!

Can be degenerate with effects of DDM



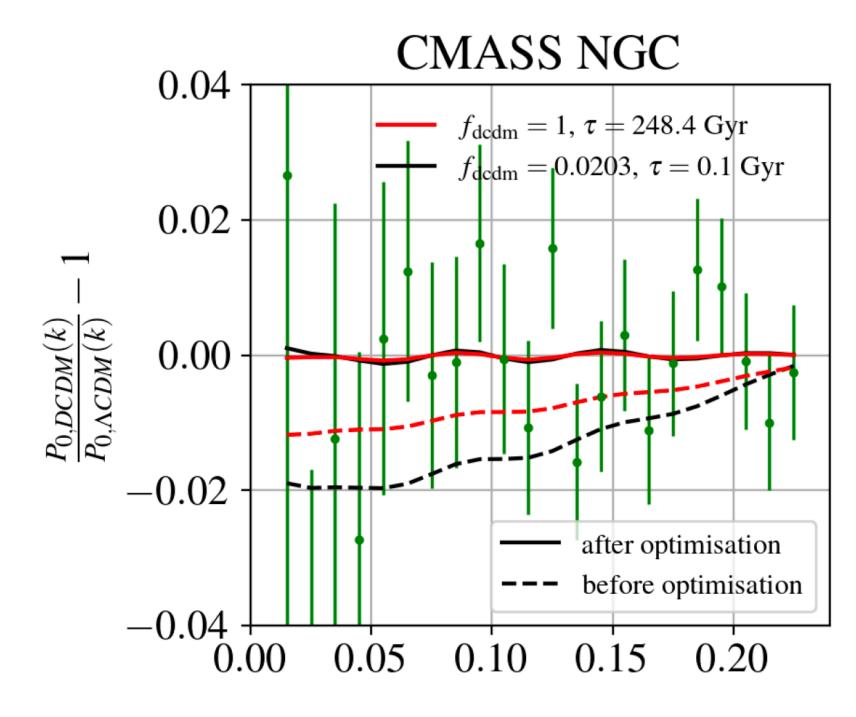
[Simon, Abellán+ 22]

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Compare results with and without EFT



[Simon, Abellán+ 22]

- Planck + Pantheon + BOSS BAO/f σ_8 + Ext-BAO Planck + Pantheon + EFTofBOSS + Ext-BAO
- 0.030.01 0.84 $\overset{\infty}{\backsim}$ 0.82 0.80 0.01 $0.02 \quad 0.03$ Γ/Gyr^{-1} S_8 $f_{
 m dcdm}$

[Simon, Abellán+ 22]

Results for CDM → **DR**

EFTofBOSS doesn't improve the constraints significantly

- Planck + Pantheon + BOSS BAO/f σ_8 + Ext-BAO

 Planck + Pantheon + EFTofBOSS + Ext-BAO
- 0.03 $f_{\text{qcdm}}^{\text{m}}$ 0.010.84 \mathcal{S}^{∞} 0.82 0.80 0.01 $0.02 \quad 0.03$ Γ/Gyr^{-1} S_8 $f_{
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Results for CDM → **DR**

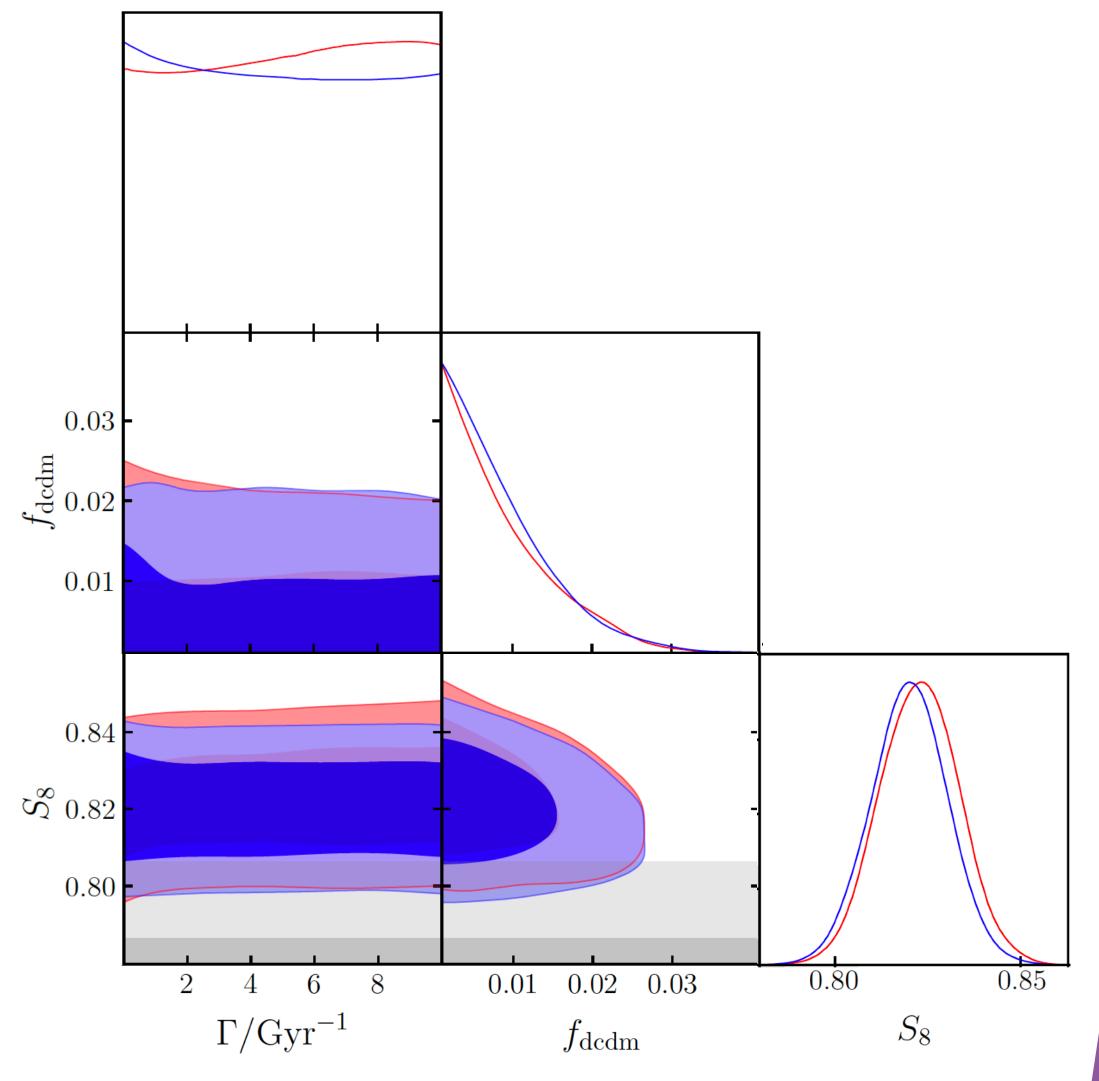
EFTofBOSS doesn't improve the constraints significantly

Most up-to-date constraints

$$\tau > 250 \text{ Gyr} \quad \text{for } f_{\text{dcdm}} = 1$$

$$f_{\rm dcdm} < 2.16\%$$
 for $\tau < t_U$

Planck + Pantheon + BOSS BAO/f σ_8 + Ext-BAO Planck + Pantheon + EFTofBOSS + Ext-BAO



[Simon, Abellán+ 22]

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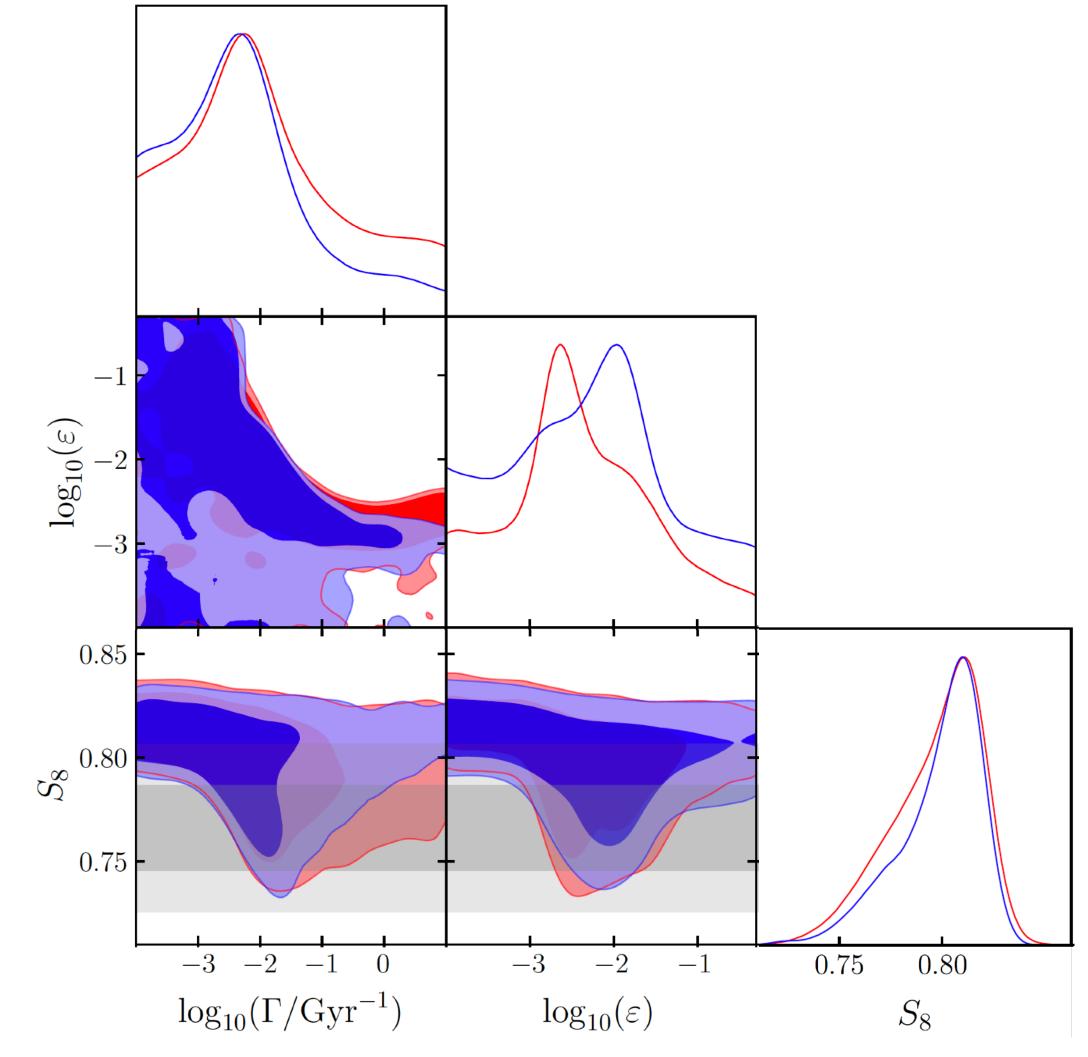
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We confirm that this model does **not resolve** the S₈ tension

- Planck + Pantheon + BOSS BAO/f σ_8 + Ext-BAO + S_8
- Planck + Pantheon + EFTofBOSS + Ext-BAO + S_8

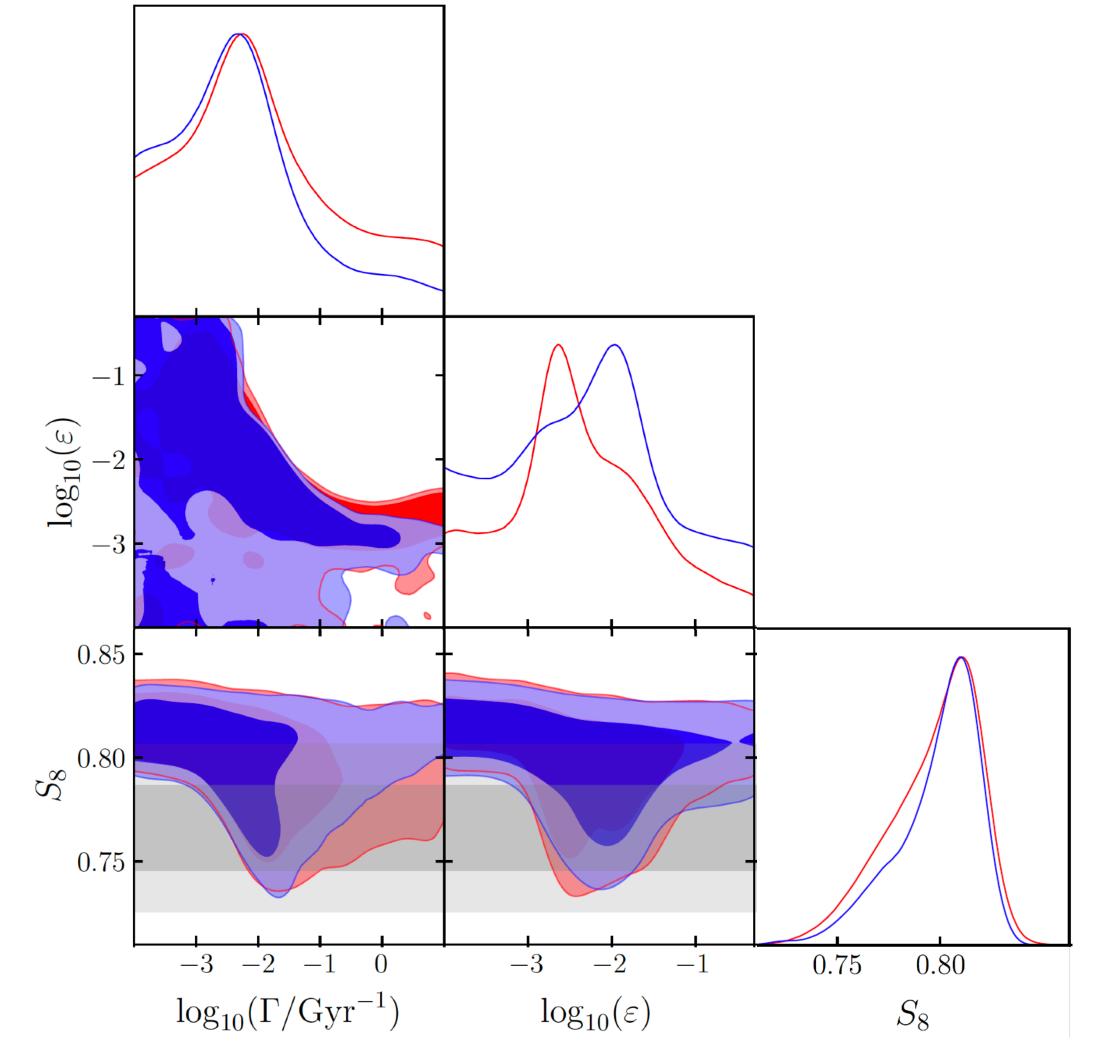


[Simon, Abellán+ 22]

Results for CDM → **DR** + **WDM**

This model can still resolve the S₈ tension

- Planck + Pantheon + BOSS BAO/f σ_8 + Ext-BAO + S_8
- Planck + Pantheon + EFTofBOSS + Ext-BAO + S_8



[Simon, Abellán+ 22]

Results for CDM → **DR** + **WDM**

- This model can still resolve the S₈ tension
- EFTofBOSS improves constraints on the lifetime

$$1.3 < \log_{10}(\tau/\text{Gyr}) < 3.8$$
 without EFT

$$1.6 < \log_{10}(\tau/\text{Gyr}) < 3.7 \text{ with EFT}$$

- Planck + Pantheon + BOSS BAO/f σ_8 + Ext-BAO + S_8 Planck + Pantheon + EFTofBOSS + Ext-BAO + S_8
- $\log_{10}(\varepsilon)$ 0.85 $\overset{\infty}{\sim} 0.80$ 0.75-2 -10.80 $\log_{10}(\Gamma/\mathrm{Gyr}^{-1})$ $\log_{10}(\varepsilon)$ S_8

[Simon, Abellán+ 22]

Results for CDM → **DR** + **WDM**

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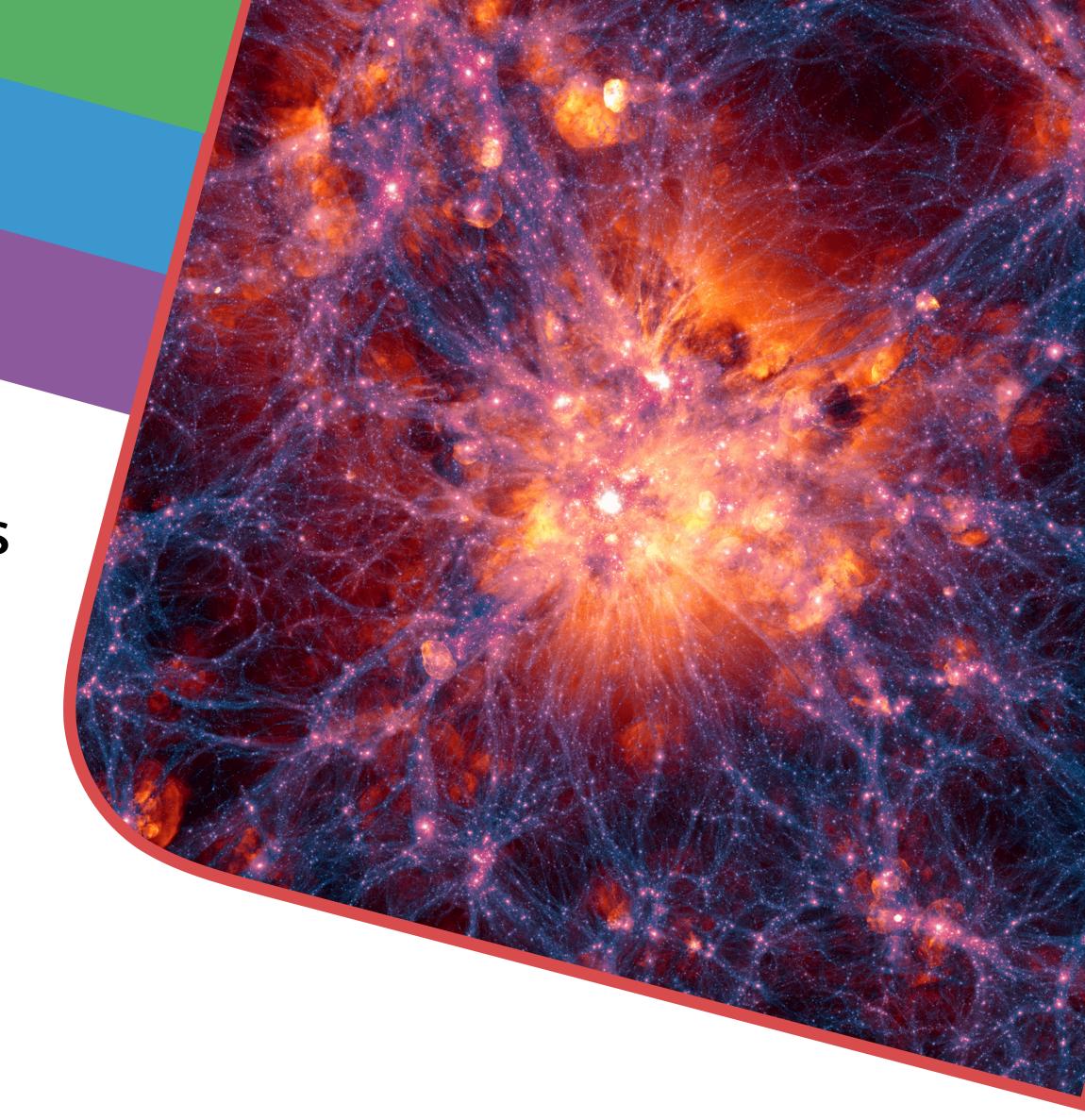
It changes the best-fit

$$\tau = 55 \text{ Gyr} \longrightarrow \tau = 120 \text{ Gyr}$$

$$\varepsilon = 0.7\% \longrightarrow \varepsilon = 1.2\%$$

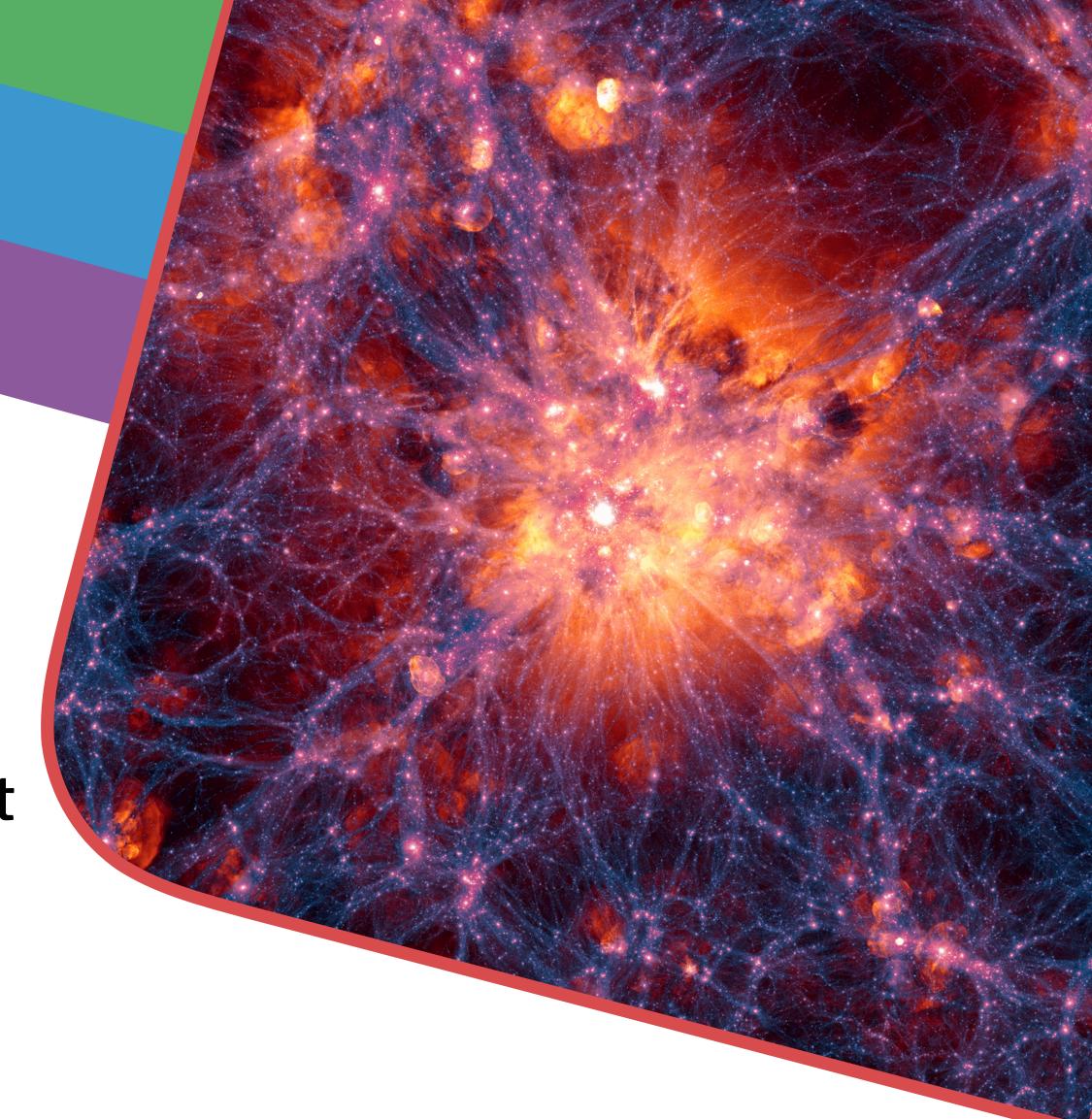
Conclusions

EFTofLSS applied to BOSS data can have **constraining power** on ΛCDM extensions



Conclusions

- **EFTofLSS** applied to BOSS data can have **constraining power** on ΛCDM extensions
- We derived most up-to-date constraints on two DDM scenarios, which are relevant for model building and the S₈ tension



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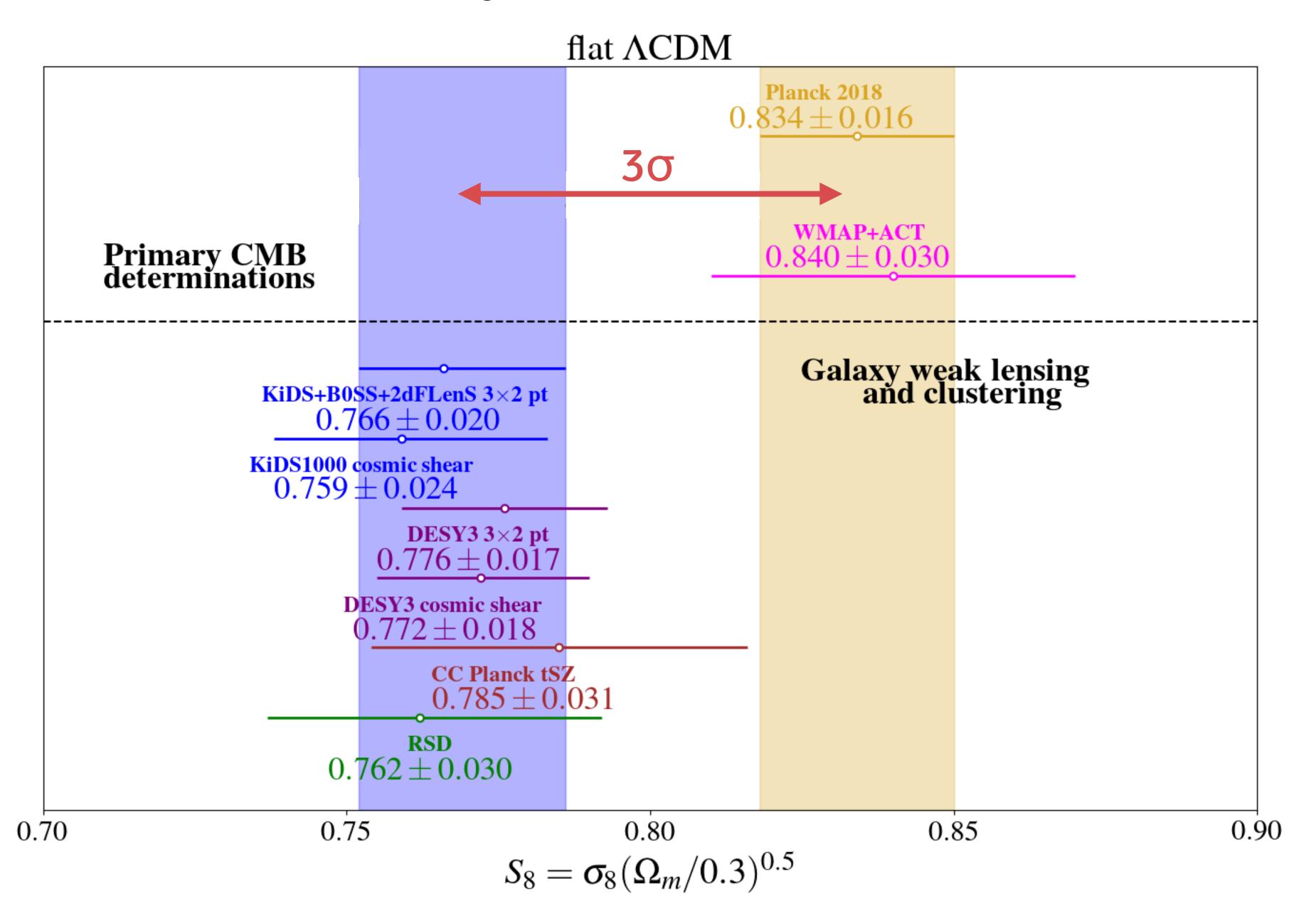
EFTofLSS applied to BOSS data can have **constraining power** on ΛCDM extensions

We derived most up-to-date constraints on two DDM scenarios, which are relevant for model building and the S₈ tension

Future LSS data (at small z) will probe further the CDM \rightarrow DR + WDM model that solves the S₈ tension

BACK-UP

A closer look at the S_8 tension



Evolution of DDM perturbations

Track δ_i , θ_i and σ_i for i = CDM, DR, WDM

- Boltzmann hierarchy of eqs., dictate evolution of p.s.d. multipoles δf_{ℓ} (q, k, T)
 - For DM and DR, momentum d.o.f. are integrated out
 - For WDM, need to follow full evolution in phase space Computationally prohibitive, $\mathcal{O}(10^8)$ ODEs to solve!

New fluid equations for the WDM species

Based on previous approximation for massive neutrinos

[Lesgourgues+ 11]

$$\delta'_{\text{wdm}} = -3aH(c_{\text{syn}}^2 - w)\delta_{\text{wdm}} - (1+w)\left(\theta_{\text{wdm}} + \frac{h'}{2}\right) + a\Gamma(1-\varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}(\delta_{\text{dm}} - \delta_{\text{wdm}})$$

$$\theta'_{\text{wdm}} = -aH(1 - 3c_a^2)\theta_{\text{wdm}} + \frac{c_{\text{syn}}^2}{1 + w}k^2\delta_{\text{wdm}} - k^2\sigma_{\text{wdm}} - a\Gamma(1 - \varepsilon)\frac{\bar{\rho}_{\text{dm}}}{\bar{\rho}_{\text{wdm}}}\frac{1 + c_a^2}{1 + w}\theta_{\text{wdm}}$$

CPU time reduced from
~ 1 day to ~ 1 minute!

Other LSS probes to test CDM → DR + WDM

Reduction in the abundance of subhalos, can be constrained by observations of MW satellites

[DES 22]

Model well compatible with Lyman-a forest data, given time-dependence of power suppression

[Fuss, Garny 22]

Can also be probed by looking at **abundance of clusters** detected with the **Sunyaev Zel'dovich** effect

[Tanimura+ 23]

EFTofLSS parameters

$$\begin{split} &P(k,\mu) = Z_{1}(\mu)^{2} P_{11}(k) \\ &+ 2 \int \frac{d^{3}q}{(2\pi)^{3}} \; Z_{2}(\mathbf{q},\mathbf{k}-\mathbf{q},\mu)^{2} P_{11}(|\mathbf{k}-\mathbf{q}|) P_{11}(q) + 6 Z_{1}(\mu) P_{11}(k) \int \frac{d^{3}q}{(2\pi)^{3}} \; Z_{3}(\mathbf{q},-\mathbf{q},\mathbf{k},\mu) P_{11}(q) \\ &+ 2 Z_{1}(\mu) P_{11}(k) \left(\mathbf{c}_{\text{ct}} \frac{k^{2}}{k_{\text{M}}^{2}} + \mathbf{c}_{r,1} \mu^{2} \frac{k^{2}}{k_{\text{M}}^{2}} + \mathbf{c}_{r,2} \mu^{4} \frac{k^{2}}{k_{\text{M}}^{2}} \right) + \frac{1}{\bar{n}_{g}} \left(\mathbf{c}_{\epsilon,0} + \mathbf{c}_{\epsilon,1} \frac{k^{2}}{k_{\text{M}}^{2}} + \mathbf{c}_{\epsilon,2} f \mu^{2} \frac{k^{2}}{k_{\text{M}}^{2}} \right) \end{split}$$

with

$$Z_{1}(\mathbf{q}_{1}) = K_{1}(\mathbf{q}_{1}) + f\mu_{1}^{2}G_{1}(\mathbf{q}_{1}) = b_{1} + f\mu_{1}^{2}$$

$$Z_{2}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mu) = K_{2}(\mathbf{q}_{1}, \mathbf{q}_{2}) + f\mu_{12}^{2}G_{2}(\mathbf{q}_{1}, \mathbf{q}_{2}) + \frac{1}{2}f\mu q \left(\frac{\mu_{2}}{q_{2}}G_{1}(\mathbf{q}_{2})Z_{1}(\mathbf{q}_{1}) + \text{perm.}\right)$$

$$Z_{3}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mathbf{q}_{3}, \mu) = K_{3}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mathbf{q}_{3}) + f\mu_{123}^{2}G_{3}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mathbf{q}_{3})$$

$$+ \frac{1}{3}f\mu q \left(\frac{\mu_{3}}{q_{3}}G_{1}(\mathbf{q}_{3})Z_{2}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mu_{123}) + \frac{\mu_{23}}{q_{23}}G_{2}(\mathbf{q}_{2}, \mathbf{q}_{3})Z_{1}(\mathbf{q}_{1}) + \text{cyc.}\right)$$

with

$$K_{1} = b_{1}$$

$$K_{2}(\mathbf{q}_{1}, \mathbf{q}_{2}) = b_{1} \frac{\mathbf{q}_{1} \cdot \mathbf{q}_{2}}{q_{1}^{2}} + b_{2} \left(F_{2}(\mathbf{q}_{1}, \mathbf{q}_{2}) - \frac{\mathbf{q}_{1} \cdot \mathbf{q}_{2}}{q_{1}^{2}} \right) + b_{4} + \text{perm.}$$

$$K_{3}(k, q) = \frac{b_{1}}{504k^{3}q^{3}} \left(-38k^{5}q + 48k^{3}q^{3} - 18kq^{5} + 9(k^{2} - q^{2})^{3} \log \left[\frac{k - q}{k + q} \right] \right)$$

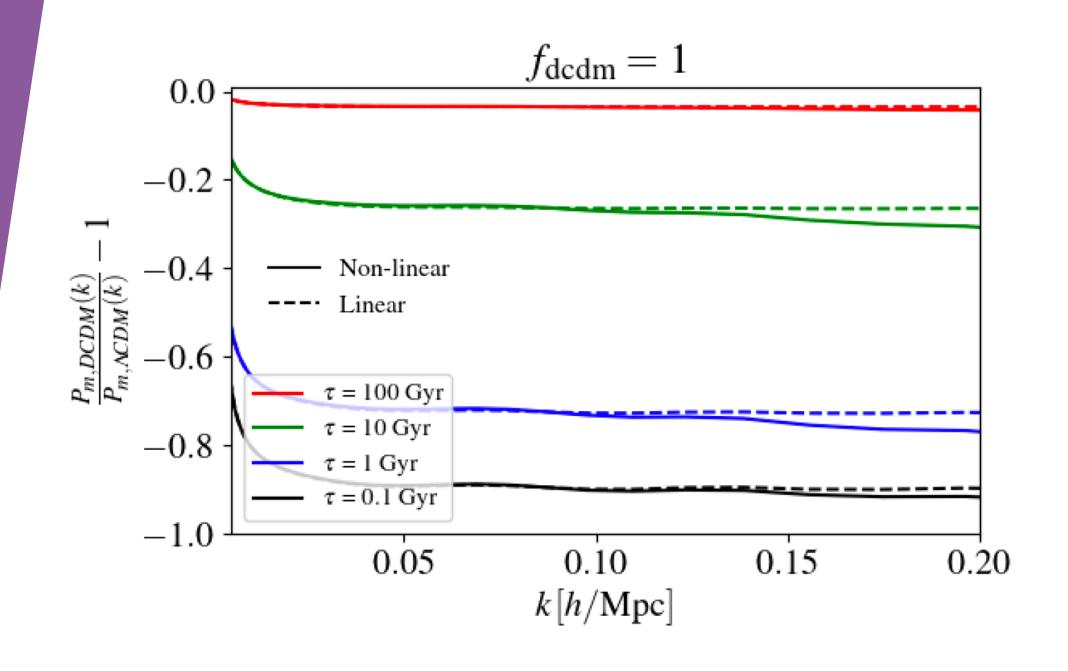
$$+ \frac{b_{3}}{756k^{3}q^{5}} \left(2kq(k^{2} + q^{2})(3k^{4} - 14k^{2}q^{2} + 3q^{4}) + 3(k^{2} - q^{2})^{4} \log \left[\frac{k - q}{k + q} \right] \right)$$

- 4 parameters to describe galaxy bias
- 3 parameters to describe counterterms
- 3 parameters to describe stochastic terms

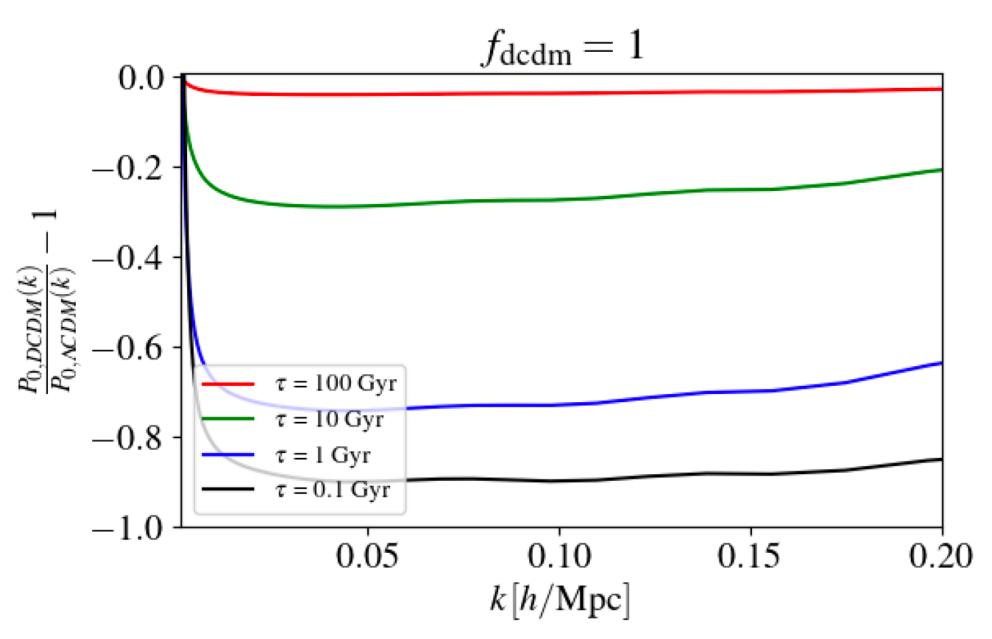
10 parameters in total, but 8 are analytically marginalized

Matter and galaxy power spectra for CDM → DR

Matter power spectrum



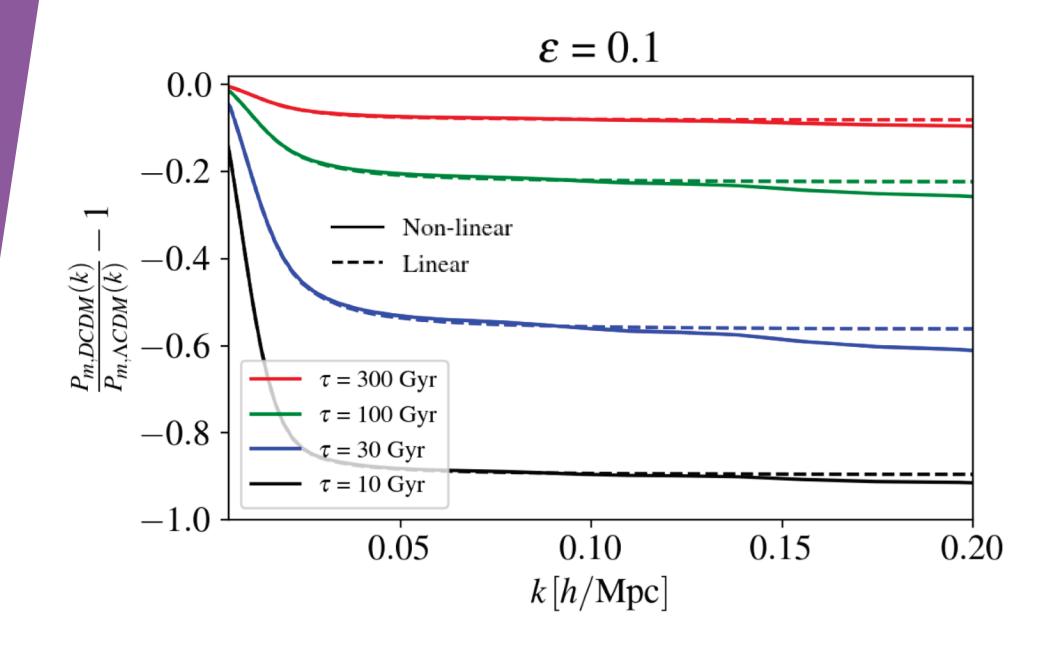
Galaxy spectrum (monopole)



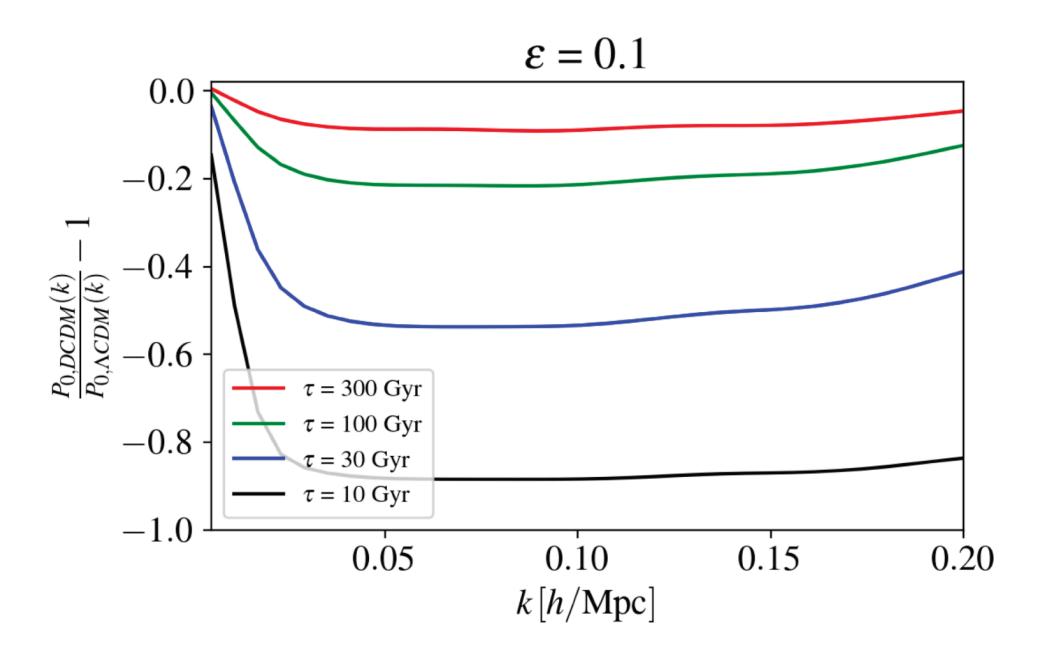
[Simon, Abellán+ 22]

Matter and galaxy power spectra for CDM → DR + WDM

Matter power spectrum



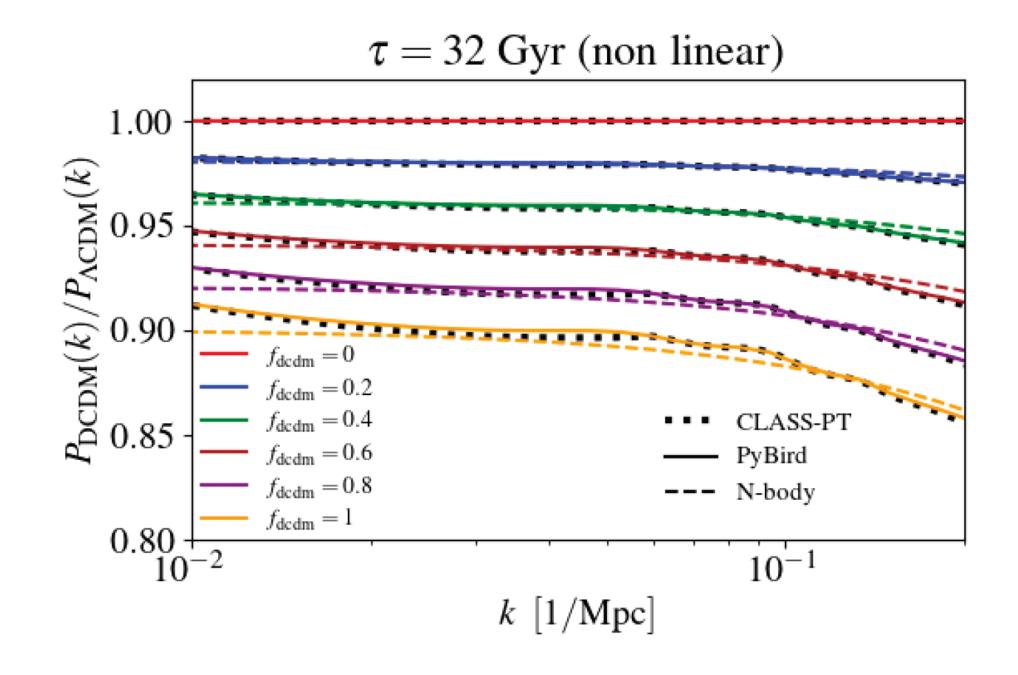
Galaxy spectrum (monopole)



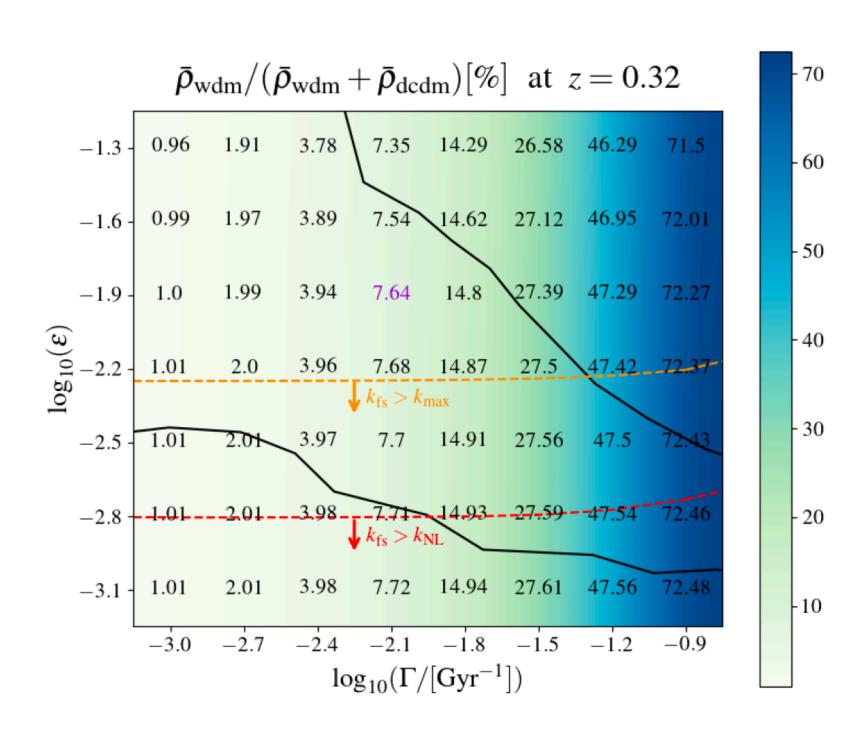
[Simon, Abellán+ 22]

Checking validity of EFTofLSS applied to DDM





CDM → DR + WDM



[Simon, Abellán+ 22]